

Supersymmetry after LHC Run 1



Brian Petersen
CERN

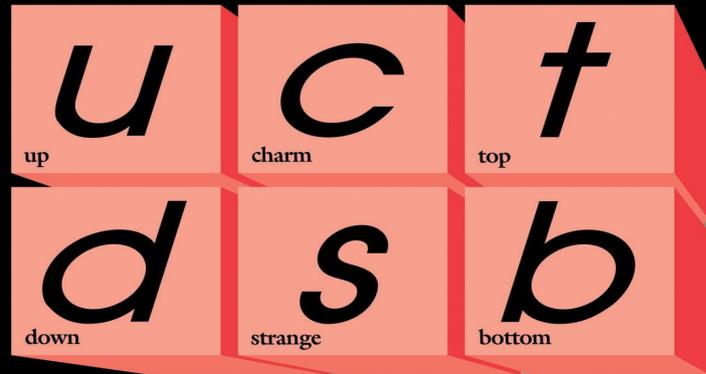
Outline

- Supersymmetry
What and why?
- Searching for Supersymmetry at the LHC
*Brief intro to LHC and the ATLAS experiment
and how we search for Supersymmetry there*
- Status of Run-1 Supersymmetry searches
Where did we look? What did we find?
Will focus mostly on some of the latest results
- Outlook for upcoming data-taking
What can we expect for 2015 and beyond

Supersymmetry

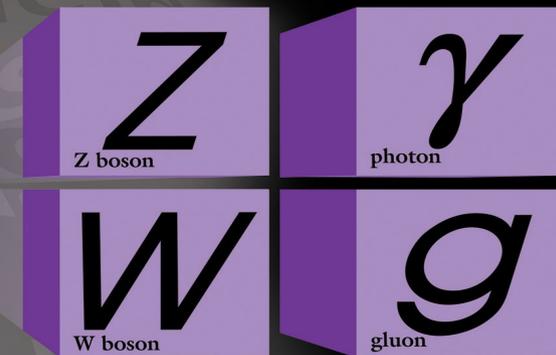
Standard Model

Quarks



Pre-LHC:
Is the Higgs Mechanism
responsible for masses?

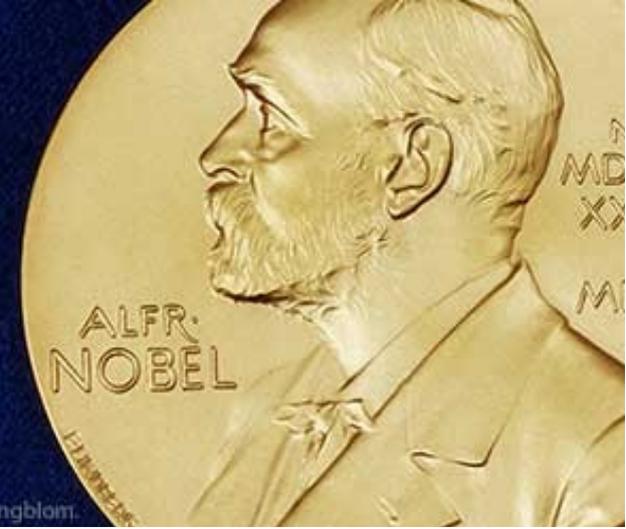
Forces



Leptons

2013 NOBEL PRIZE IN PHYSICS

François Englert Peter W. Higgs



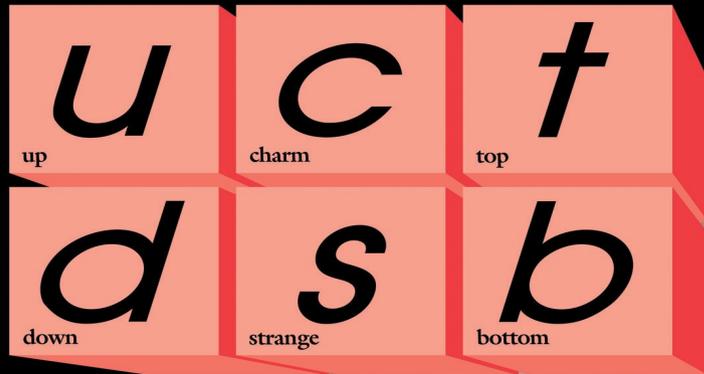
© © The Nobel Foundation, Photo: Lovisa Engblom.

The Nobel Prize in Physics 2013 was awarded jointly to François Englert and Peter W. Higgs

"for the theoretical discovery of a mechanism that contributes to our understanding of the origin of mass of subatomic particles, and which recently was confirmed through the discovery of the predicted fundamental particle, by the ATLAS and CMS experiments at CERN's Large Hadron Collider"

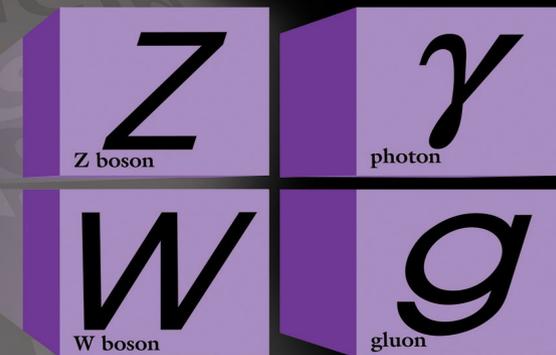
Standard Model Complete?

Quarks



Post-LHC run 1:
Is it the Standard Model
Higgs Boson?

Forces

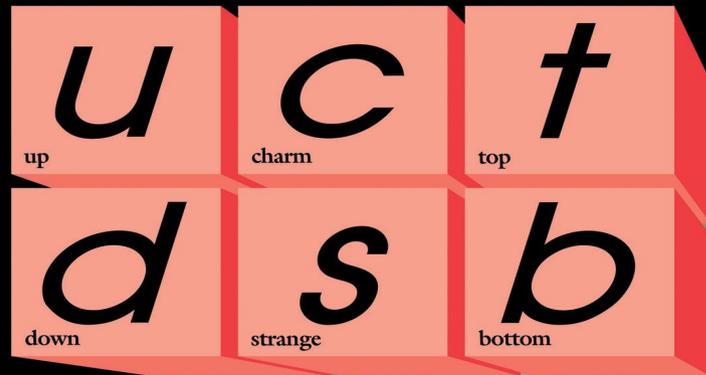


Leptons



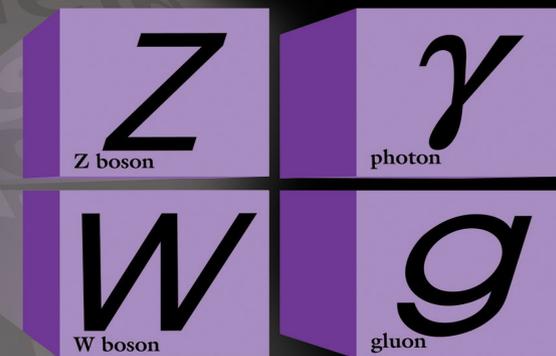
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Post-LHC run 1:
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H
Higgs boson



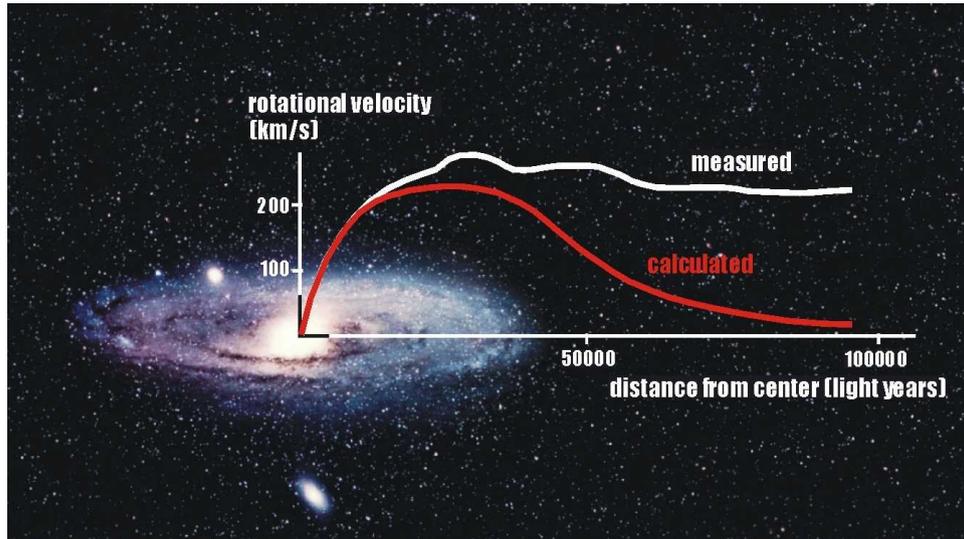
Leptons

Is there more than this?

Dark Matter

Strong evidence in astrophysics for presence of dark matter

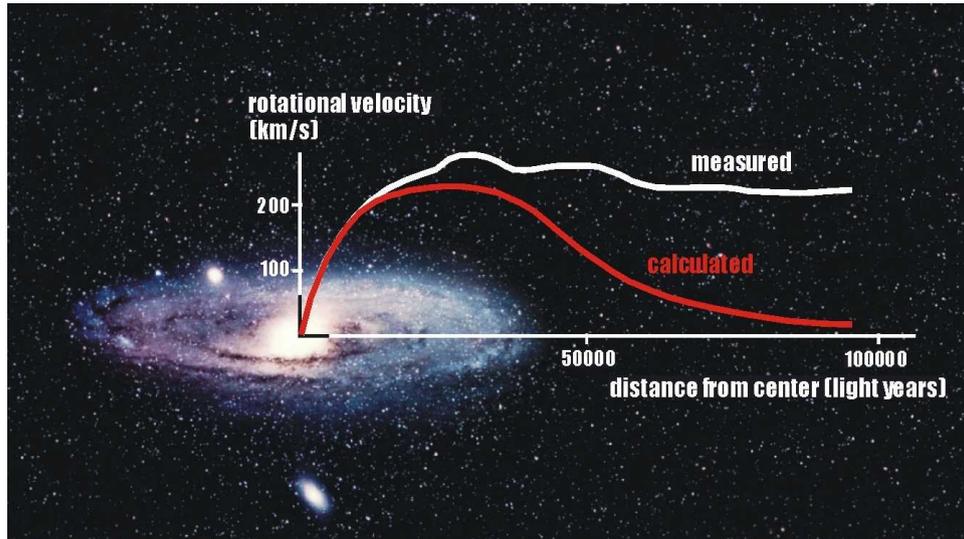
Galaxy rotational curves



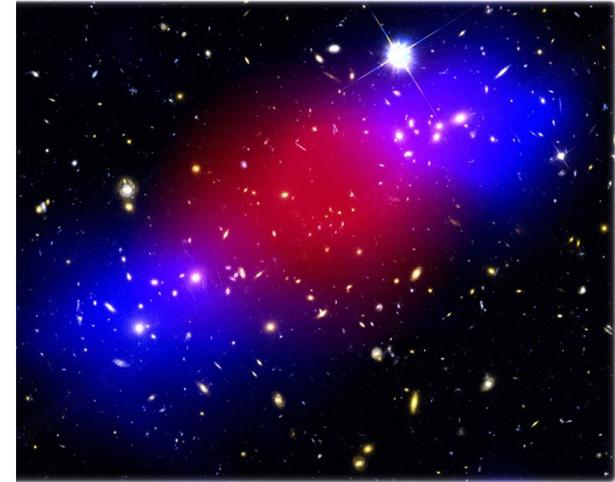
Dark Matter

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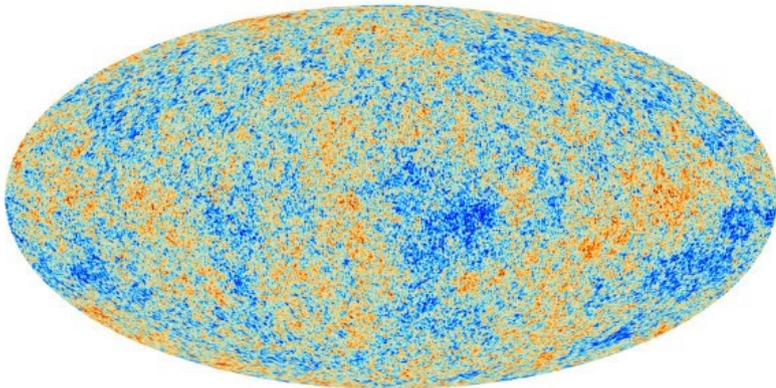
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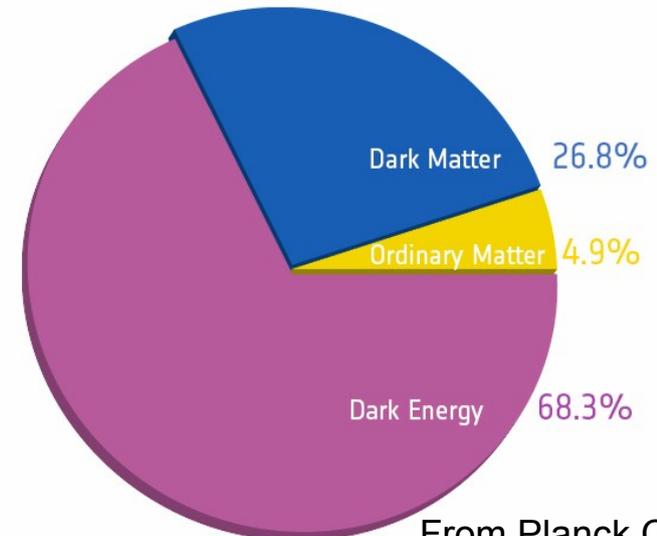
Gravitational lensing



Cosmic microwave background



Combined cosmological fit



From Planck Collaboration

Possible Sources of Dark Matter

Ordinary non-luminous matter?

- A lot of the 5% ordinary matter is not luminous, but present in form of gas etc.
- Cannot account for non-baryonic DM

Neutrinos?

- Can only have small contribution due to large-scale structure formation

Modifying gravity

- Very difficult to explain all the different measurements

Axions?

- Particle resulting from the Peccei-Quinn solution to the strong CP problem
- O(keV) mass – good DM candidate

WIMPs?

- Weakly Interacting Massive Particle
- EW-scale (100-1000 GeV) gives right DM density

Other exotic particles

- Many other options have been proposed

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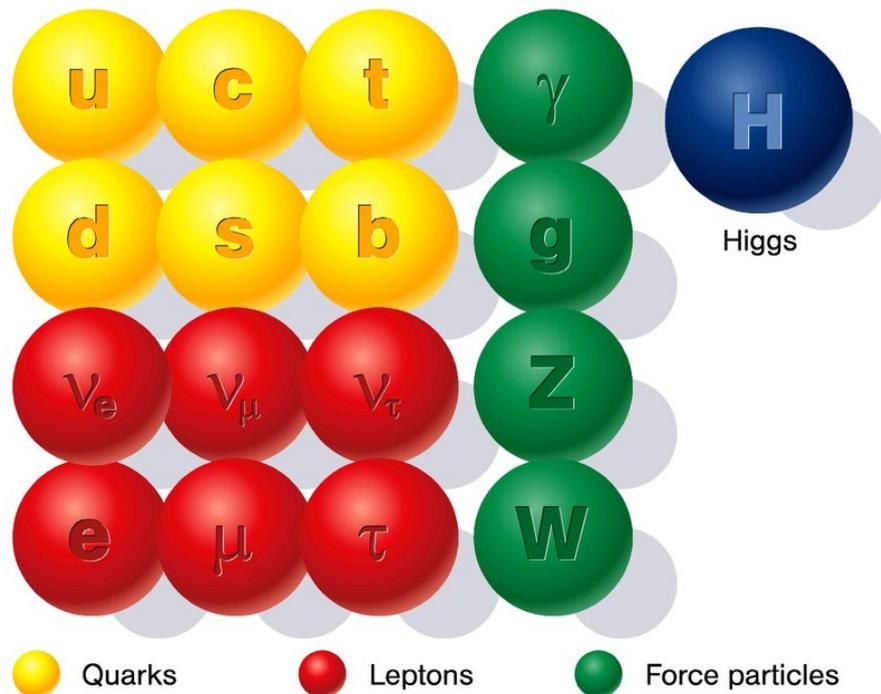
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Supersymmetry

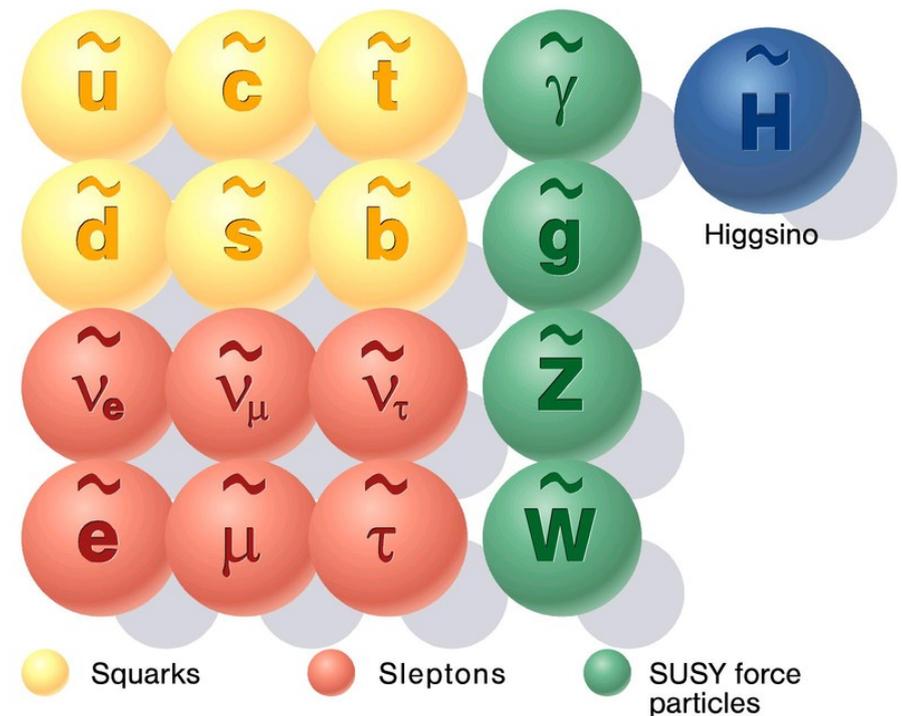
Supersymmetry (SUSY) is a favored source for WIMP

- New symmetry between bosons and fermions
- For every SM particle, introduces partner with $\Delta\text{spin}=\frac{1}{2}$

Standard particles



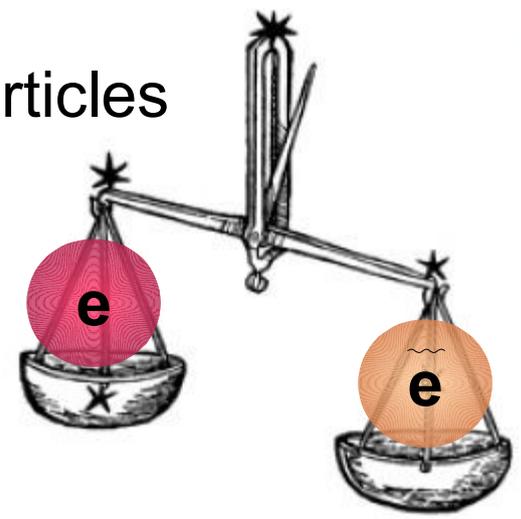
SUSY particles



Supersymmetry

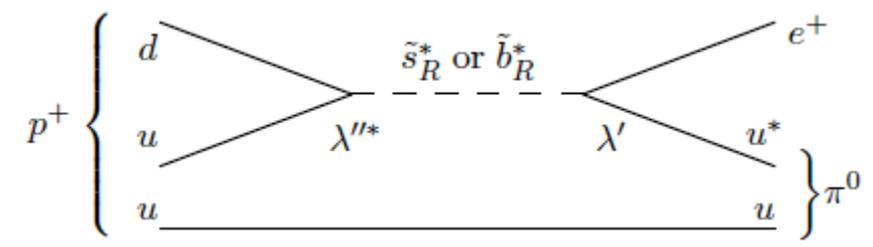
Supersymmetry must be broken symmetry

- If unbroken, would have same mass as SM particles
- Introduce soft breaking terms
- In minimal supersymmetry (MSSM) this introduces 124 new parameters



R-parity

- MSSM introduces baryon and lepton number violating processes
- Protect against proton decays by introducing new discrete symmetry



$$R = (-1)^{2S+3B+L} = +1 \text{ } (-1) \text{ for SM (SUSY) particles}$$

- Implies SUSY particles always produced in pairs
- Lightest SUSY particle (LSP) has to be stable

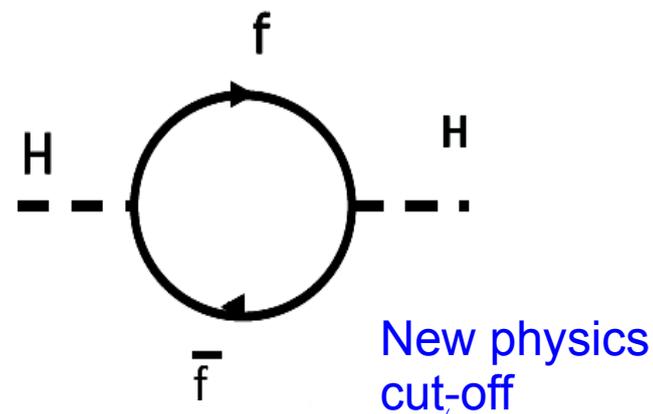
This is an excellent Dark Matter candidate

Supersymmetry Motivation

Dark matter is not the original or only motivation for SUSY
 Also provides solutions to several particle theory issues

Hierarchy Problem

- Very large loop corrections to Higgs mass in any SM extension



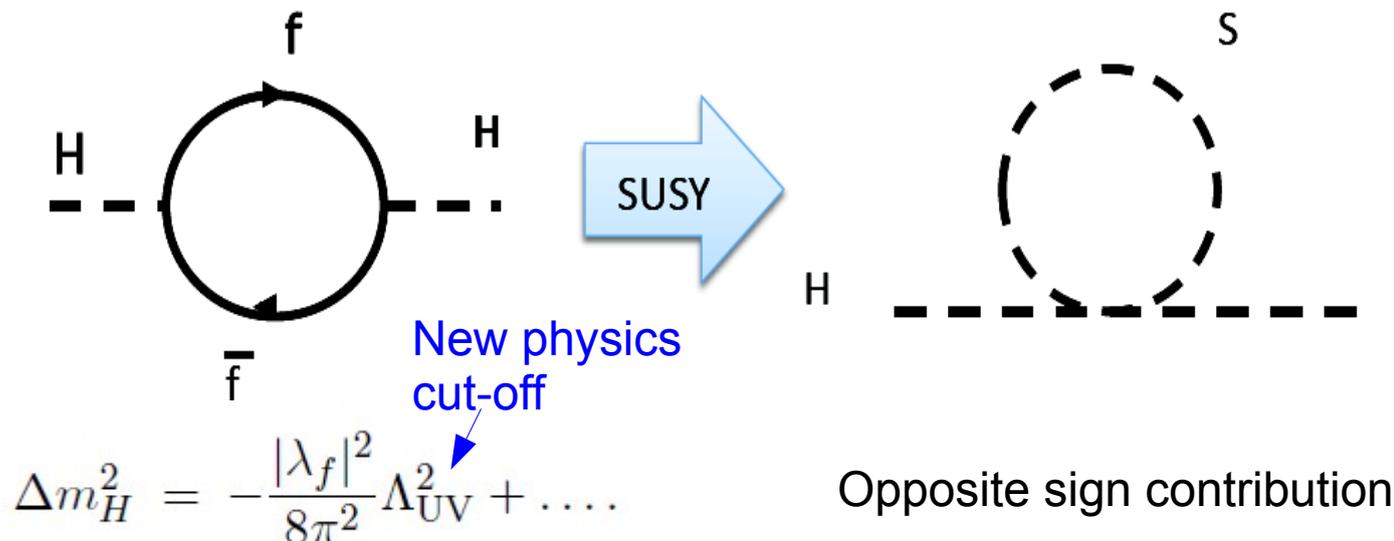
$$\Delta m_H^2 = -\frac{|\lambda_f|^2}{8\pi^2} \Lambda_{UV}^2 + \dots$$

Supersymmetry Motivation

Dark matter is not the original or only motivation for SUSY
Also provides solutions to several particle theory issues

Hierarchy Problem

- Very large loop corrections to Higgs mass in any SM extension
- SUSY can stabilize Higgs boson mass without extreme fine-tuning



Requires SUSY particles not to be much heavier, $O(\text{TeV})$, than the Standard Model particles

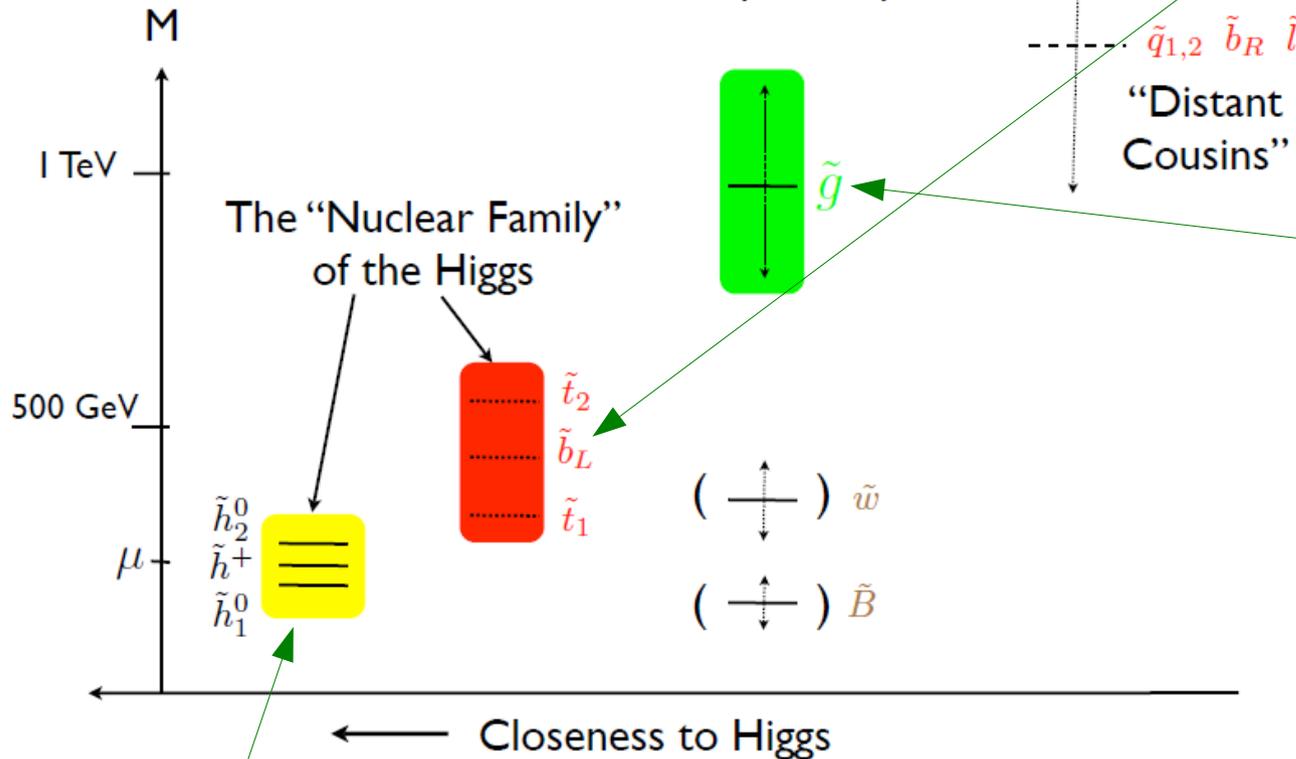
Natural SUSY

Not all SUSY particles need to be light
Focus on the ones that have to be light (“Natural SUSY”)

L. Hall, LBNL workshop, Oct 2011

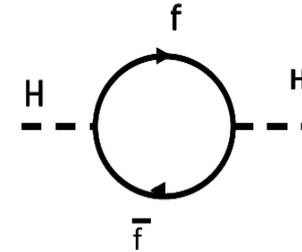
A Natural Spectrum

General “bottom-up” viewpoint



Light higgsinos as $\mu \sim m_z$ at tree level

Only light stop needed to regularize light Higgs boson



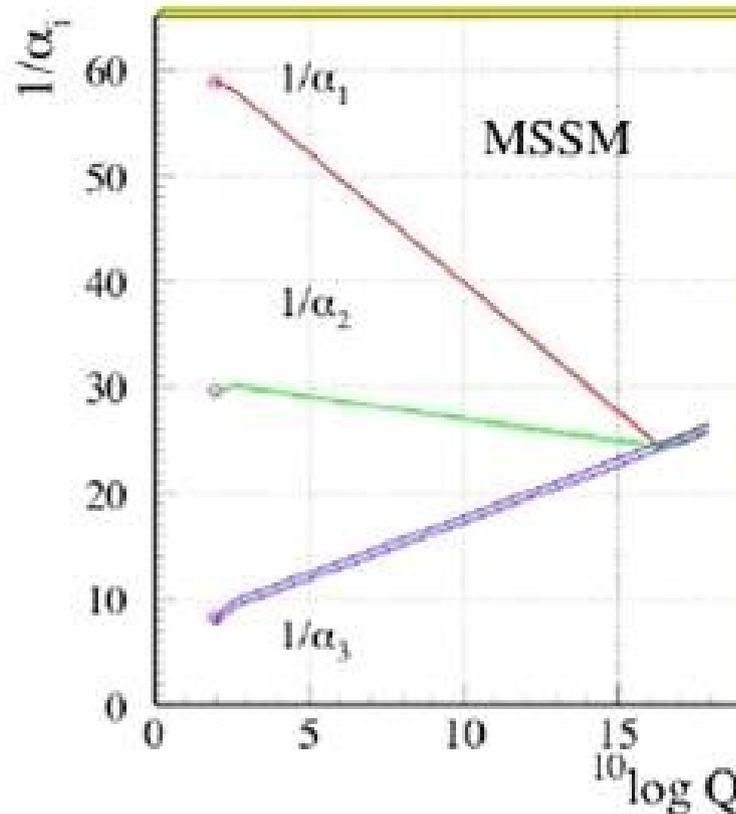
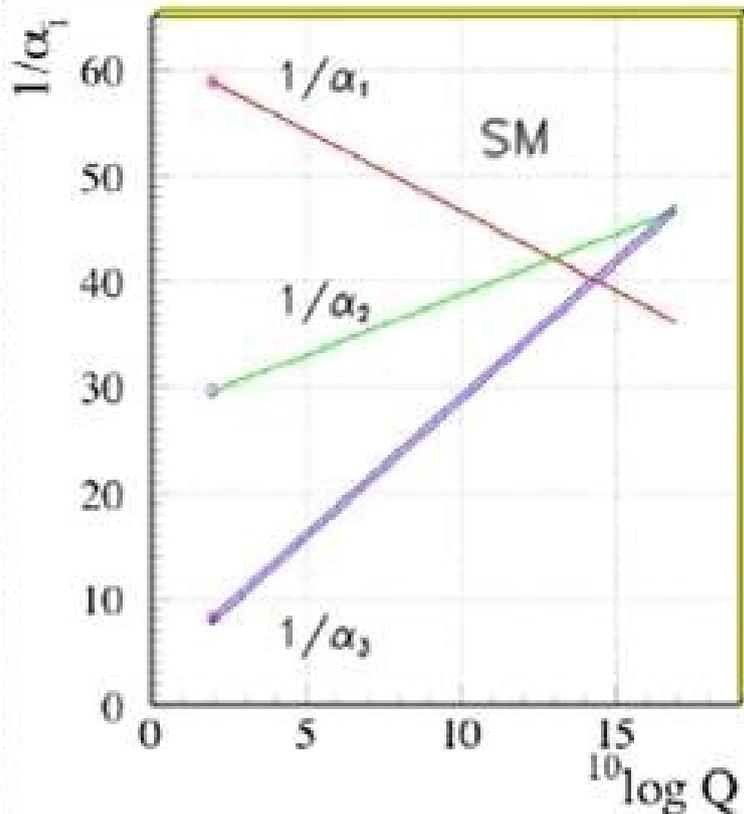
The gluino should not be too heavy either

Much attention has been devoted to the search for these in the LHC searches

Supersymmetry Motivation

Gauge-coupling Unification

- Couplings run in SM, but do not quite meet as one would expect in a Grand Unified Theory (GUT)
- SUSY changes running of couplings to allow unification of couplings at GUT scale

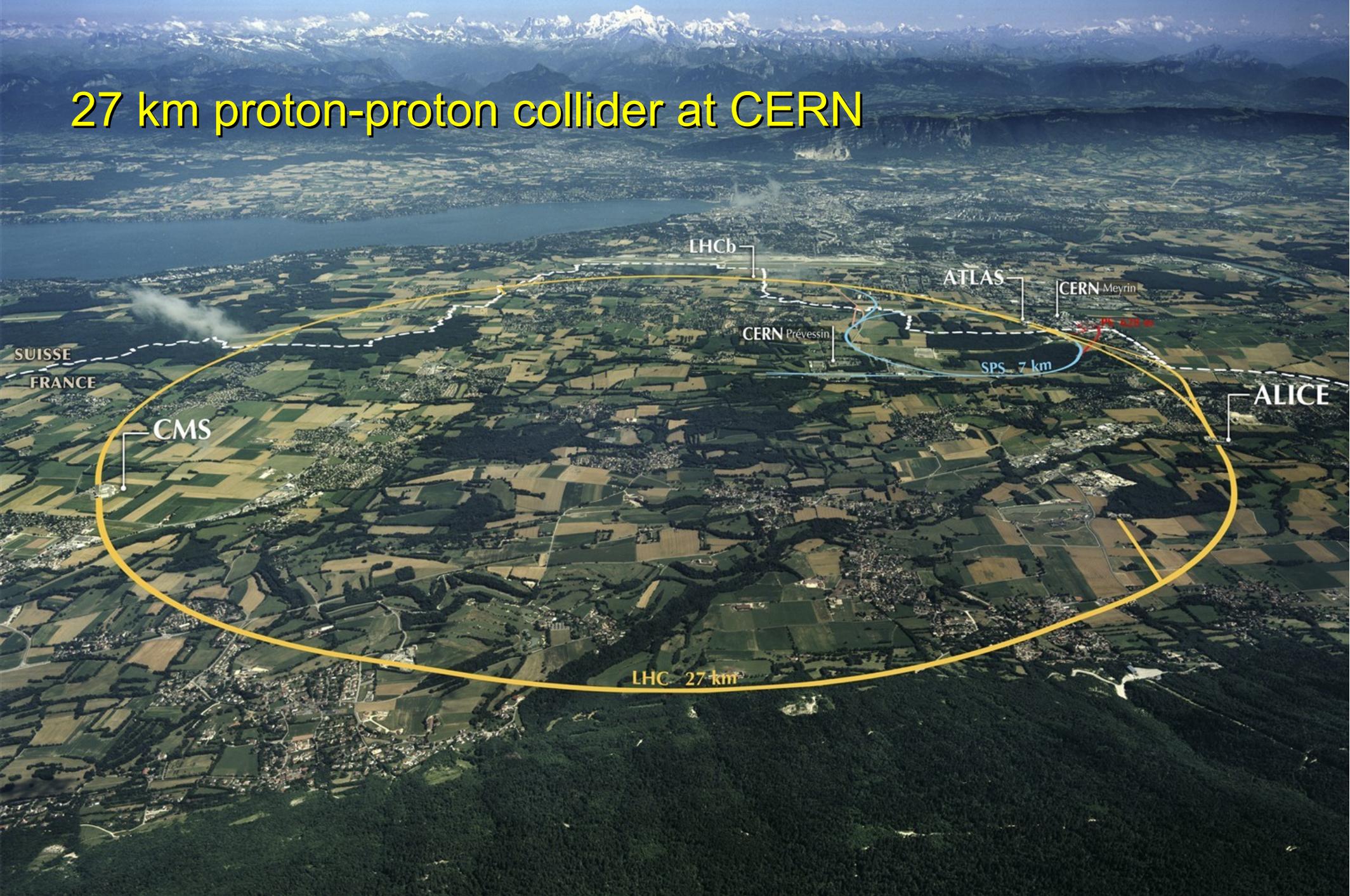


Again requires
SUSY at TeV
scale

Searching for Supersymmetry at the LHC

Large Hadron Collider

27 km proton-proton collider at CERN



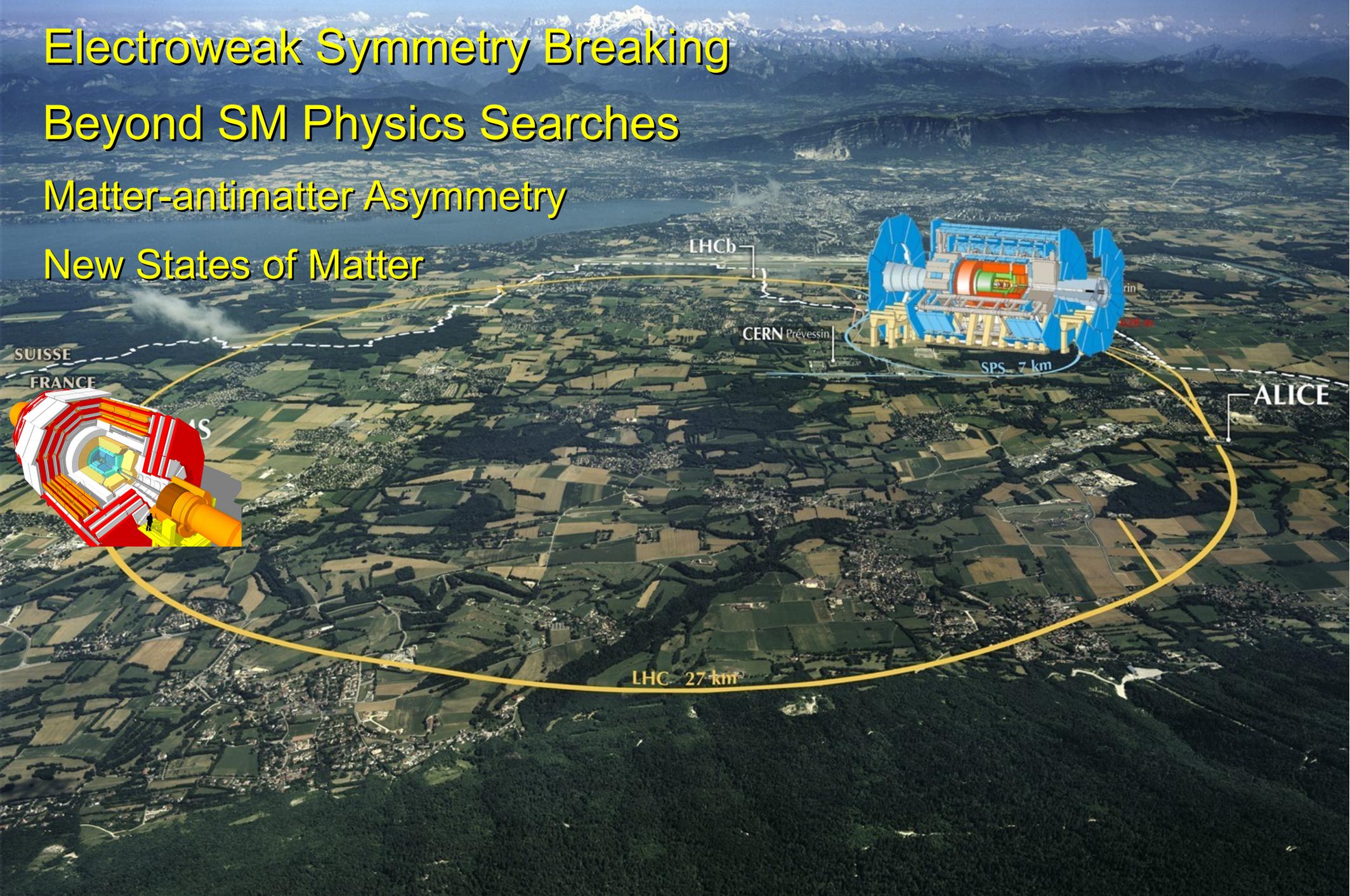
Large Hadron Collider Goals

Electroweak Symmetry Breaking

Beyond SM Physics Searches

Matter-antimatter Asymmetry

New States of Matter



LHCb

CERN Prévessin

SPS 7 km

ALICE

LHC 27 km

SUISSE
FRANCE

S

The ATLAS Experiment

Inner Detector

- Pixel (pixel detector)
- SCT (silicon strip detector)
- TRT (transition radiation tracker)

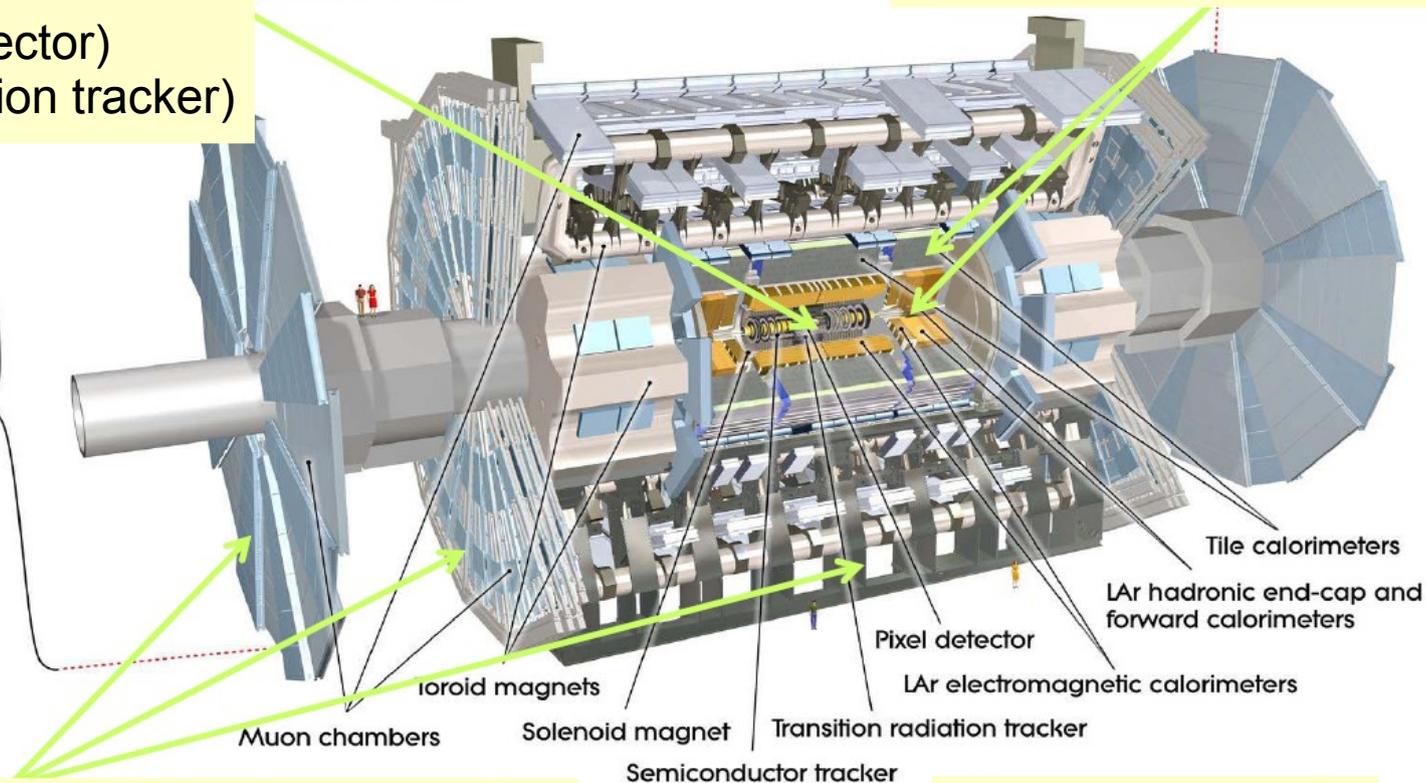
Calorimeter

- LAr (EM calorimeter)
- Tile (Fe/Scintillator tile)

7 TeV proton

25m

44m



Muon Spectrometer

- MDT,CSC (precise momentum measurement)
- RPC,TGC (trigger chambers)

Magnet System

- 2 T solenoid
- 0.5 T toroid

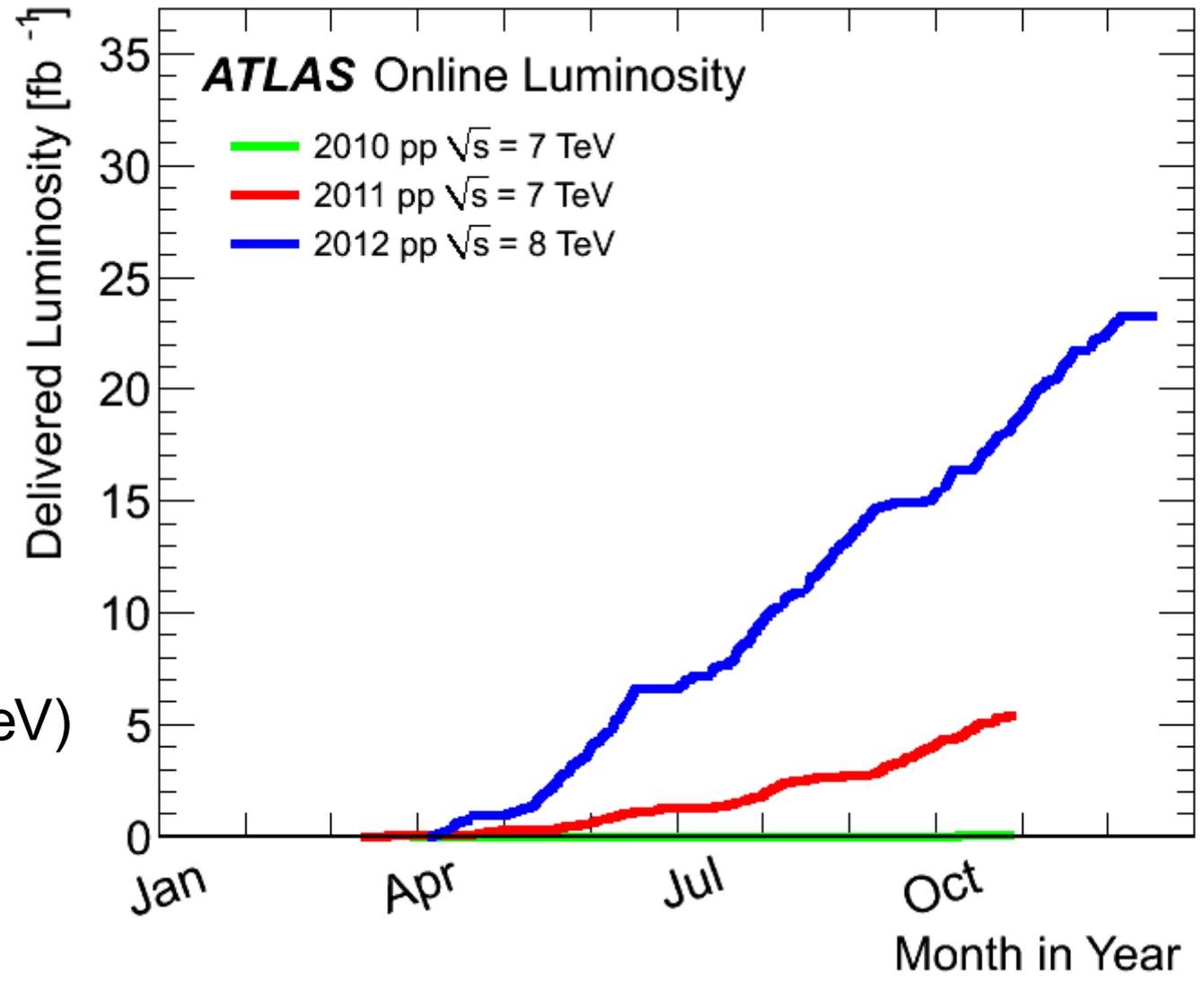
Luminosity

Critical LHC parameter is delivered luminosity:

$$\text{Rate} = \text{Cross-section} \times \text{Efficiency} \times \text{Luminosity}$$

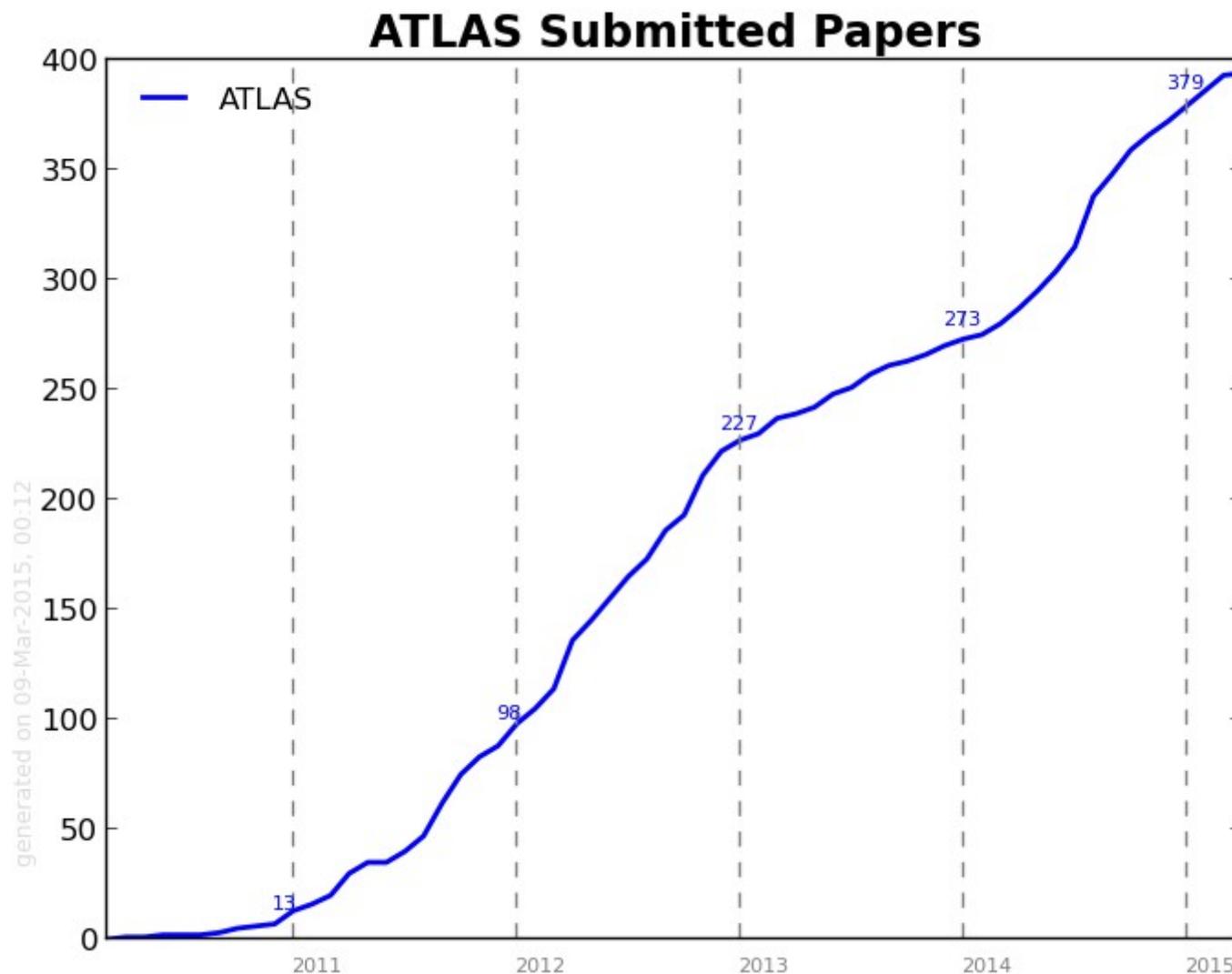
Excellent LHC performance:

- Peak luminosity of $7.7 \times 10^{33} \text{ cm}^{-2}\text{s}^{-1}$
- Big jumps in integrated luminosity:
2010: 0.048 fb^{-1}
2011: 5.6 fb^{-1}
2012: 23.3 fb^{-1}
(also increase in energy from 7 to 8 TeV)
- Full sample is called "Run-1"



LHC – A Paper Factory

~400 papers for the first 3 years of data-taking
In addition more than 600 preliminary results



ATLAS Total: 395

on SUSY: 74

CMS Total: 378

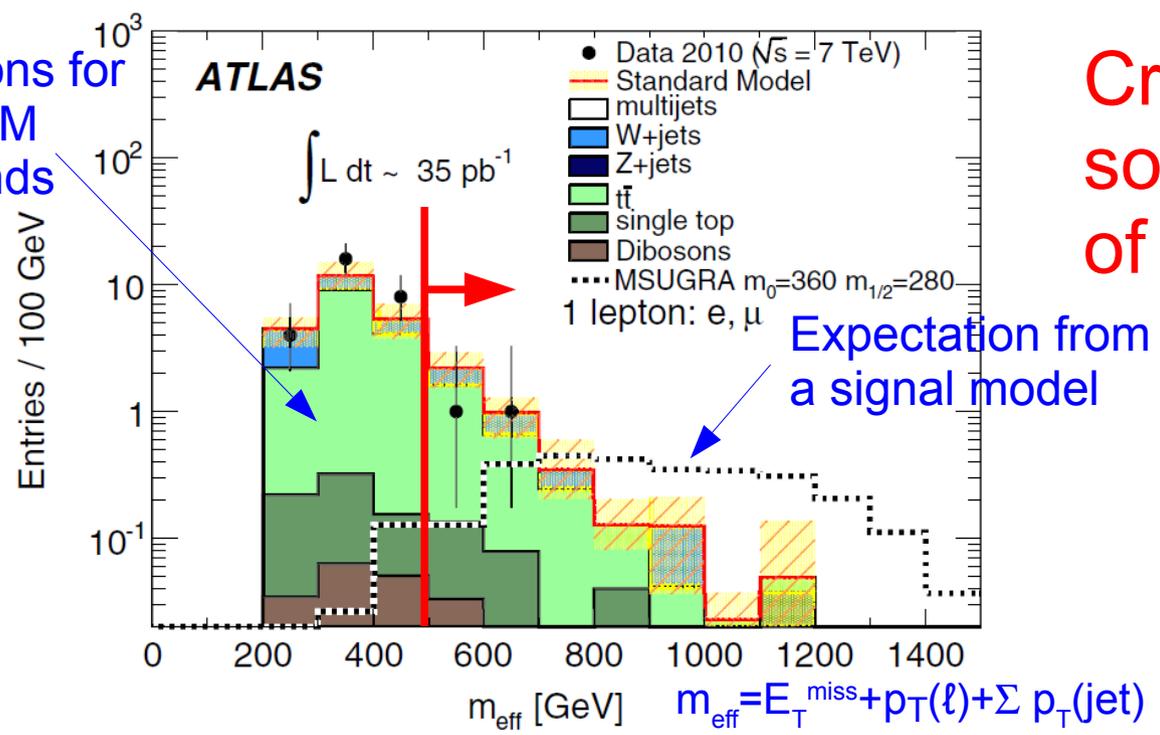
on SUSY: 48

SUSY Analysis and Backgrounds

Typical ATLAS supersymmetry search:

- Count events that has a characteristic SUSY signature and is unlikely to be due to Standard Model processes
 - More than ... jets and/or leptons
 - Measured energy above ... GeV
 - Missing transverse energy above ... GeV
- ...
- *Are there more events than expected from the Standard Model?*

Expectations for different SM backgrounds



Critical to have solid understanding of SM backgrounds

Estimating Backgrounds

Each analysis typically has multiple source of backgrounds
 Estimation of each depends on the nature of the background

False signal events due to
 e.g. fake E_T^{miss} or fake leptons

Standard Model
 Top, multijets
 V , VV , VVV , Higgs
 & combinations of these

SM events with same
 final state as the signal events

Combined fit of
 all regions and
 backgrounds and
 incl. systematic
 exp. and theor.
 uncertainties as
 nuisance
 parameters

Reducible backgrounds

Determined from data
 Backgrounds and methods
 depend on analyses

Irreducible backgrounds

*Dominant sources: normalise
 MC in data control regions*
Subdominant sources: MC

Validation

Validation regions used to
 cross check SM predictions
 with data

blinded

Signal regions

blinded

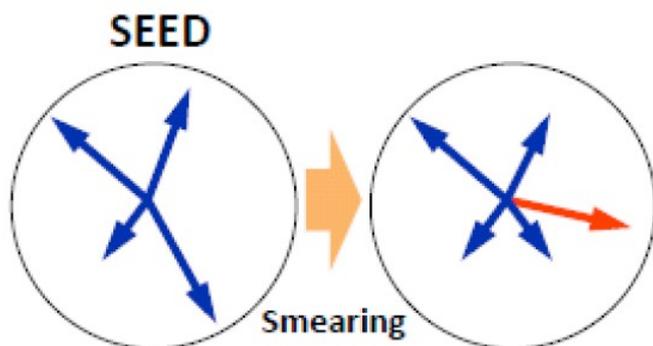
Reducible Background – Example

Jet mis-measurement can give large fake E_T^{miss}

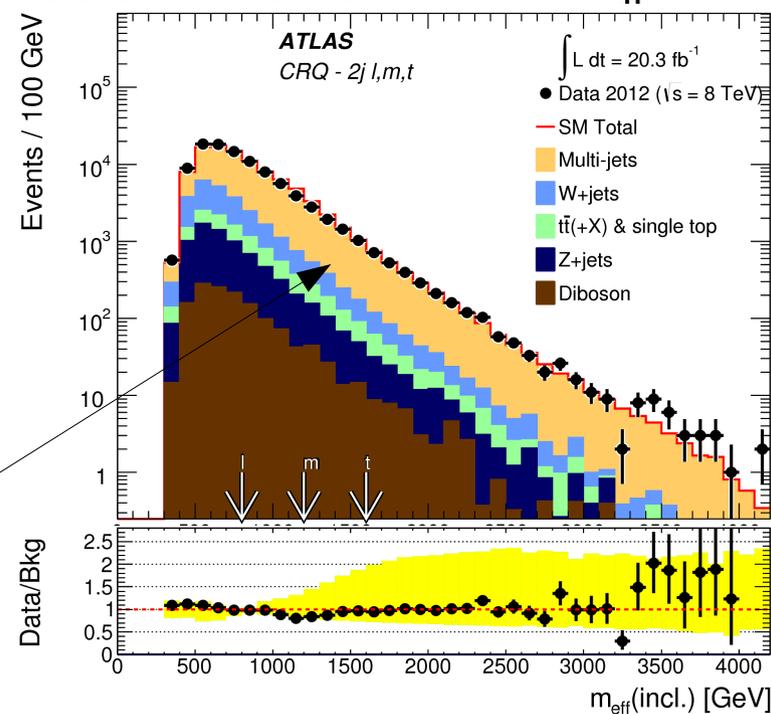
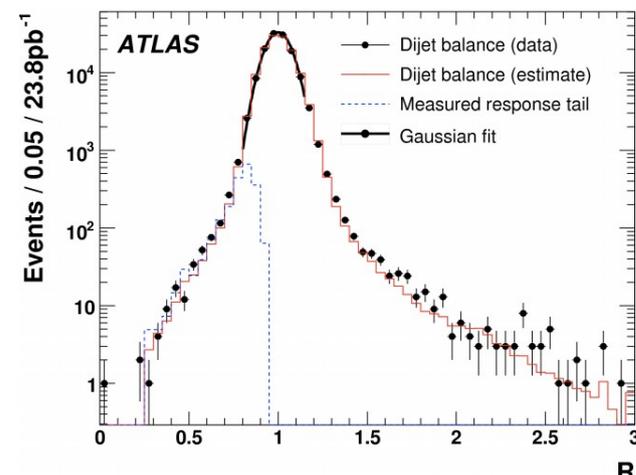
Makes E_T^{miss} -less multi-jet or $Z \rightarrow \ell\ell$ events look like signal

Jet-smearing method

- Jet-response function derived from MC and adjusted to match data in di-jet and three-jet events
- Select low E_T^{miss} events in data and smear the jets with a response function



- Provides good estimate of fake E_T^{miss} background as seen in control-region

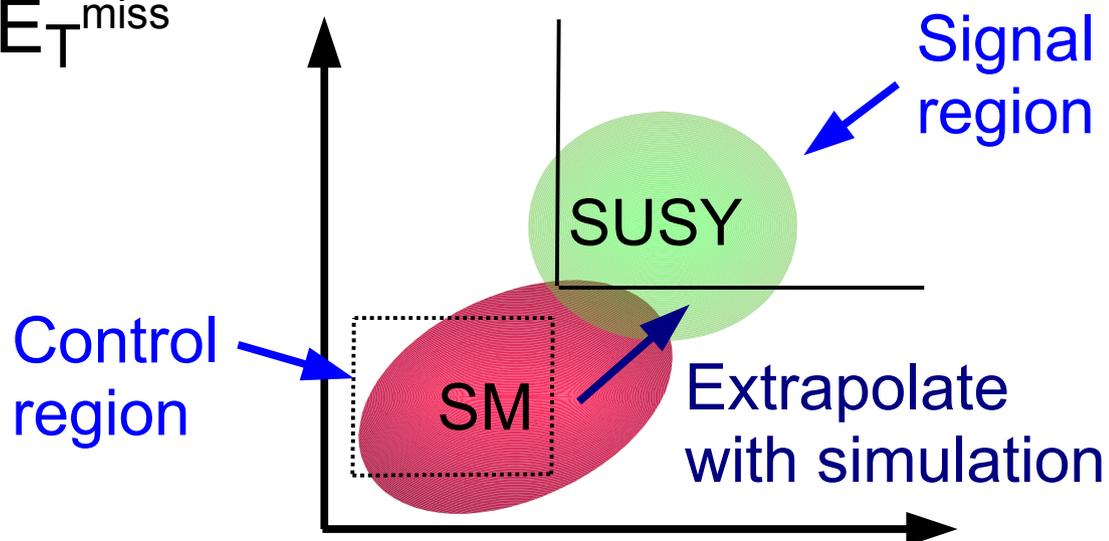


Irreducible Backgrounds

SM backgrounds with real E_T^{miss}
such as $Z \rightarrow \nu\nu$ and $W \rightarrow l\nu$

Rely (partly) on simulation

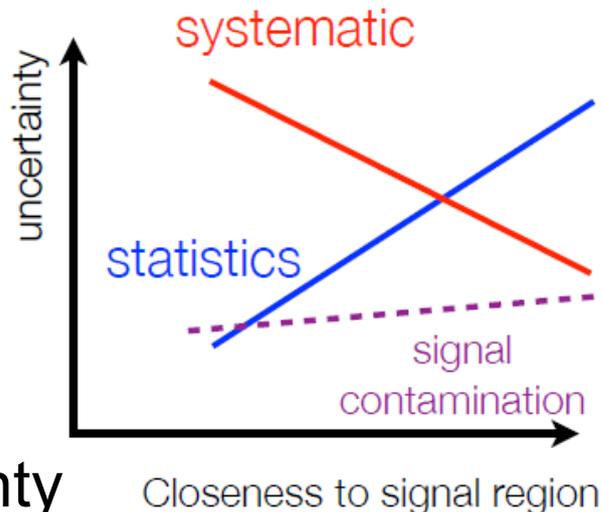
- Small contributions taken directly from simulation
- Larger one normalized to data in signal free region and extrapolated to signal region



$$N_{SR}^i = \frac{N_{SR}^{i,MC}}{N_{CR}^{i,MC}} (N_{CR}^{i,data} - \sum_{j=process} N_{CR}^{j,MC}) = T (N_{CR}^{i,data} - \sum_{j=process} N_{CR}^{j,MC})$$

Minimize uncertainty on background

- Systematic uncertainty on extrapolation factor from both theory (modelling) and experimental effects (efficiencies, etc.)
- In some cases can derive extrapolation factor from data or at least correct using data to reduce systematic uncertainty



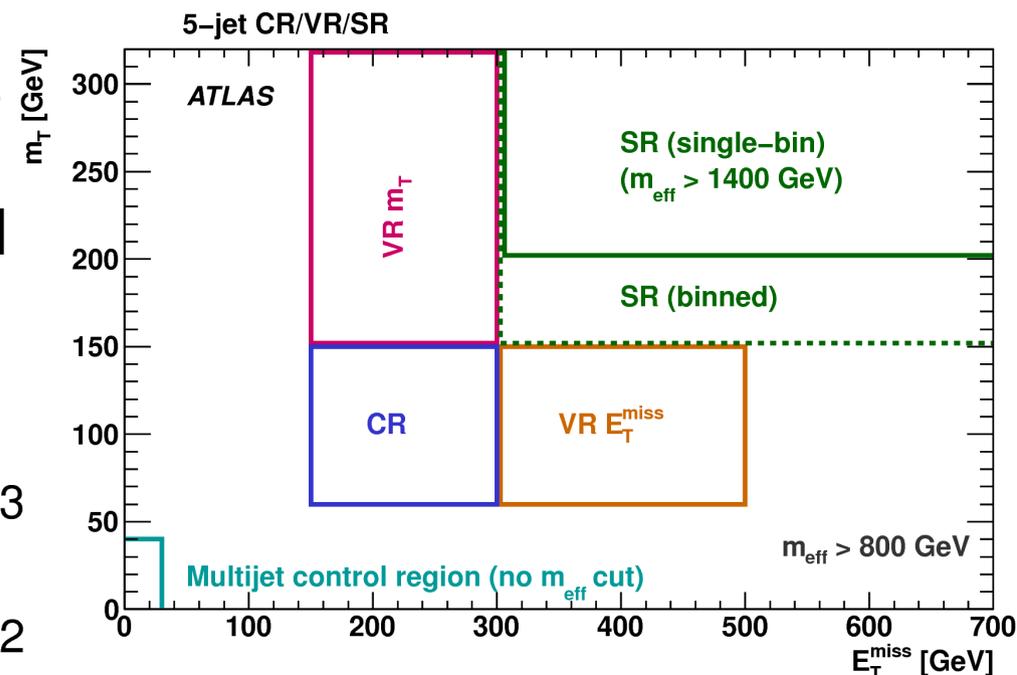
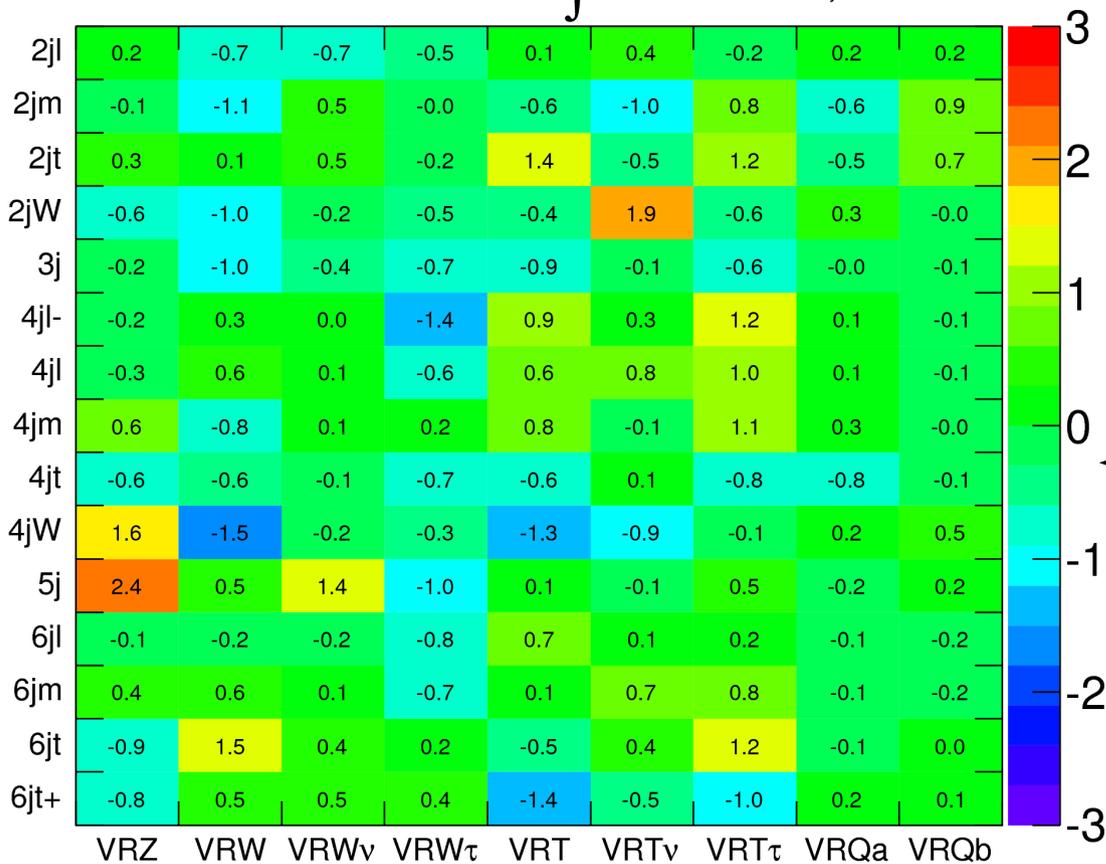
Background Validation

Validate background predictions with dedicated validation regions

- Extrapolate to regions close to signal region, but signal depleted
- If prediction matches the data have confidence in extrapolation

ATLAS

$$\int L dt = 20.3 \text{ fb}^{-1}, \sqrt{s}=8\text{TeV}$$



Can get complicated

- 9 validation regions for each of the 15 signal regions in “0-lepton” search
- Deviation from prediction shown in “ σ ”

LHC Supersymmetry Search Results

Short Version

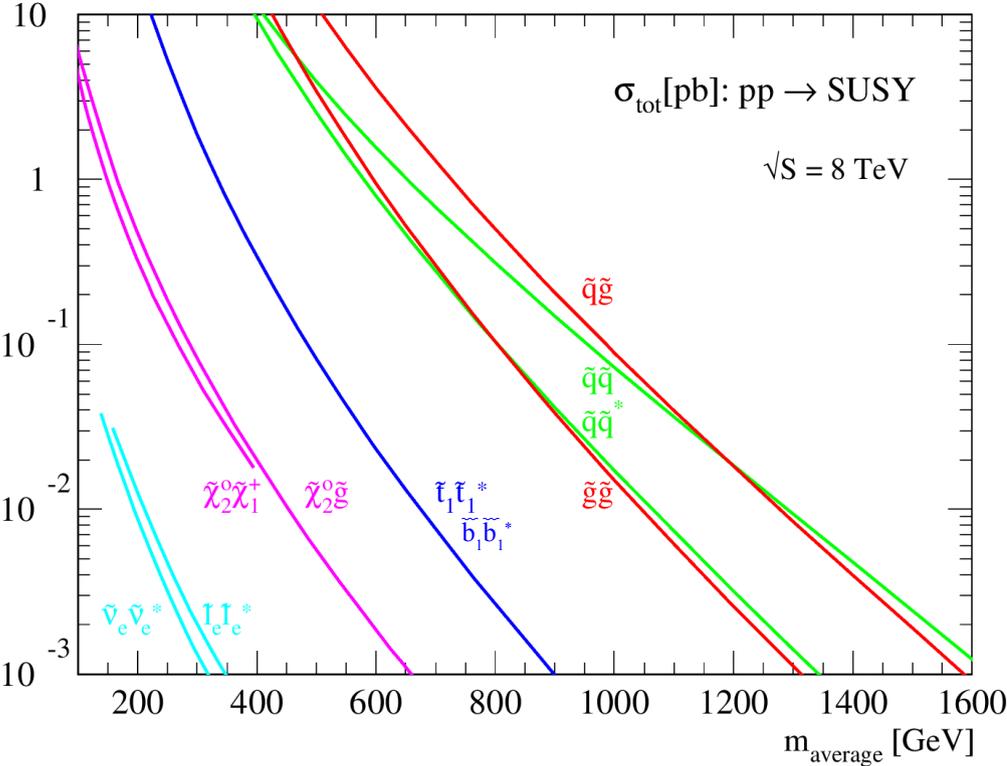


Did We Really Look Everywhere?



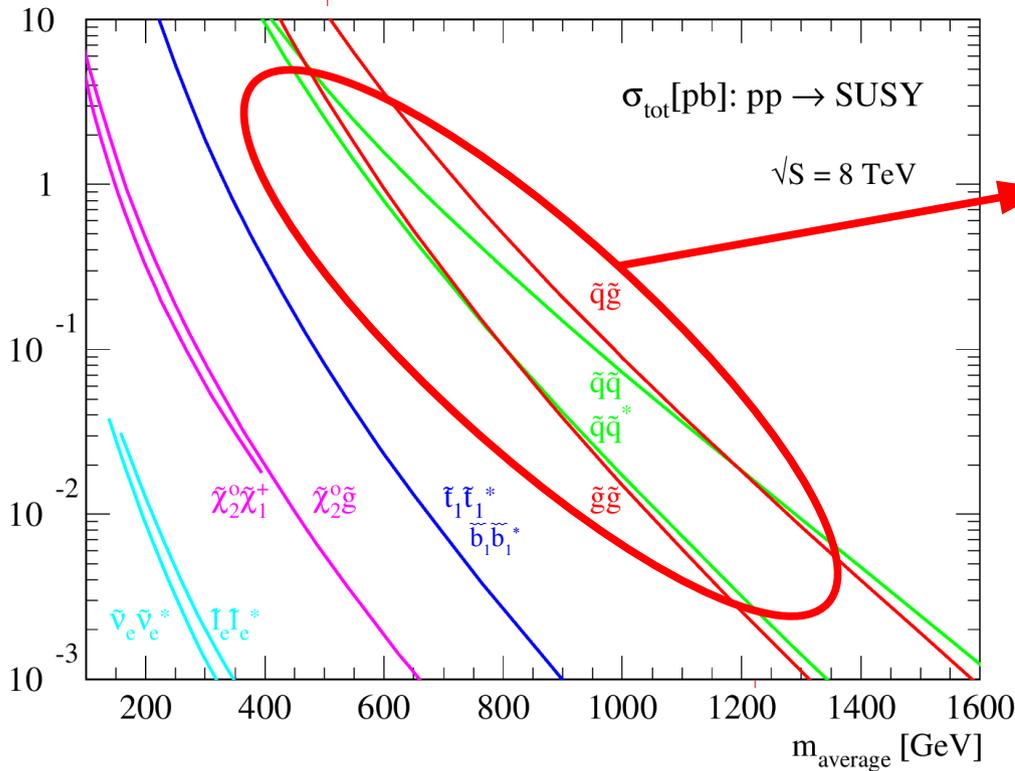
LHC SUSY Searches

At LHC can search for production for almost all SUSY particles, but with different sensitivity as production cross-sections vary



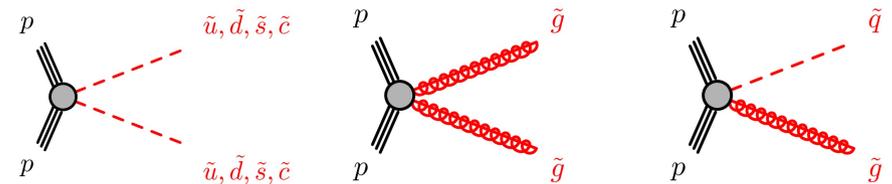
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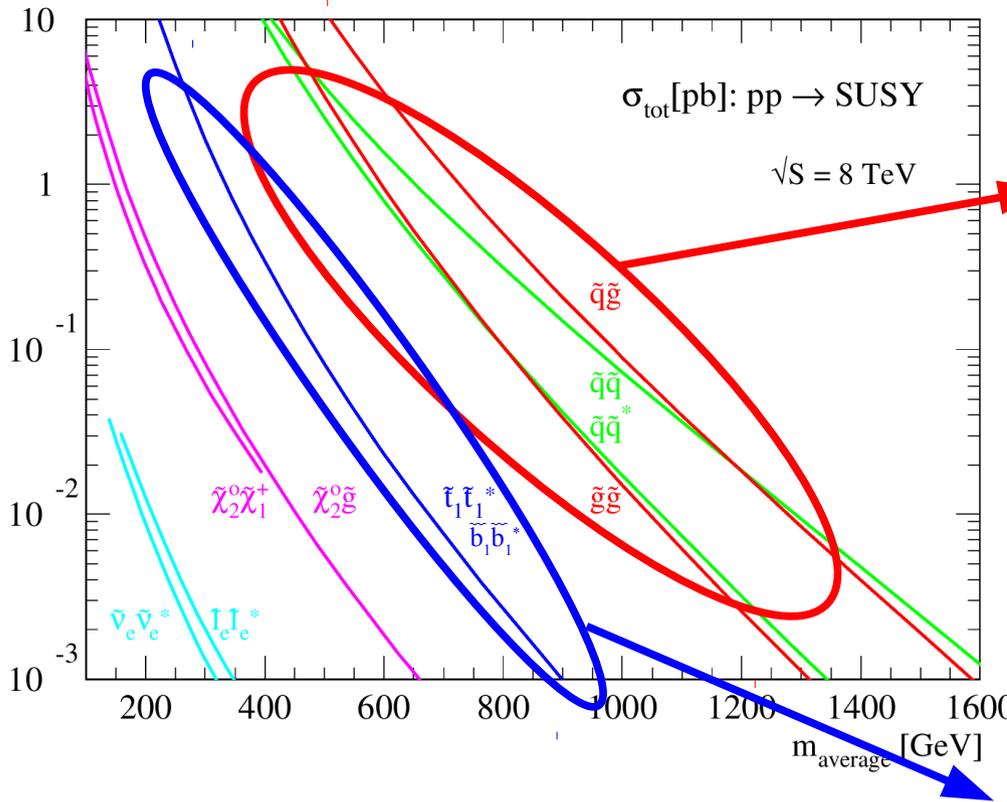
Production of squarks/gluinos:

- Strong coupling gives large sensitivity at the LHC
- Primary signature is jets and missing energy



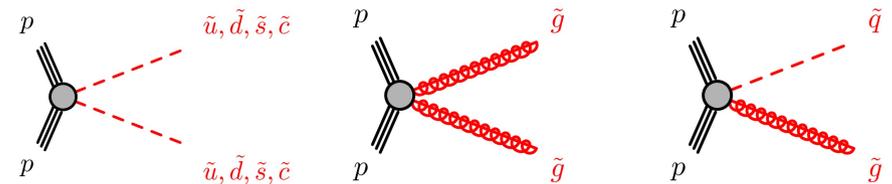
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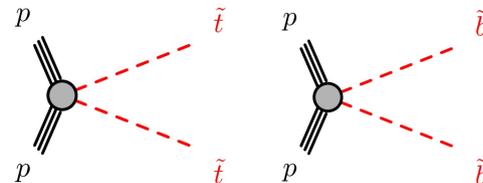
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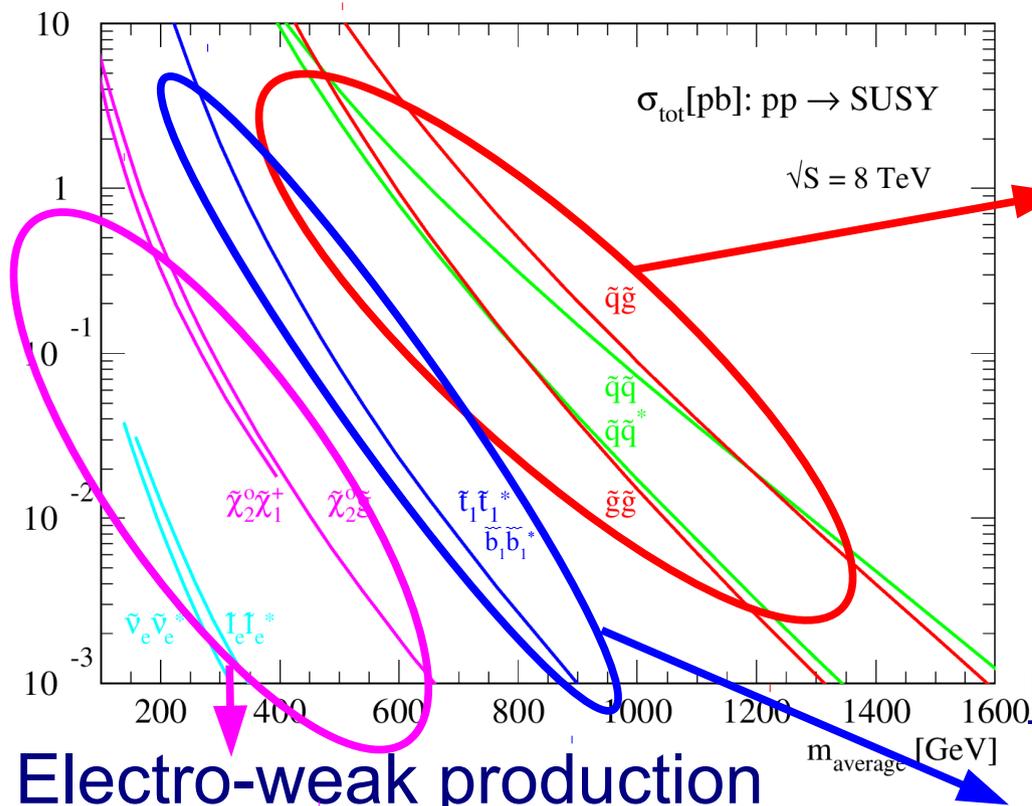
Prod. of 3rd generation squarks:

- Lower cross section,
- Decays chains frequently has top and bottom quarks
- Of great interest to “Natural SUSY”



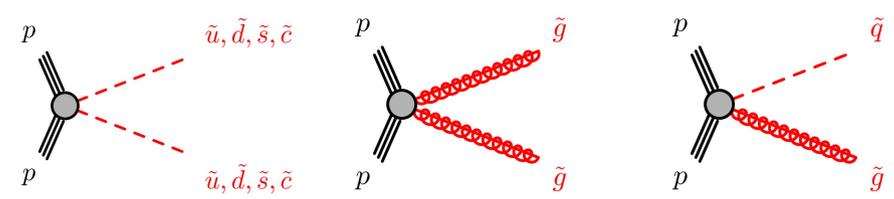
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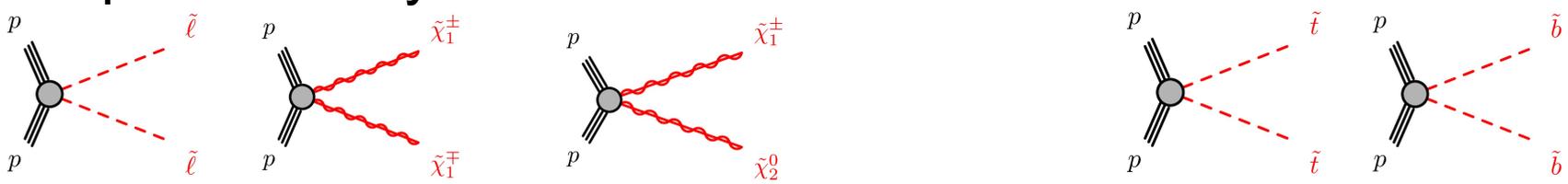


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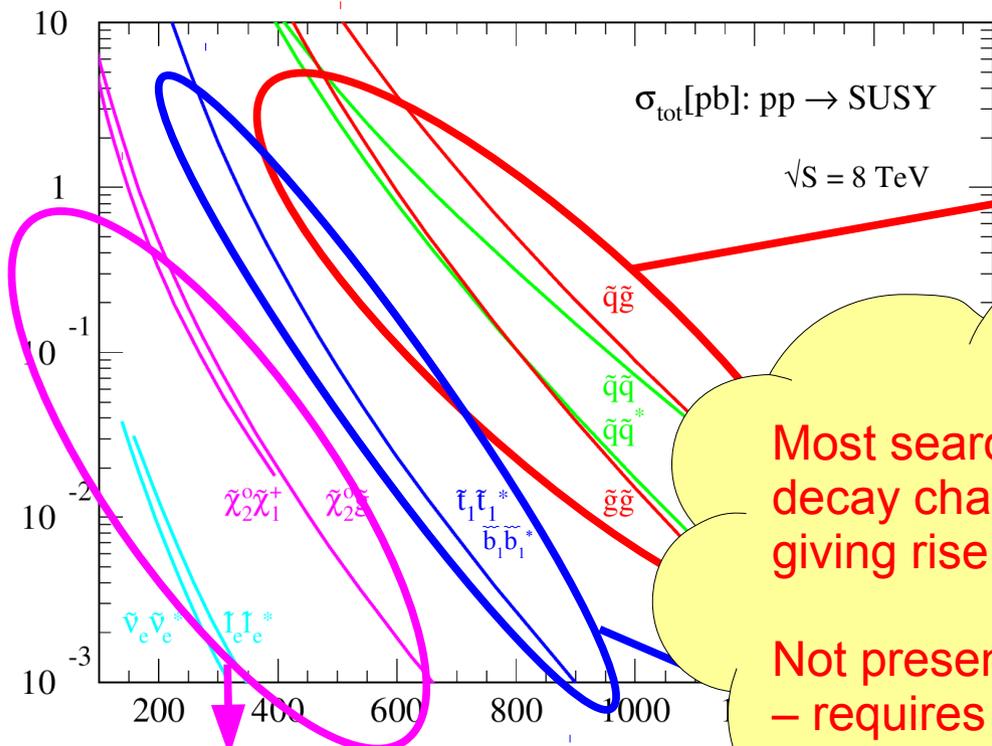
Electro-weak production

- Slepton, neutralino and charginos
- Detectable mostly through leptonic decays



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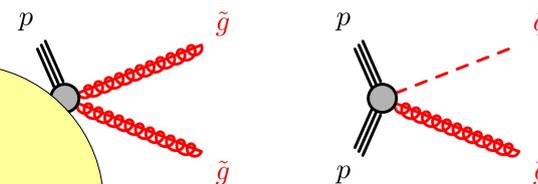


Production of squarks/gluinos:

- Strong coupling gives large sensitivity at the LHC
- Primary signature is jets and missing energy

Most searches assume decay chain to stable LSP giving rise to missing energy

Not present in all SUSY models – requires dedicated searches to cover this



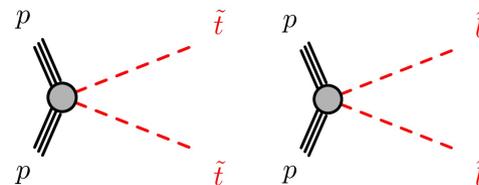
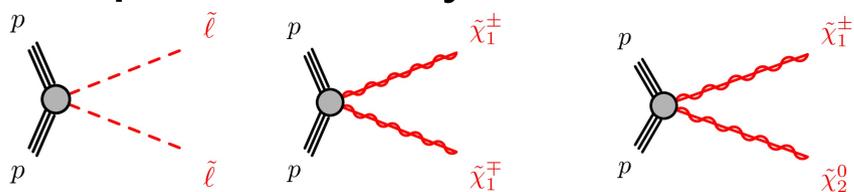
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Production squarks:

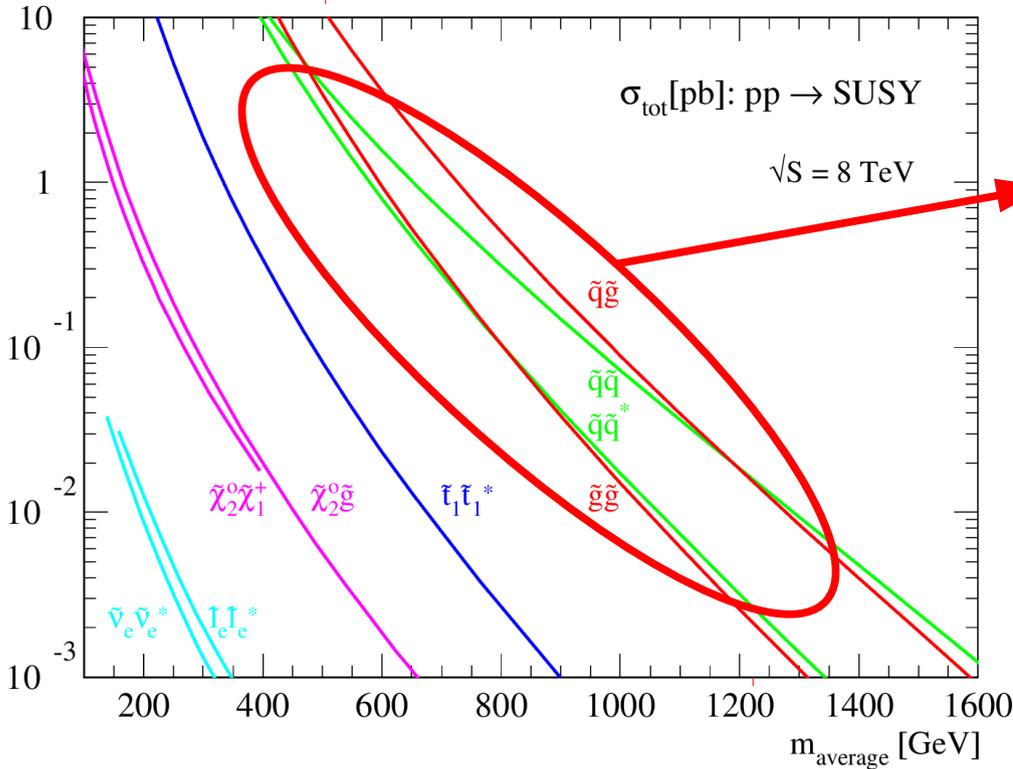
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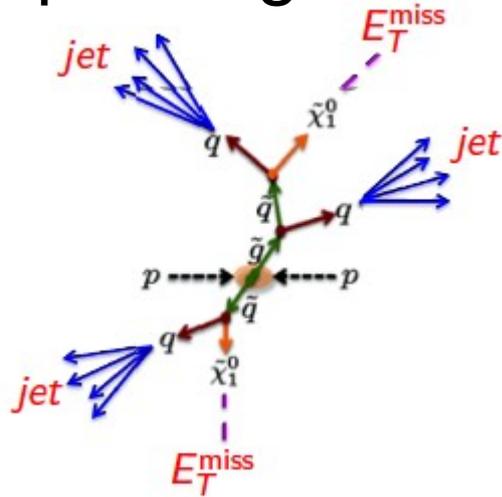


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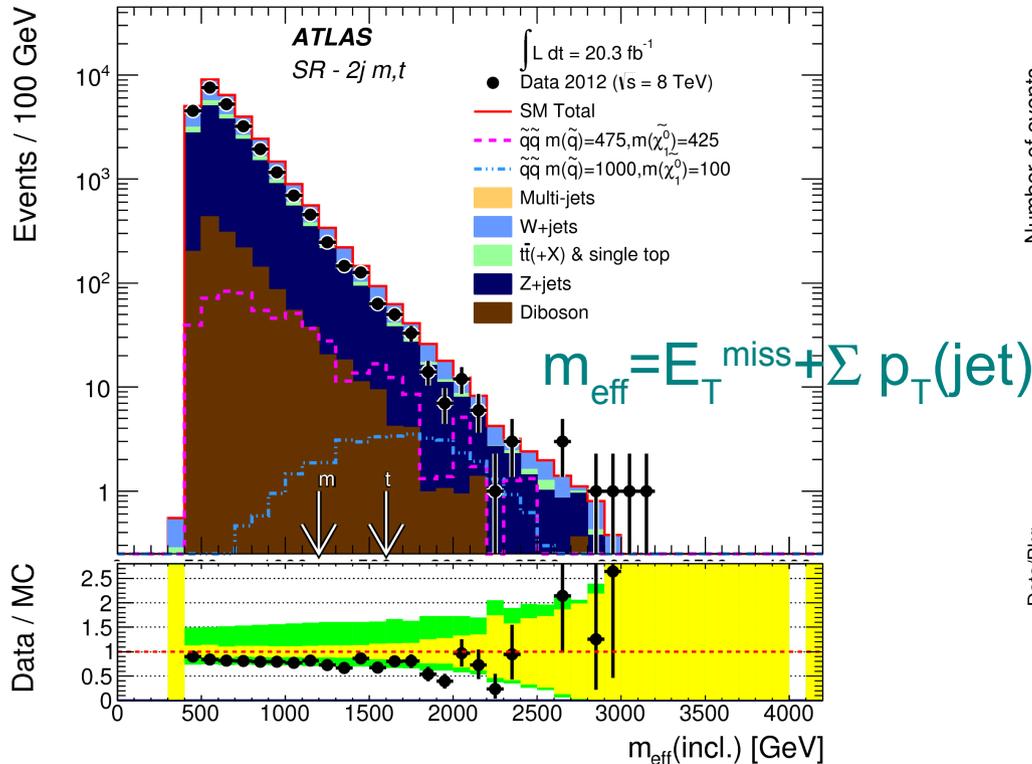
Inclusive 0-Lepton Analysis

Work-horse for searches for squarks/gluinos

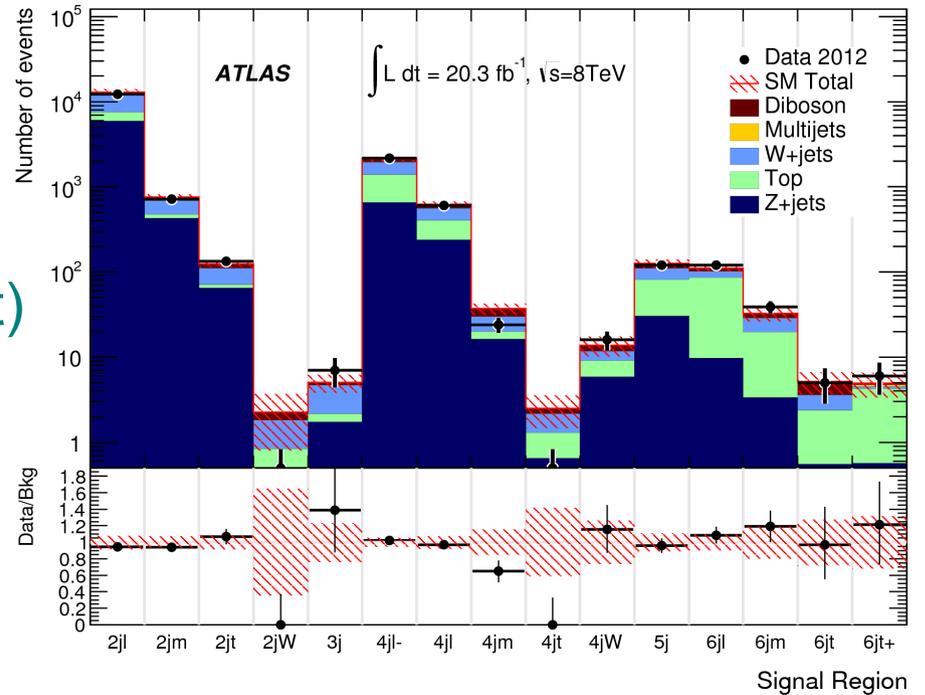


Inclusive selection

- Between 2-6 jets with $p_T > 60$ GeV
- $\Delta\phi(j_i, E_T^{\text{miss}}) > 0.4$ (0.2 for $i > 3$)
- No leptons with $p_T > 10$ GeV
- Large E_T^{miss} wrt m_{eff} ($> 0.15 - 0.4$)
- Signal region split by #jets and m_{eff} with two SRs dedicated to $W \rightarrow$ jets

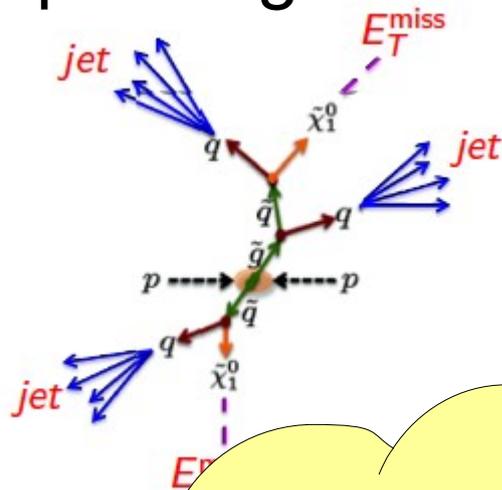


Background and data yields in 15 SRs



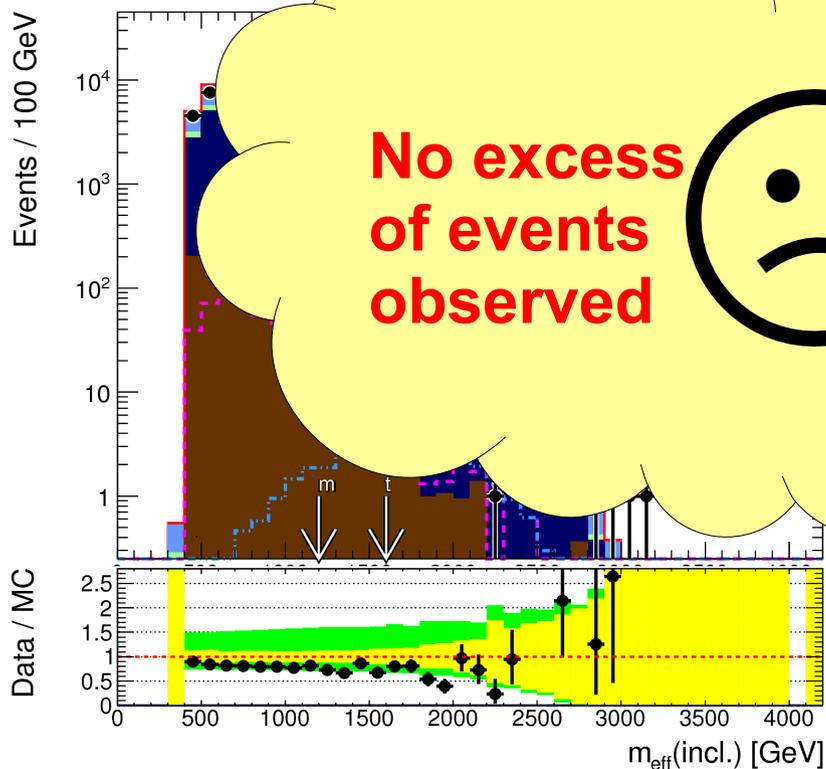
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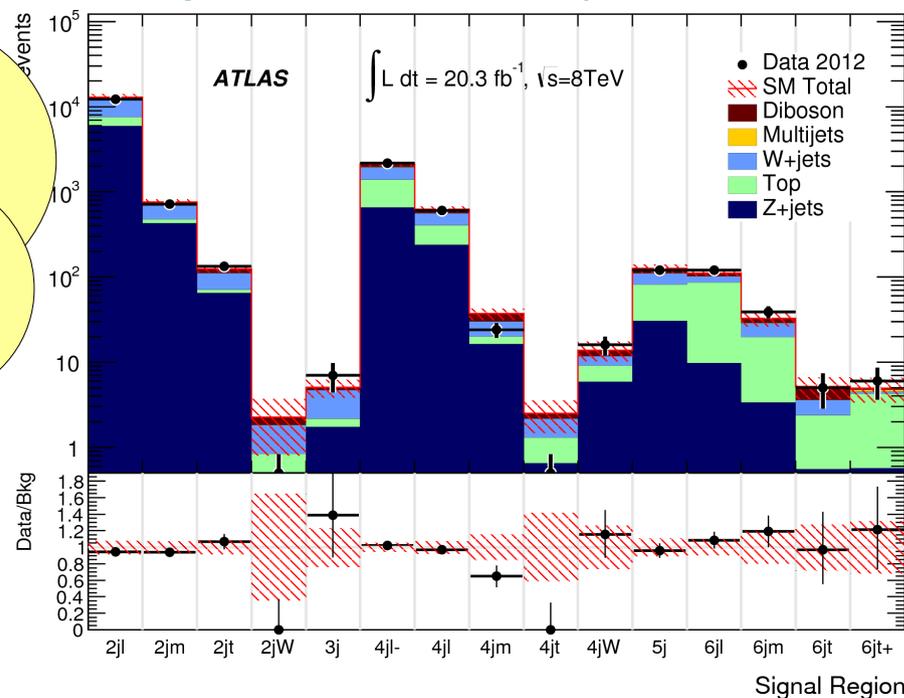
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No excess of events observed



Background and data yields in 15 SRs



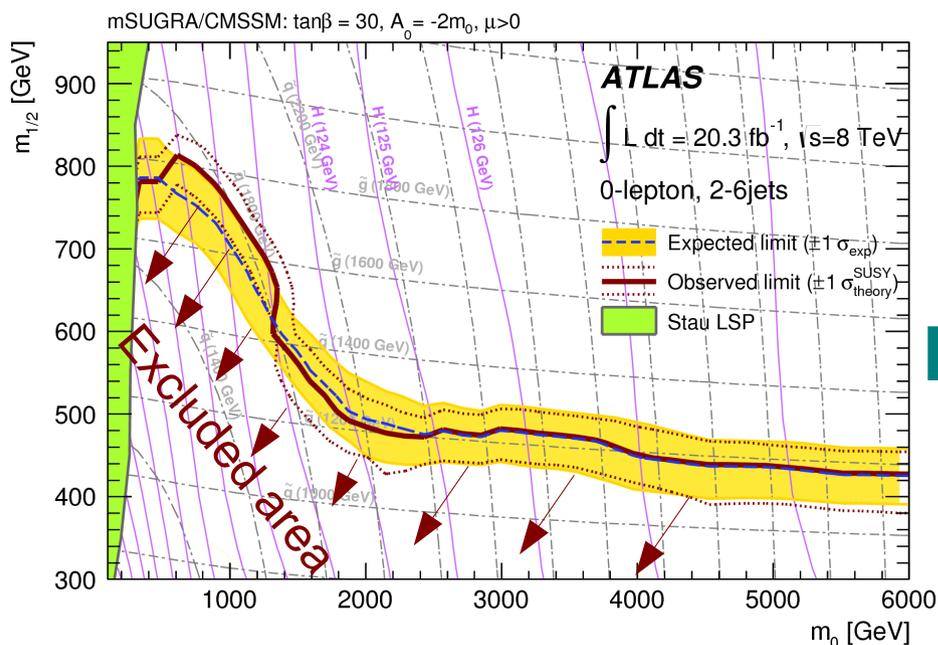
Inclusive 0-Lepton Analysis Results

With no signals observed, proceed to set limits on SUSY models

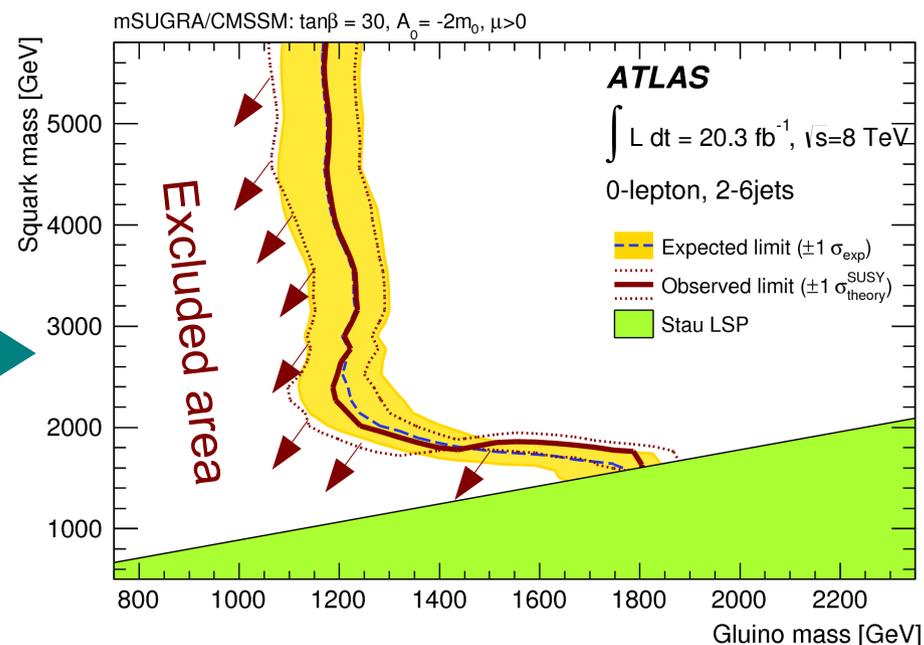
Limit setting procedure

- For each SUSY model calculate expected number of events in each signal region and any leakage to control regions
- Obtain p_0 from simultaneous fit to the signal and control regions
- Limit quoted from best expected signal region
- Typically do this in scan over several SUSY model parameters

Limits on mSUGRA parameters



Limits on squark/gluino masses



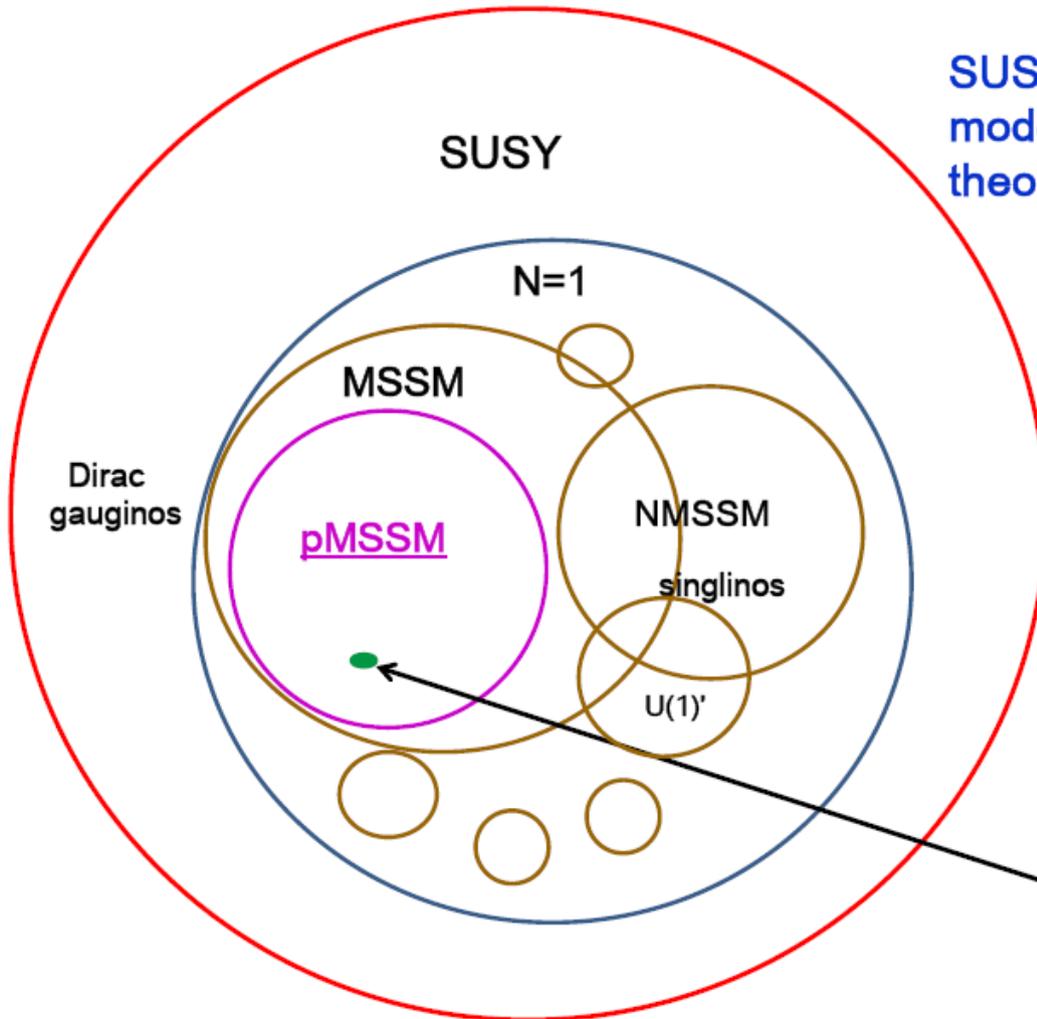
Inclusive 0-Lepton Analysis Results

With no signals observed, proceed to set limits on SUSY models

Limit setting procedure

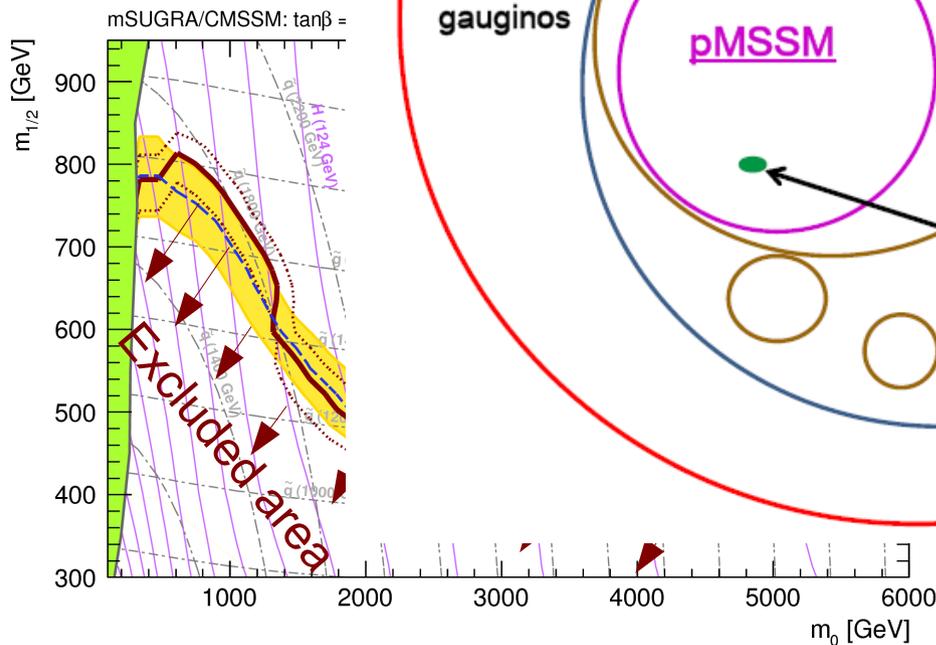
- For each model
- of every parameter
- Obtain the cross-section
- Limit $q_{\text{obs}} < q_{\text{th}}$
- Typical $q_{\text{obs}} \approx 100 \text{ fb}^{-1}$

SUSY is *not* a single model but a very large theoretical framework

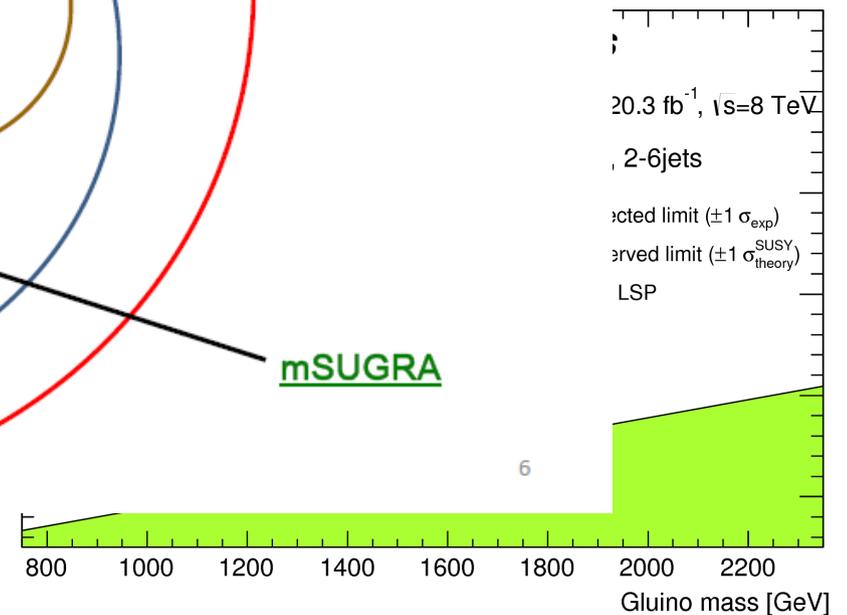


parameters

Limits on $m_{1/2}$



processes

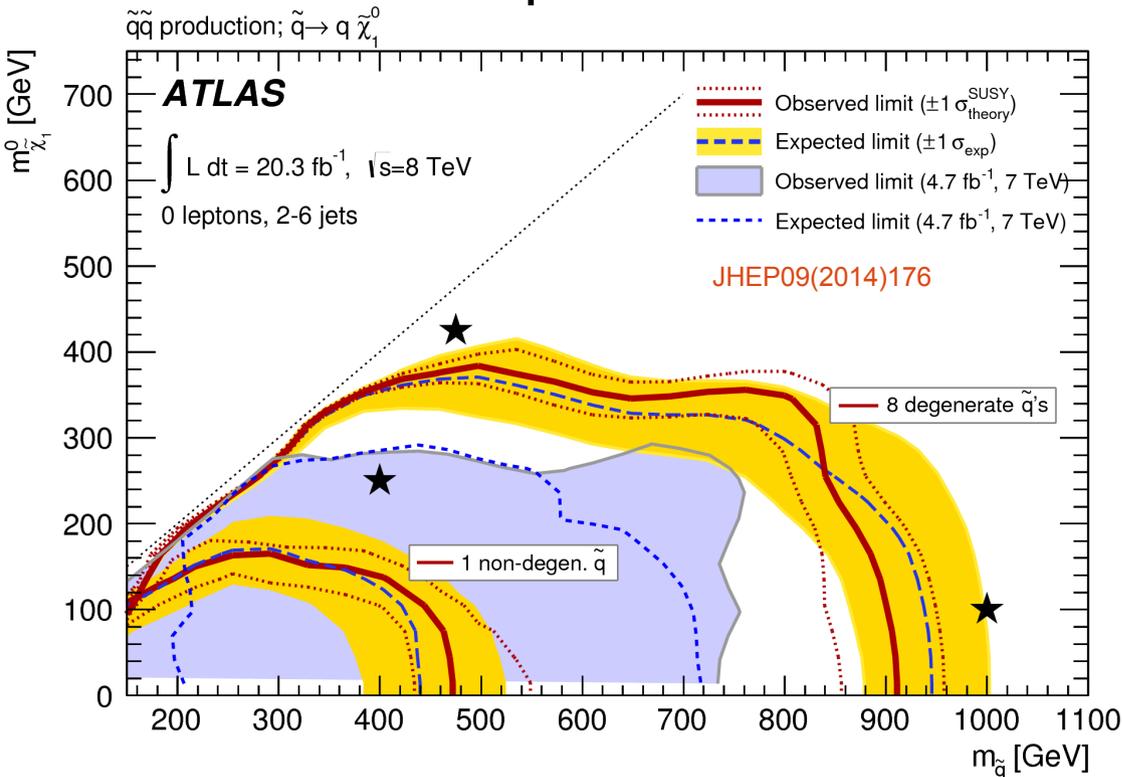
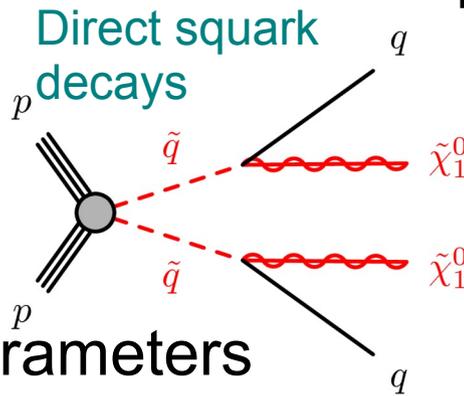


Simplified Model Limits

Present results in form of simplified models

Simplified Models

- 1 production channel
- 1 decay channel (BF=100%)
- Minimal amount of parameters
- Assume SUSY production σ , but can reinterpret broader

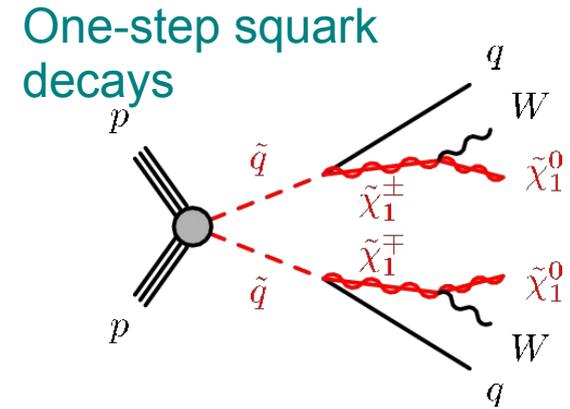
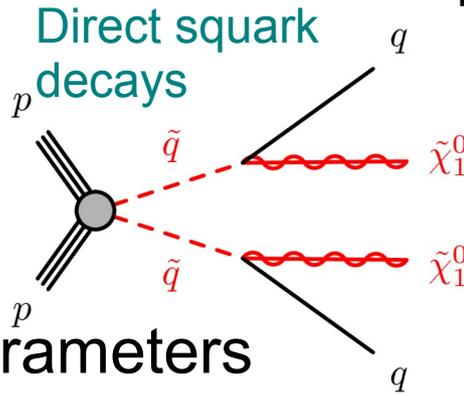


Simplified Model Limits

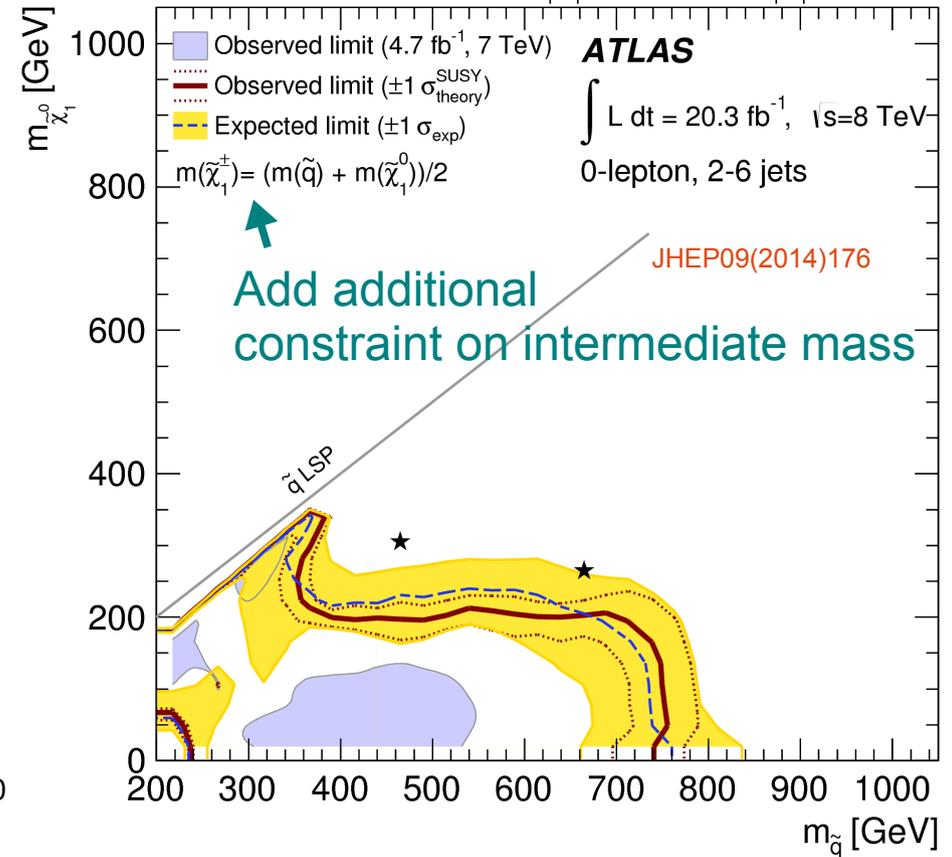
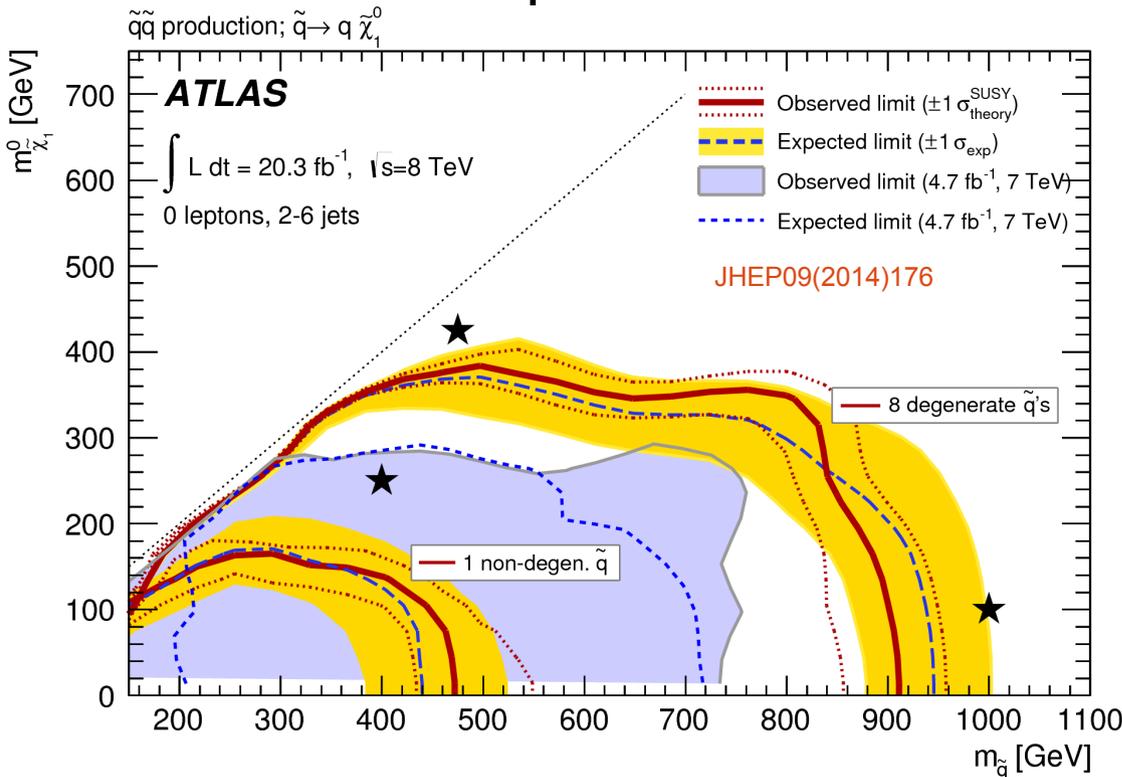
Present results in form of simplified models

Simplified Models

- 1 production channel
- 1 decay channel (BF=100%)
- Minimal amount of parameters
- Assume SUSY production σ , but can reinterpret broader



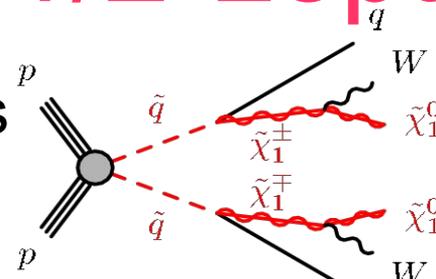
Simplified model, $\tilde{q}\tilde{q}^* \rightarrow q\bar{q} \tilde{\chi}_1^+ \tilde{\chi}_1^- \rightarrow q\bar{q} W^+ W^- \tilde{\chi}_1^0 \tilde{\chi}_1^0$



New

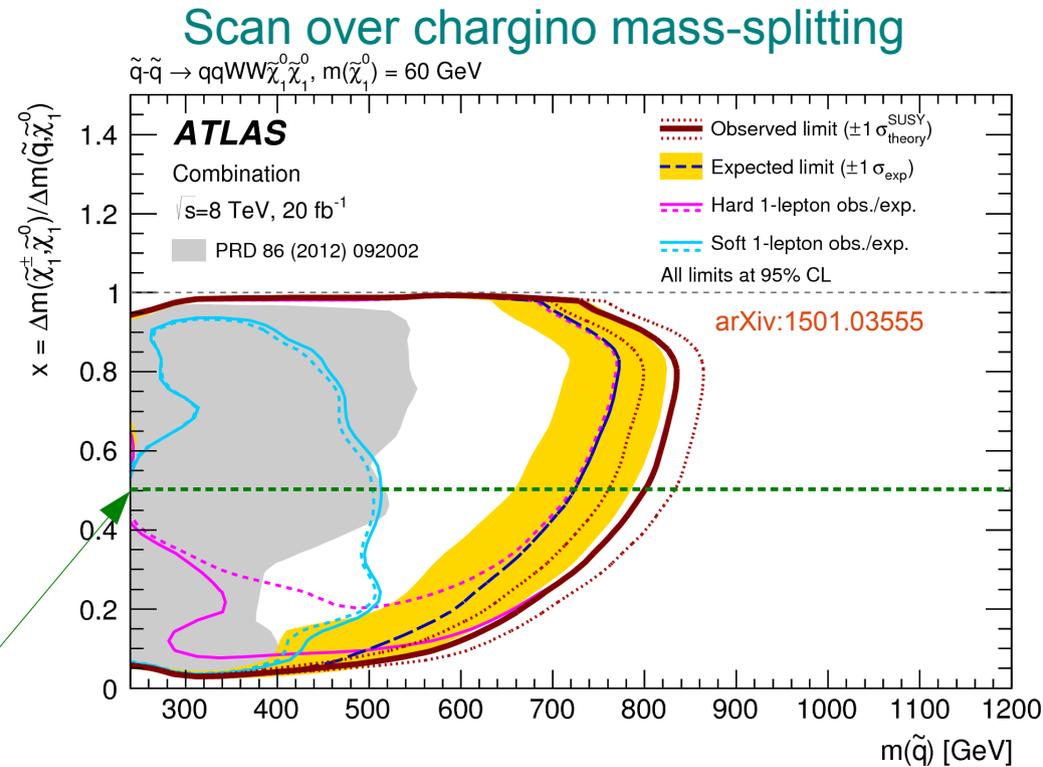
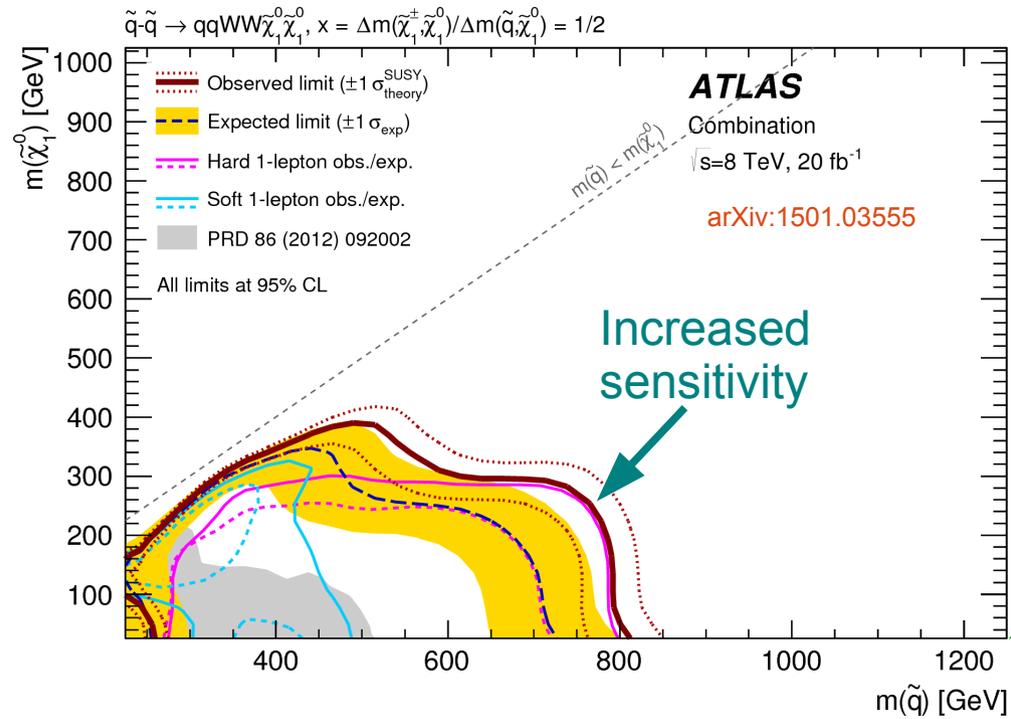
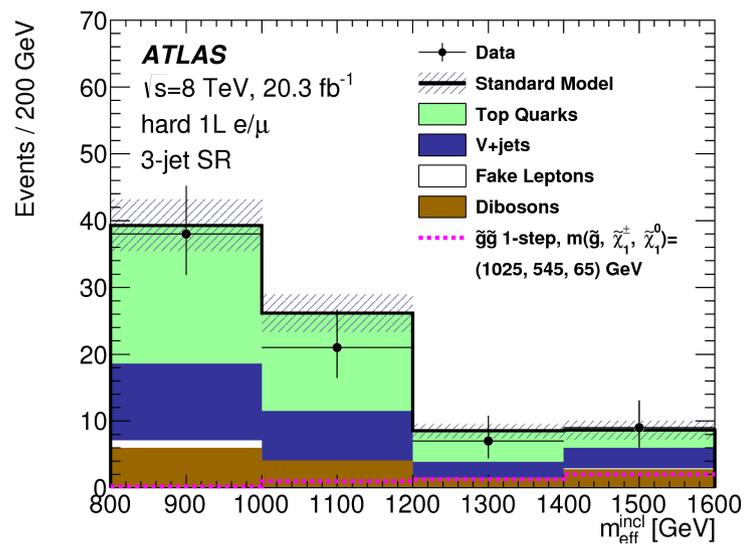
Inclusive 1/2-Lepton Analysis

Longer decay chains gives more possible signatures



1/2-lepton analysis

- 1 or 2 leptons with $p_T > 6$ GeV, split in hard ($p_T > 25$ GeV) and soft ($6 < p_T < 25$ GeV)
- 3, 5 and 6 jet signal regions (also 2 jets for 2-lepton)
- Large transverse mass, $m_T = \sqrt{2p_T^\ell E_T^{\text{miss}}(1 - \cos[\Delta\phi(\vec{\ell}, \vec{p}_T^{\text{miss}})])}$ to suppress W and top bkg
- Simultaneous fit to multiple bins in m_{eff} or $E_T^{\text{miss}}/m_{\text{eff}}$ to increase sensitivity

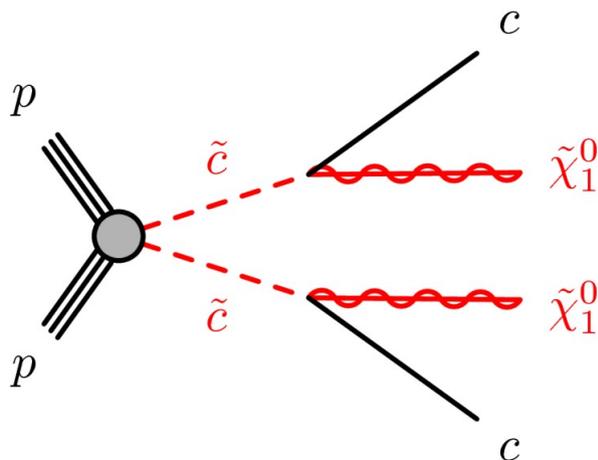
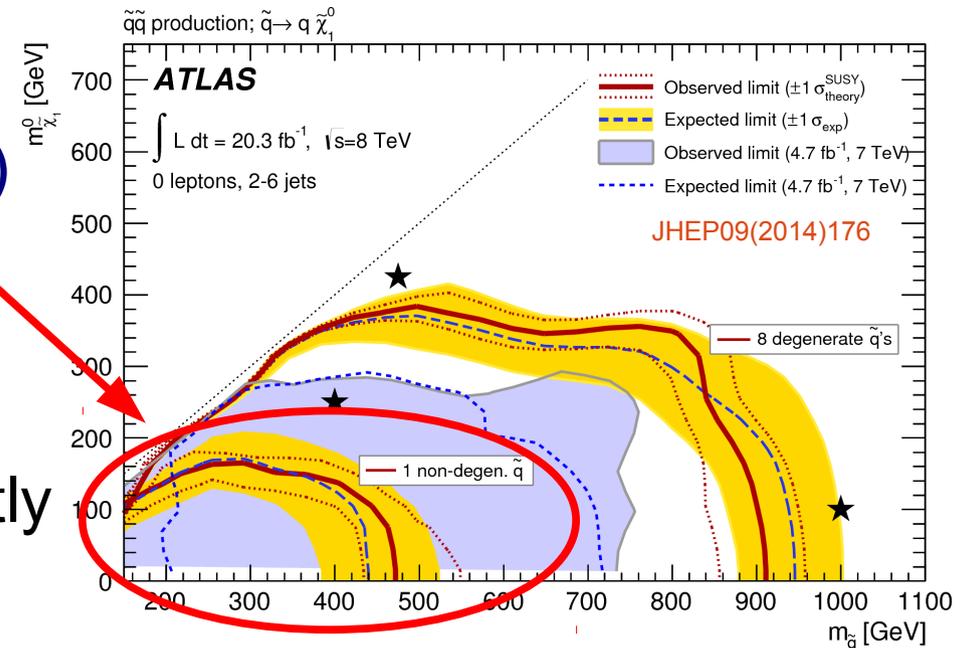


New

Scalar Charm Search

Much weaker squark limit if only one light light-flavor squark ($1/8\sigma$)

- Cannot have any light-flavor squark on its own as it would violate flavor-physics constraints
- However, scharm is not significantly constrained from flavor physics and could be lighter than others



New, dedicated scharm search

- Search for direct decay to charm
- Signal is two high- p_T charm jets, high E_T^{miss} and no leptons
- Use dedicated charm-tagger for jets to suppress other jet types

New

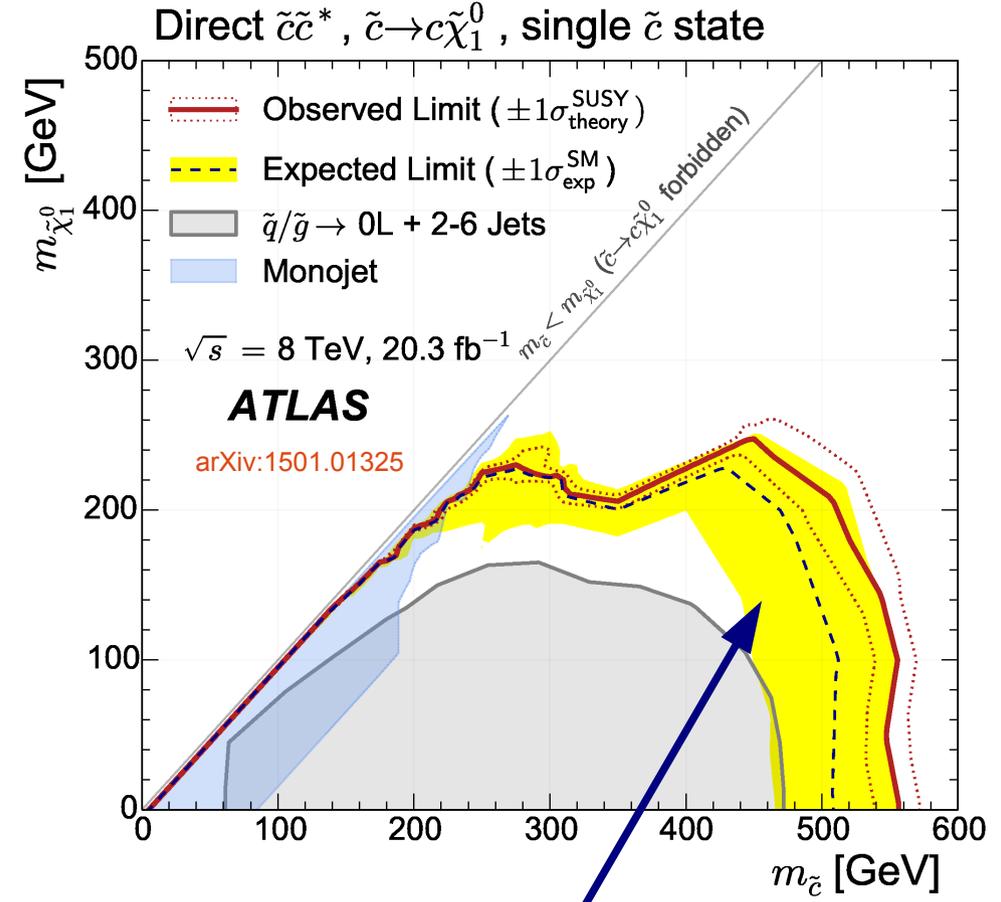
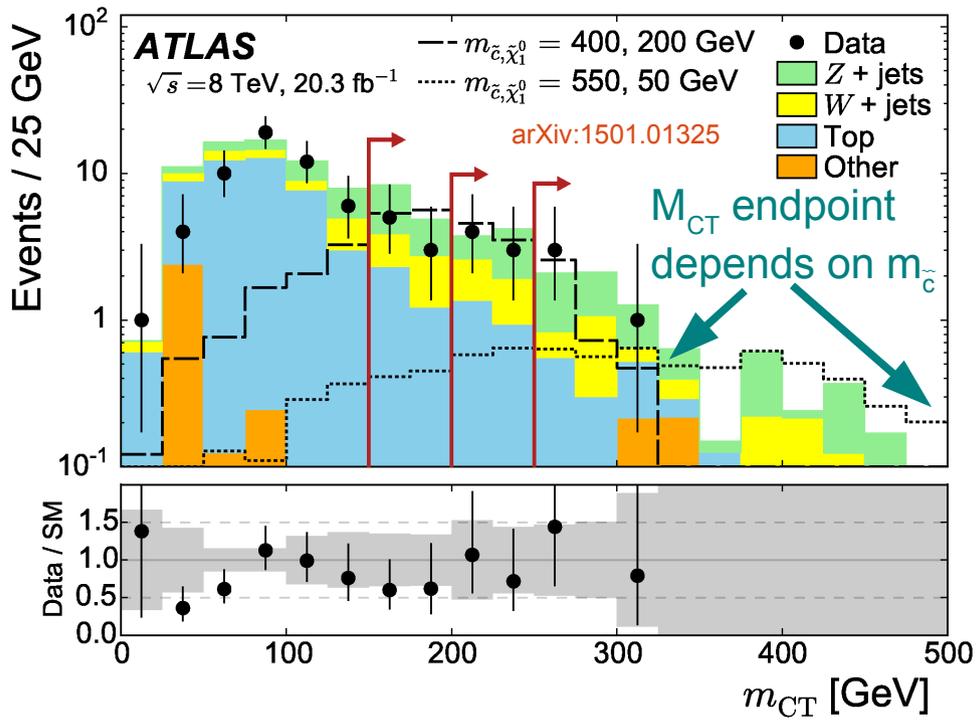
Scalar Charm Search

- Kinematic selection using contransverse mass m_{CT} and charm pair mass m_{CC}

$$m_{CT}^2(j_1, j_2) = [E_{T,1} + E_{T,2}]^2 - [\mathbf{p}_{T,1} - \mathbf{p}_{T,2}]^2$$

$$E_T = \sqrt{p_T^2 + m^2}$$

$$m_{CT}^{max} = \frac{m^2(\tilde{c}) - m^2(\tilde{\chi}_1^0)}{m(\tilde{c})}$$

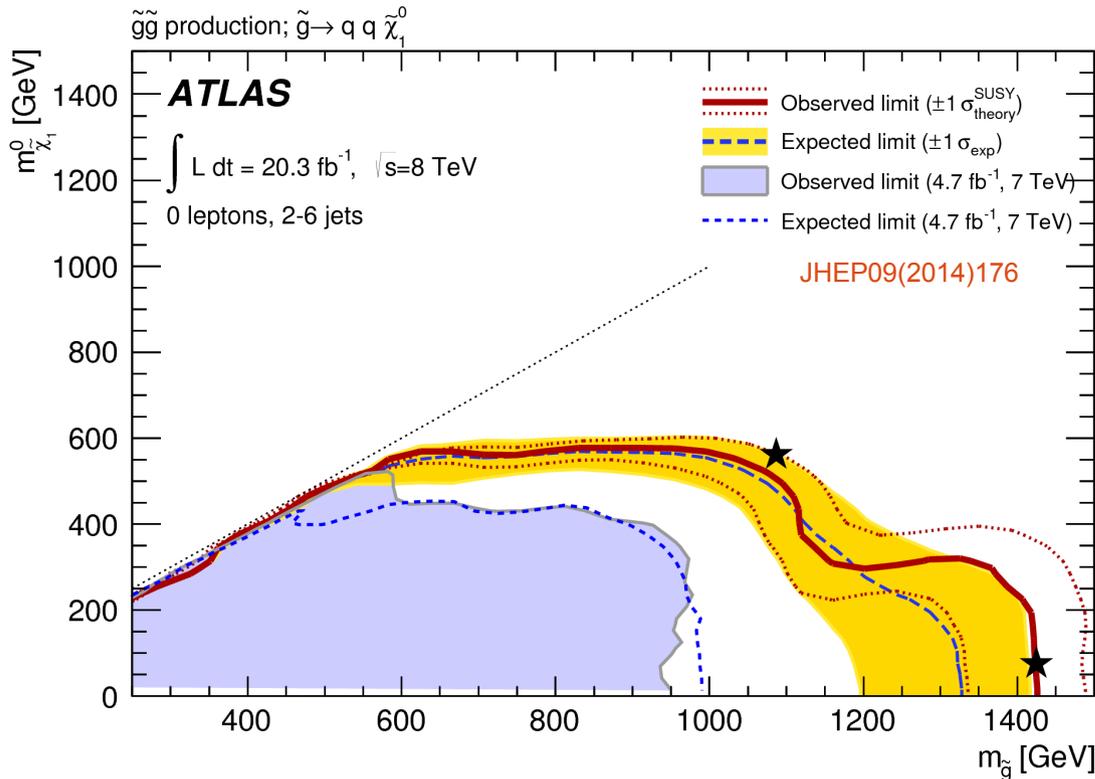
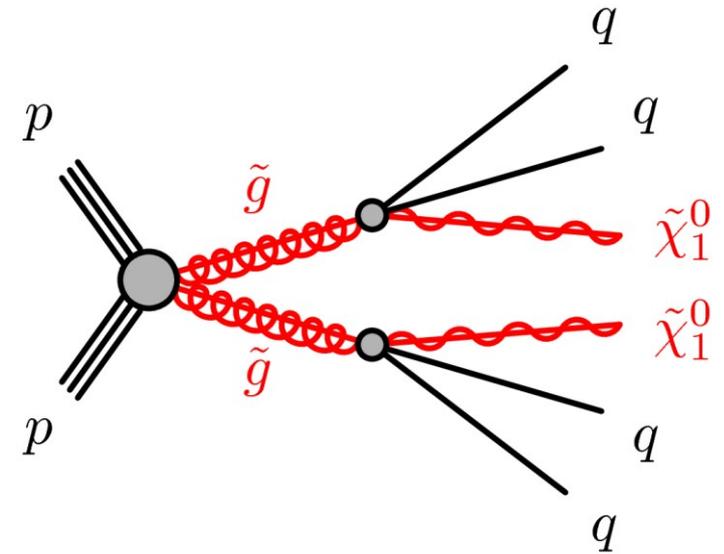


~100 GeV improvement in sensitivity w.r.t. inclusive search

Limits on Gluino Production

Even stronger limits on gluinos

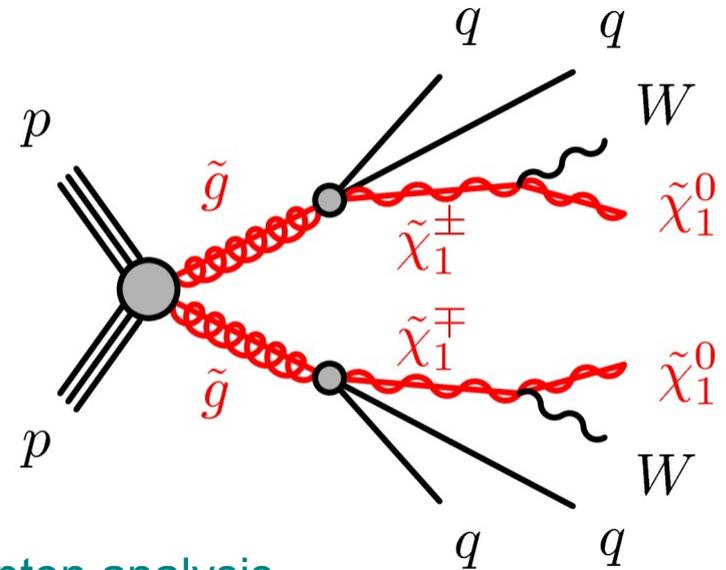
- Several quarks(=jets) in each decay
- Inclusive 0-lepton analysis sensitive up to 1.3 TeV for direct gluino decays



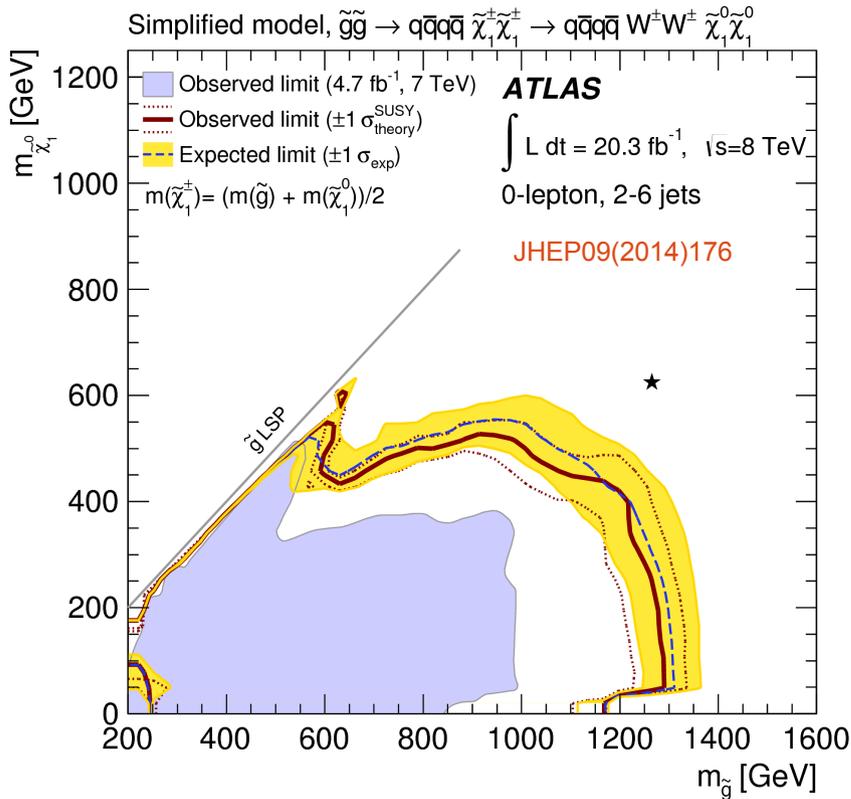
Limits on Gluino Production

Even stronger limits on gluinos

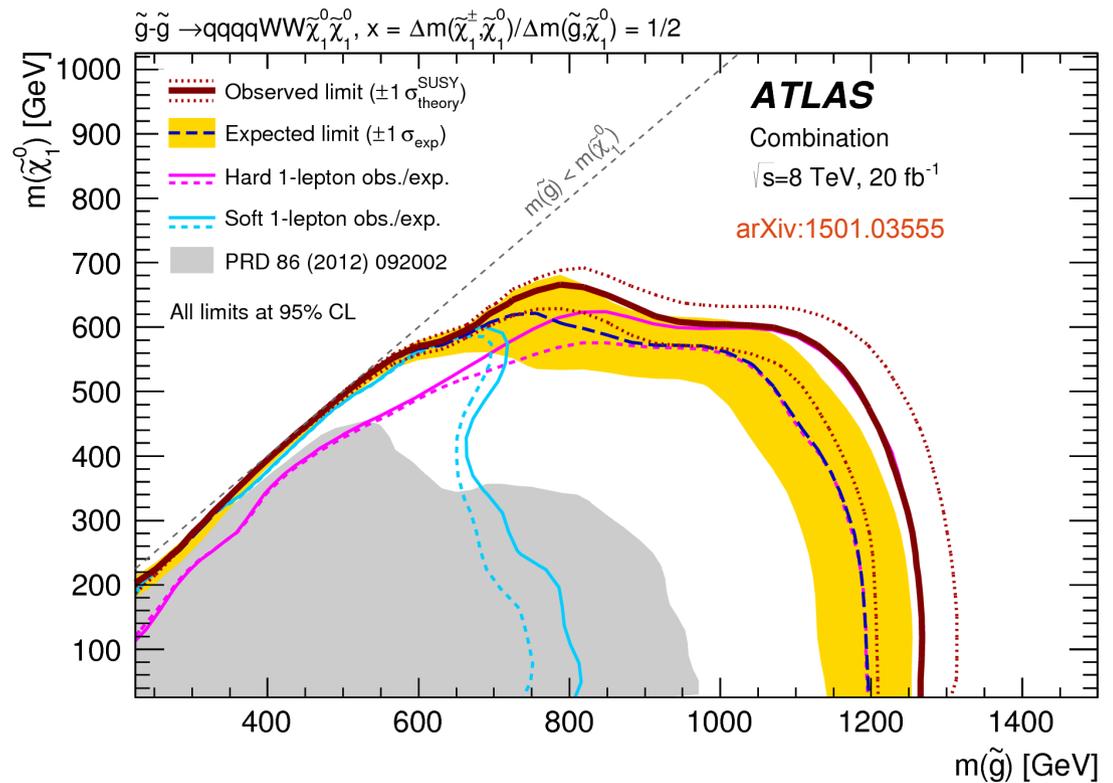
- Several quarks(=jets) in each decay
- Inclusive 0-lepton analysis sensitive up to 1.3 TeV for direct gluino decays
- For one-step decay, good complementarity with 1-lepton analysis



0-lepton analysis



1-lepton analysis

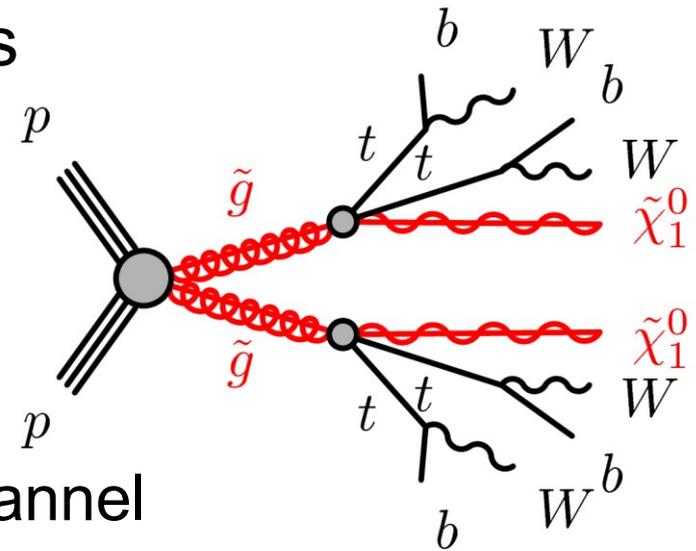


Limits on Gluino Production

Many more decays possible for gluinos

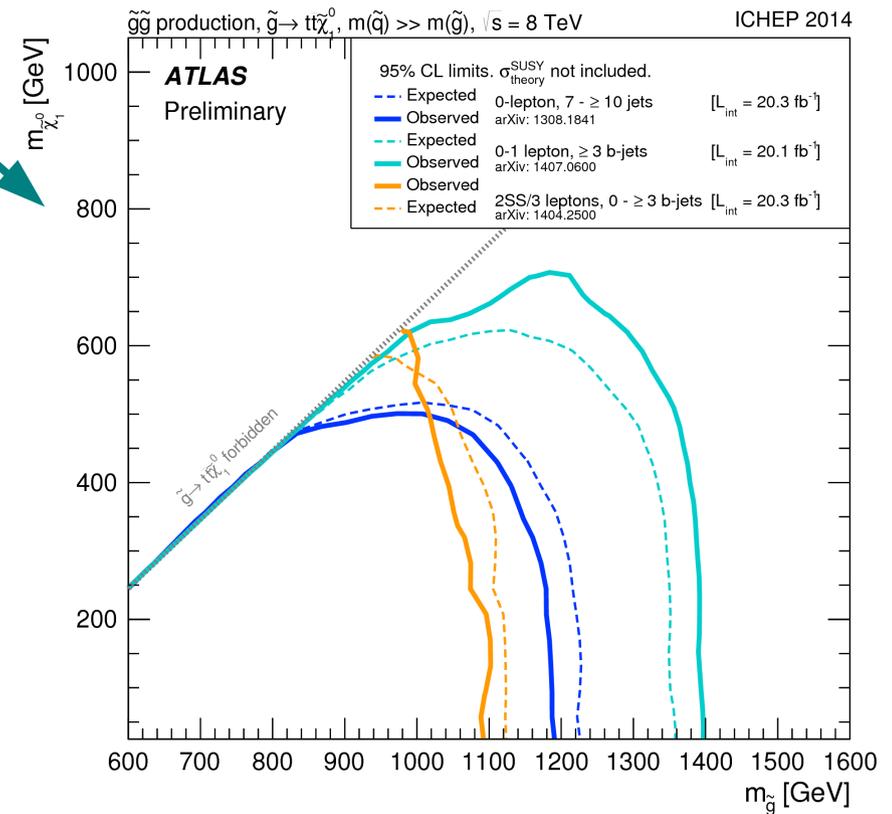
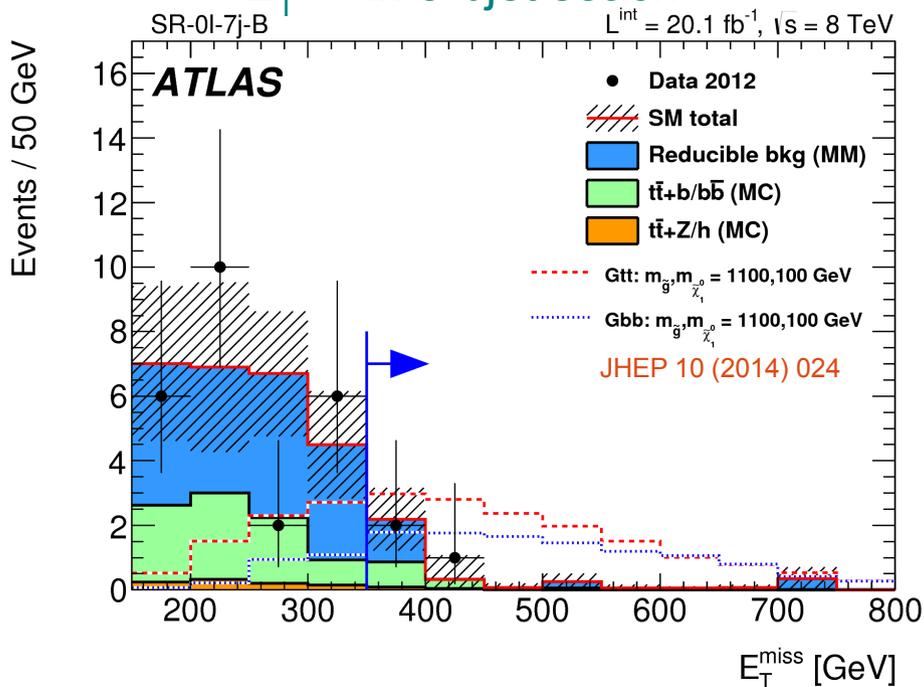
Decays through stop

- If scalar top lighter than other squarks, gluinos will decay to multiple top quarks
- Provides a very rich final state: 4-bjets, up to 4 leptons, up to 12 jets
- Many SUSY searches sensitive to this channel
- Best limit from a search for 3-bjets:



Gluino heavier than ~1.4 TeV

E_T^{miss} in 3-bjet search





Search for Di-lepton Final-state

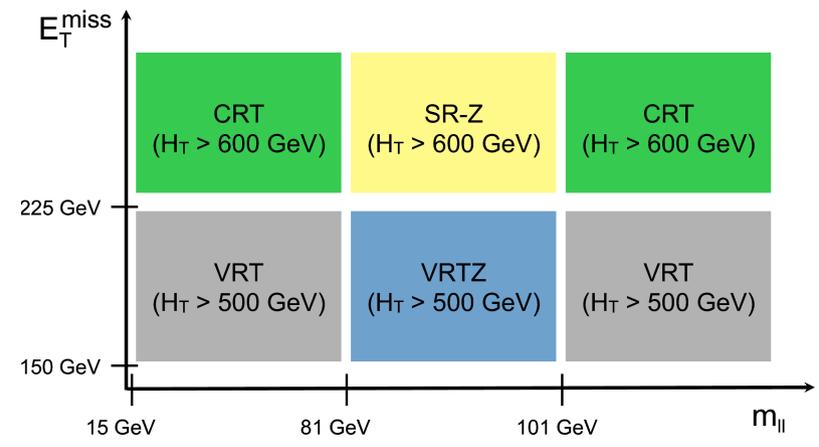
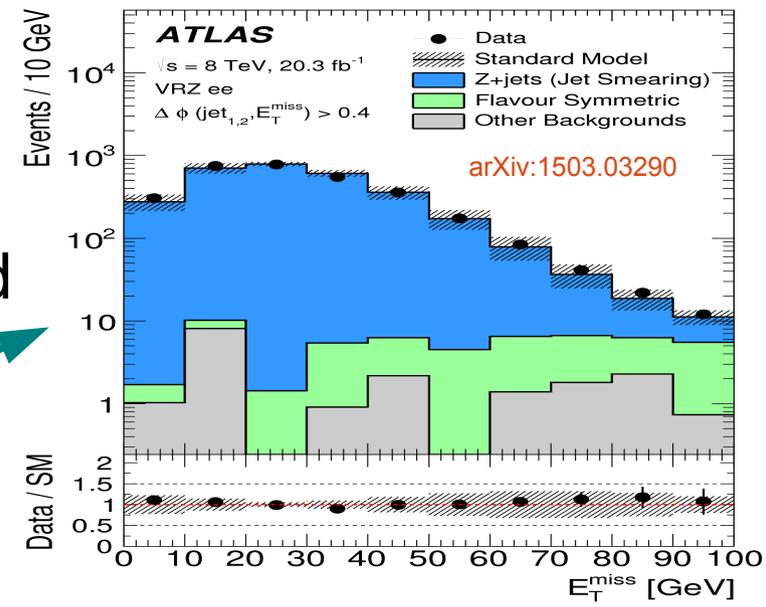
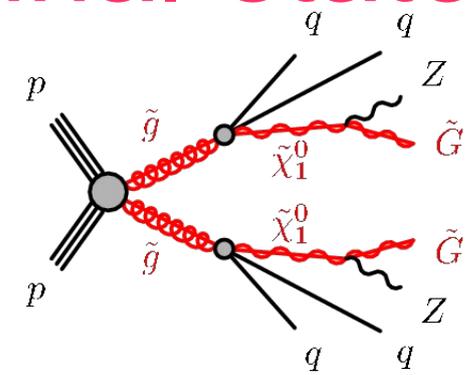
Many more decays possible for gluinos

Search for decays through $Z \rightarrow ll$:

- Select di-leptons with $|m_{ll} - m_Z| < 10$ GeV
- At least two jets ($p_T > 35$ GeV)
- $E_T^{miss} > 225$ GeV
- $t\bar{t}$ background estimated from data using $e\mu$ events
- Cross check in $|m_{ll} - m_Z| > 10$ GeV sideband
- Z+fake MET from jet smearing method
- Multiple validation regions at low E_T^{miss}

Signal, control and validation regions

On-Z Region	E_T^{miss} [GeV]	H_T [GeV]	n_{jets}	$m_{\ell\ell}$ [GeV]	SF/DF	E_T^{miss} sig. [$\sqrt{\text{GeV}}$]	f_{ST}	$\Delta\phi(\text{jet}_{1,2}, E_T^{miss})$
Signal regions								
SR-Z	> 225	> 600	≥ 2	$81 < m_{\ell\ell} < 101$	SF	-	-	> 0.4
Control regions								
Seed region	-	> 600	≥ 2	$81 < m_{\ell\ell} < 101$	SF	< 0.9	< 0.6	-
CRe μ	> 225	> 600	≥ 2	$81 < m_{\ell\ell} < 101$	DF	-	-	> 0.4
CRT	> 225	> 600	≥ 2	$m_{\ell\ell} \notin [81, 101]$	SF	-	-	> 0.4
Validation regions								
VRZ	< 150	> 600	≥ 2	$81 < m_{\ell\ell} < 101$	SF	-	-	-
VRT	150-225	> 500	≥ 2	$m_{\ell\ell} \notin [81, 101]$	SF	-	-	> 0.4
VRTZ	150-225	> 500	≥ 2	$81 < m_{\ell\ell} < 101$	SF	-	-	> 0.4



New

Excess in Z+jets Final-state

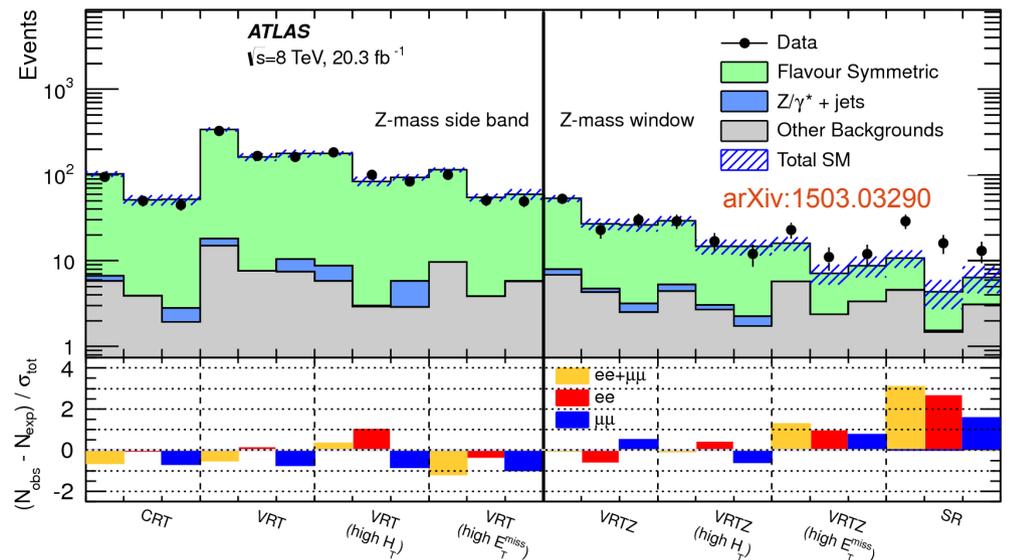
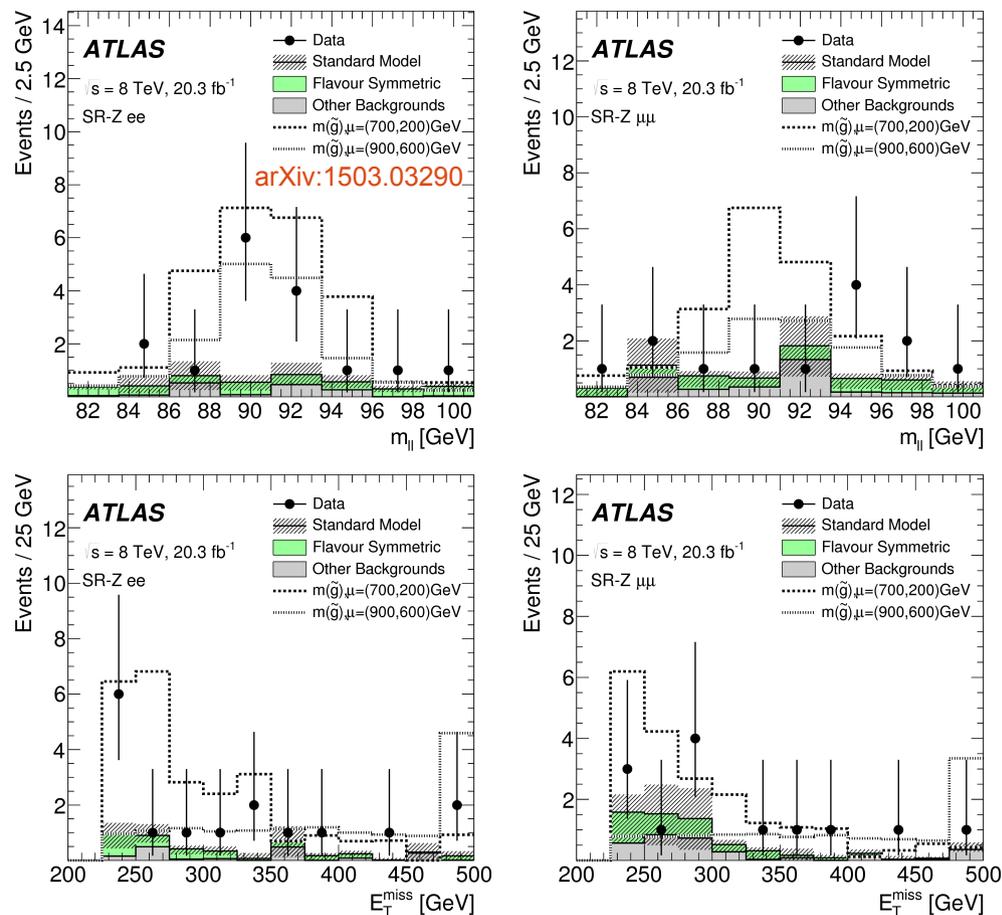
Observed events and expected backgrounds

Observe an excess in data:

- 3σ significance in ee channel
- 1.7σ in $\mu\mu$ channel
- 3σ significance when combining the two channels

Channel	SR-Z ee	SR-Z $\mu\mu$	SR-Z same-flavour combined
Observed events	16	13	29
Expected background events	4.2 ± 1.6	6.4 ± 2.2	10.6 ± 3.2
Flavour-symmetric backgrounds	2.8 ± 1.4	3.3 ± 1.6	6.0 ± 2.6
Z/ γ^* + jets (jet-smearing)	0.05 ± 0.04	$0.02^{+0.03}_{-0.02}$	0.07 ± 0.05
Rare top	0.18 ± 0.06	0.17 ± 0.06	0.35 ± 0.12
WZ/ZZ diboson	1.2 ± 0.5	1.7 ± 0.6	2.9 ± 1.0
Fake leptons	$0.1^{+0.7}_{-0.1}$	$1.2^{+1.3}_{-1.2}$	$1.3^{+1.7}_{-1.3}$

Excess of events only in signal region:



New

Excess in Z+jets Final-state

Observed events and expected backgrounds

Observe an excess in data:

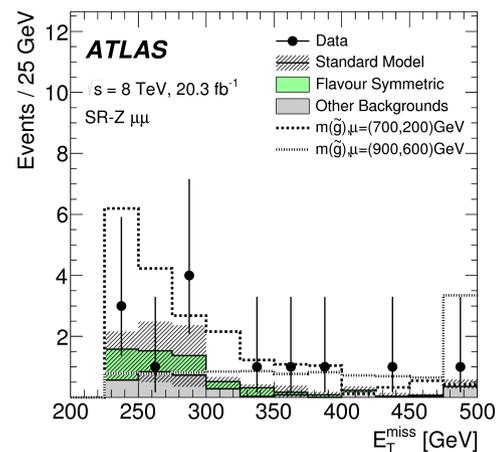
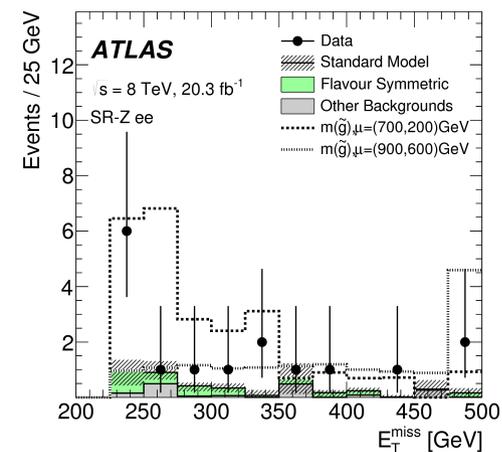
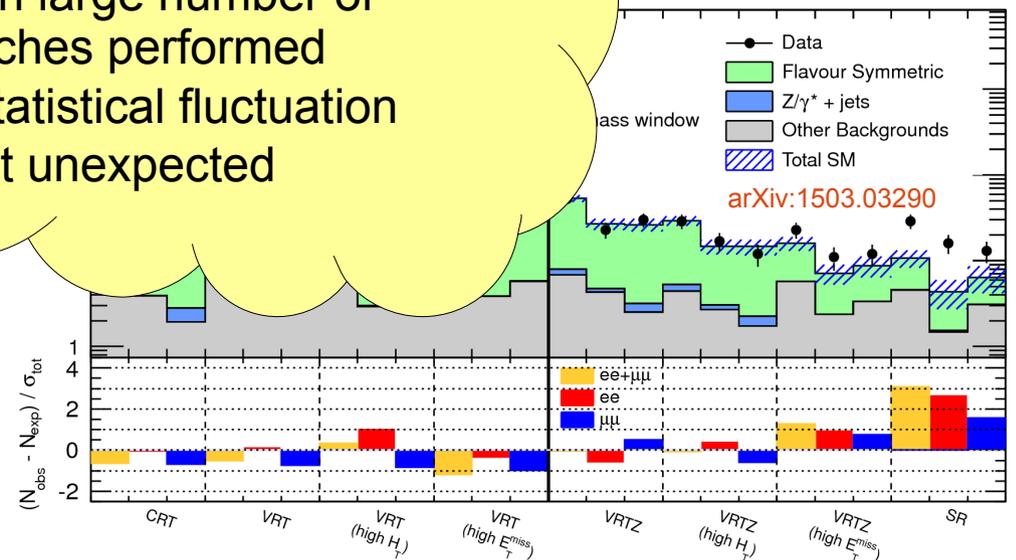
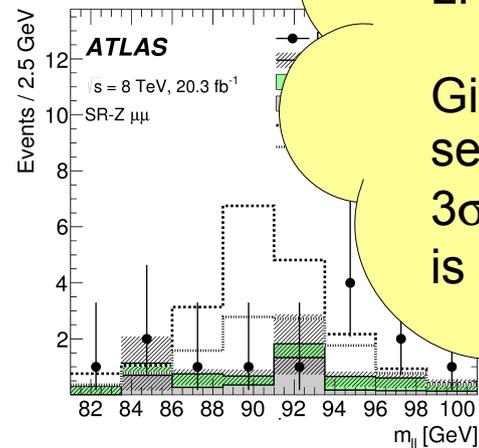
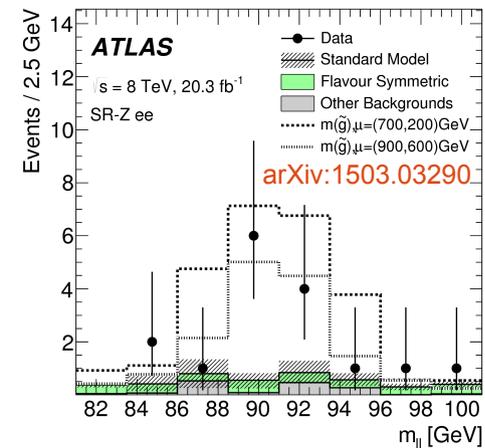
- 3σ significance in ee channel
- 1.7σ in $\mu\mu$ channel
- 3σ significance when combining the two channels

Channel	SR-Z ee	SR-Z $\mu\mu$	SR-Z same-flavour combined
Observed events	16	13	29
Expected background events	4.2 ± 1.6	6.4 ± 2.2	10.6 ± 3.2
Flavour-symmetric backgrounds	2.8 ± 1.4	3.3 ± 1.6	6.0 ± 2.6
Z/ γ^* + jets (jet-sm)	0.05 ± 0.04	$0.02^{+0.03}_{-0.02}$	0.07 ± 0.05
Other Backgrounds	0.18 ± 0.06	0.17 ± 0.06	0.35 ± 0.12
Total SM	1.2 ± 0.5	1.7 ± 0.6	2.9 ± 1.0
	$1.2^{+0.7}_{-0.1}$	$1.2^{+1.3}_{-1.2}$	$1.3^{+1.7}_{-1.3}$

Largest excess seen in LHC SUSY searches

Given large number of searches performed
 3σ statistical fluctuation is not unexpected

in signal region:



New

Does CMS See Excess?

Recent similar search in CMS:

- Similar $Z \rightarrow \ell\ell$ selection
- Counts in bins of E_T^{miss}
- ≥ 2 (or ≥ 3) jets above 40 GeV
- No selection on H_T (sum of jet p_T)

≥ 2 jets

E_T^{miss} (GeV)	100–200	200–300	>300
DY background	336 ± 89	28.6 ± 8.6	7.7 ± 3.6
FS background	868 ± 57	45.9 ± 7.3	5.1 ± 2.3
Total background	1204 ± 106	74.5 ± 11.3	12.8 ± 4.3
Data	1187	65	7

GMSB signal yields

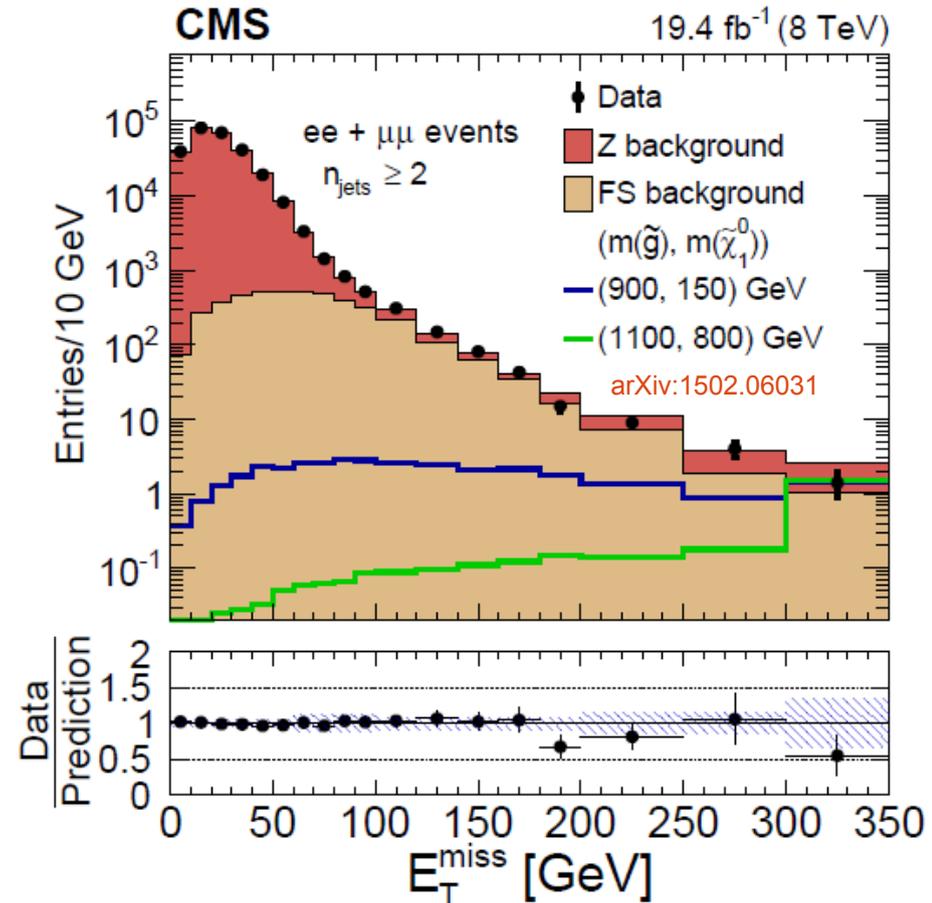
$m_{\tilde{g}} = 900, m_{\tilde{\chi}_1^0} = 150$	22.1 ± 0.4	11.1 ± 0.3	7.2 ± 0.2
$m_{\tilde{g}} = 1100, m_{\tilde{\chi}_1^0} = 800$	1.1 ± 0.04	1.6 ± 0.05	7.6 ± 0.1

≥ 3 jets

E_T^{miss} (GeV)	100–200	200–300	>300
DY background	124 ± 33	12.7 ± 3.8	3.2 ± 1.8
FS background	354 ± 28	26.5 ± 5.4	2.0 ± 1.4
Total background	478 ± 43	39.2 ± 6.6	5.3 ± 2.3
Data	490	35	6

GMSB signal yields

$m_{\tilde{g}} = 900, m_{\tilde{\chi}_1^0} = 150$	22.0 ± 0.4	11.0 ± 0.3	7.1 ± 0.2
$m_{\tilde{g}} = 1100, m_{\tilde{\chi}_1^0} = 800$	1.1 ± 0.04	1.5 ± 0.05	7.4 ± 0.1



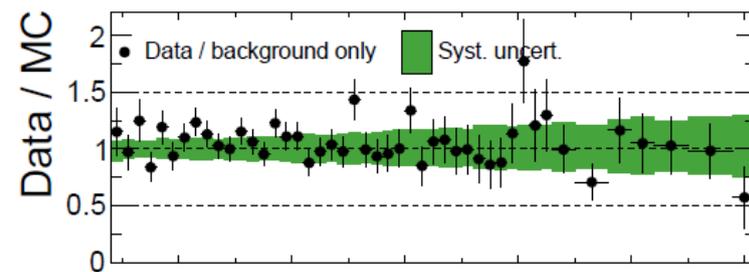
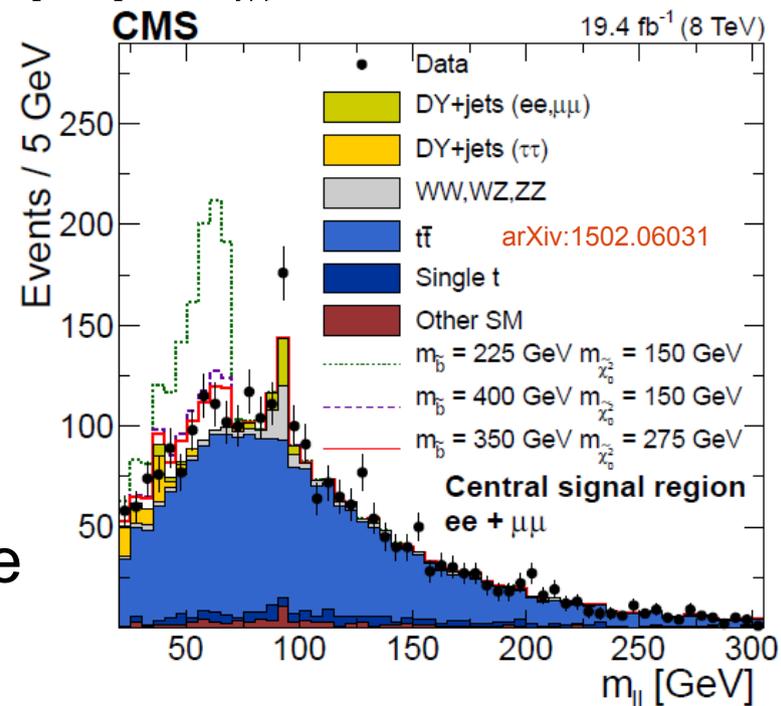
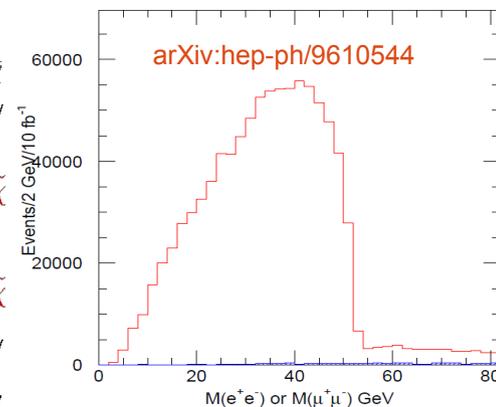
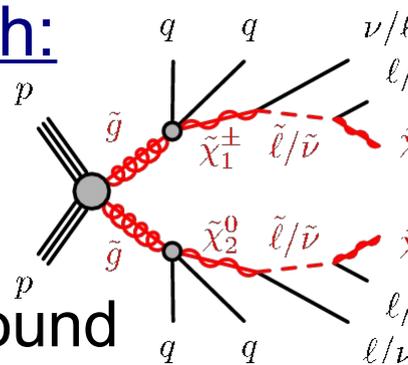
No excess seen by CMS
 However, only ~30% overlap
 with ATLAS selection

New

CMS Excess in Dilepton Search

CMS excess in dilepton edge search:

- Decay of heavy neutralino through slepton or off-shell Z gives triangular m_{ll} mass distribution with characteristic edge and kinematic bound
- Split in central leptons ($|\eta| < 1.4$) and forward leptons ($|\eta| > 1.6$)
- 2 jets and $E_T^{\text{miss}} > 150$ GeV, or 3 jets and $E_T^{\text{miss}} > 100$ GeV
- Fit to “edge” shape across m_{ll} and count in 3 m_{ll} bins
- See 2.6σ excess at low m_{ll} in central case



	Low-mass		On-Z		High-mass	
	Central	Forward	Central	Forward	Central	Forward
Observed	860	163	487	170	818	368
Flavor-symmetric	$722 \pm 27 \pm 29$	$155 \pm 13 \pm 10$	$355 \pm 19 \pm 14$	$131 \pm 12 \pm 8$	$768 \pm 28 \pm 31$	$430 \pm 22 \pm 27$
Drell-Yan	8.2 ± 2.6	2.5 ± 1.0	116 ± 21	42 ± 9	2.5 ± 0.8	1.1 ± 0.4
Total estimated	730 ± 40	158 ± 16	471 ± 32	173 ± 17	771 ± 42	431 ± 35
Observed-estimated	130^{+48}_{-49}	5^{+20}_{-20}	16^{+37}_{-38}	-3^{+20}_{-21}	47^{+49}_{-50}	-62^{+37}_{-39}
Significance	2.6σ	0.3σ	0.4σ	$<0.1\sigma$	0.9σ	$<0.1\sigma$

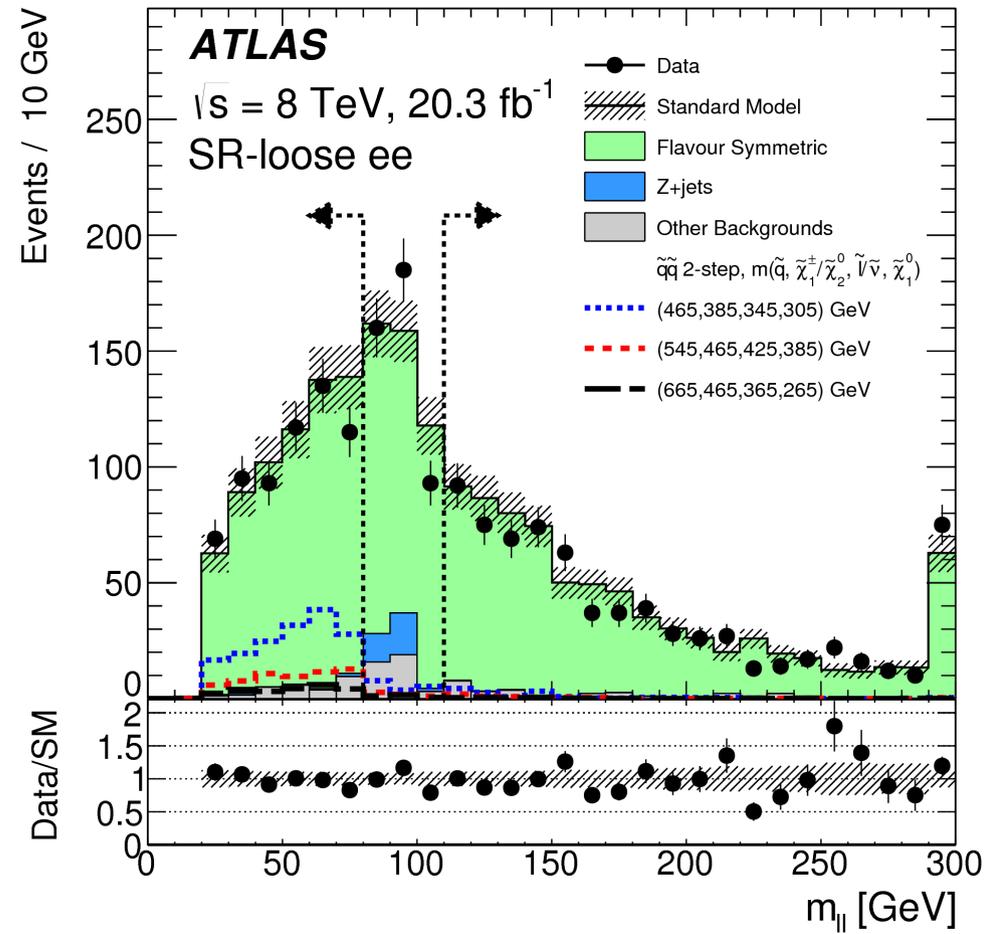
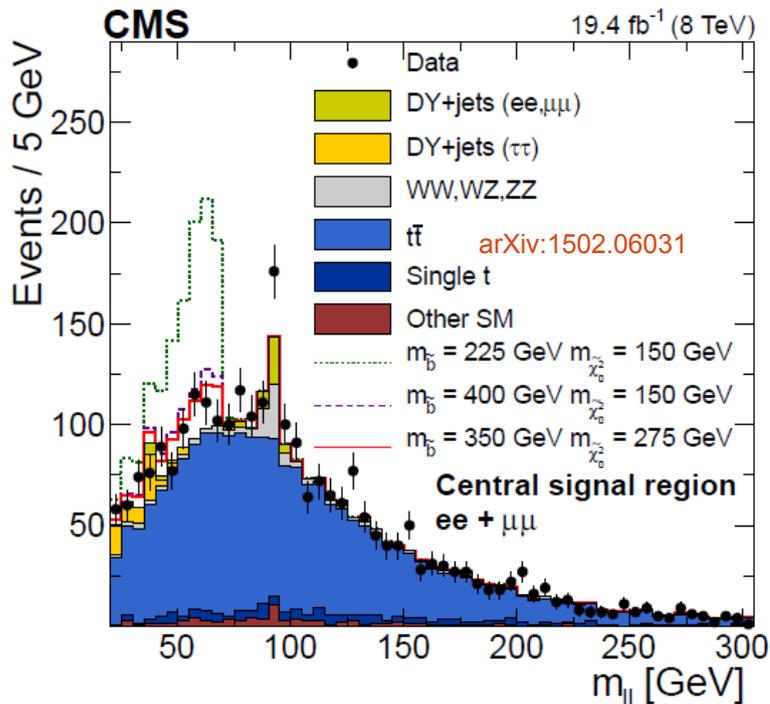
New

Dilepton Edge in ATLAS

ATLAS also has edge search:

- One set of selection almost identical to CMS
- Additional SRs split in b-jets and no bjets and minimum 2 vs 4 jets
- No excess seen in ATLAS

Below-Z ($20 < m_{\ell\ell} < 70$ GeV)	SR-loose ee	SR-loose $\mu\mu$	SR-loose same-flavour combined
Observed events	509	624	1133
Expected background events	$510 \pm 20 \pm 40$	$680 \pm 20 \pm 50$	$1190 \pm 40 \pm 70$
Flavour-symmetric backgrounds	$490 \pm 20 \pm 40$	$650 \pm 20 \pm 50$	$1140 \pm 40 \pm 70$
$Z/\gamma^* + \text{jets}$	$2.5 \pm 0.8 \pm 3.2$	$8 \pm 2 \pm 5$	$11 \pm 2 \pm 7$
Rare top	$0.3 \pm 0.0 \pm 0.0$	$0.4 \pm 0.0 \pm 0.0$	$0.7 \pm 0.0 \pm 0.0$
WZ/ZZ	$1.1 \pm 0.3 \pm 0.1$	$1.2 \pm 0.2 \pm 0.4$	$2.4 \pm 0.4 \pm 0.4$
Fake leptons	$16 \pm 4 \pm 2$	$23 \pm 5 \pm 1$	$38 \pm 6 \pm 4$



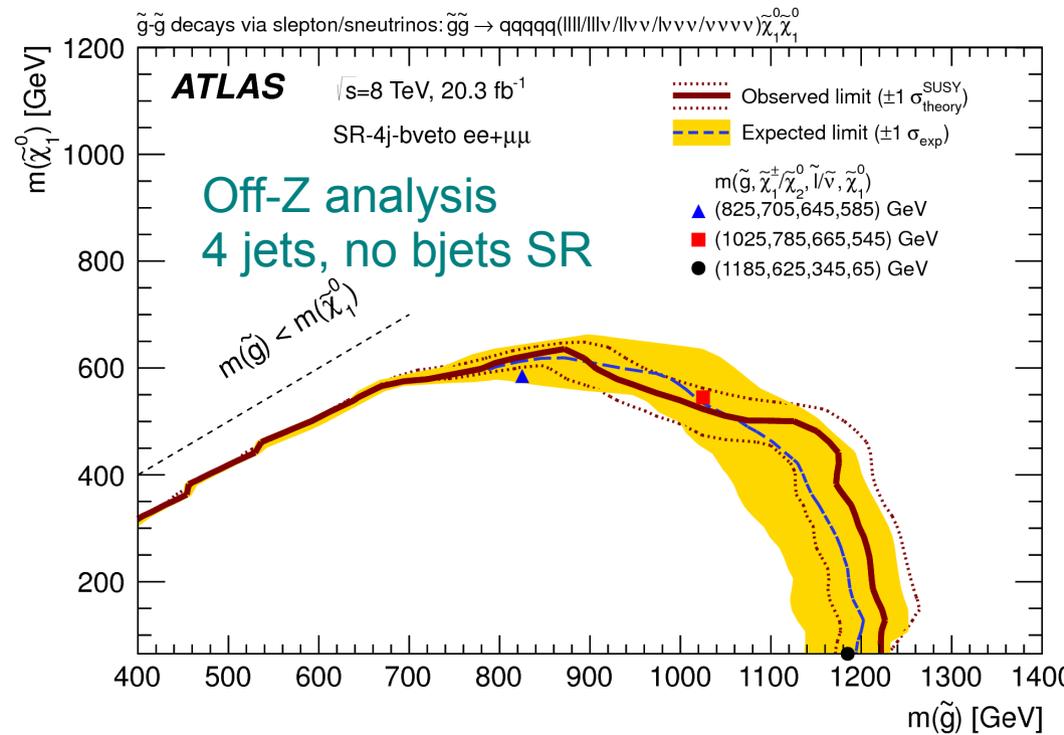
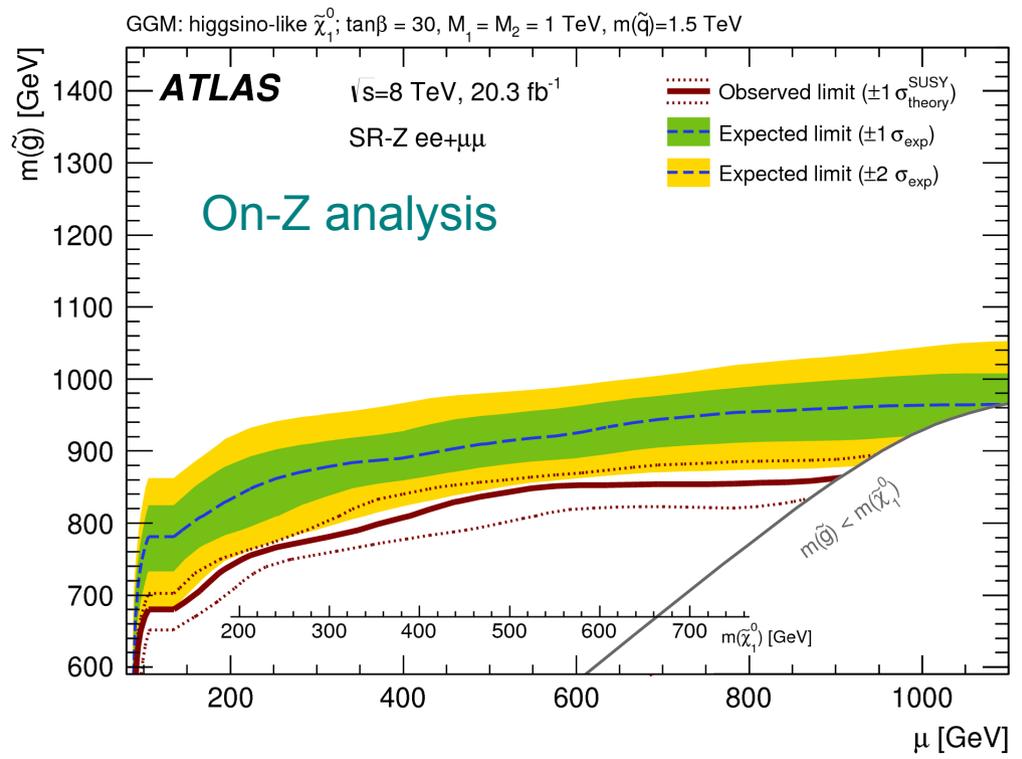
New Gluino Limits with Dileptons

Still put limits on gluinos:

- Analysis results interpreted in both simplified models and Generalized gauge mediated supersymmetry models (GGM)
- Limits weakened due to excess in on-Z search

	Dilepton edge	Z+MET
ATLAS	No excess	3.0 σ
CMS	2.6 σ	No excess

The ATLAS and CMS edge selections are the same (by design) but the Z+MET are different, only ~30% of our events enter the CMS selection

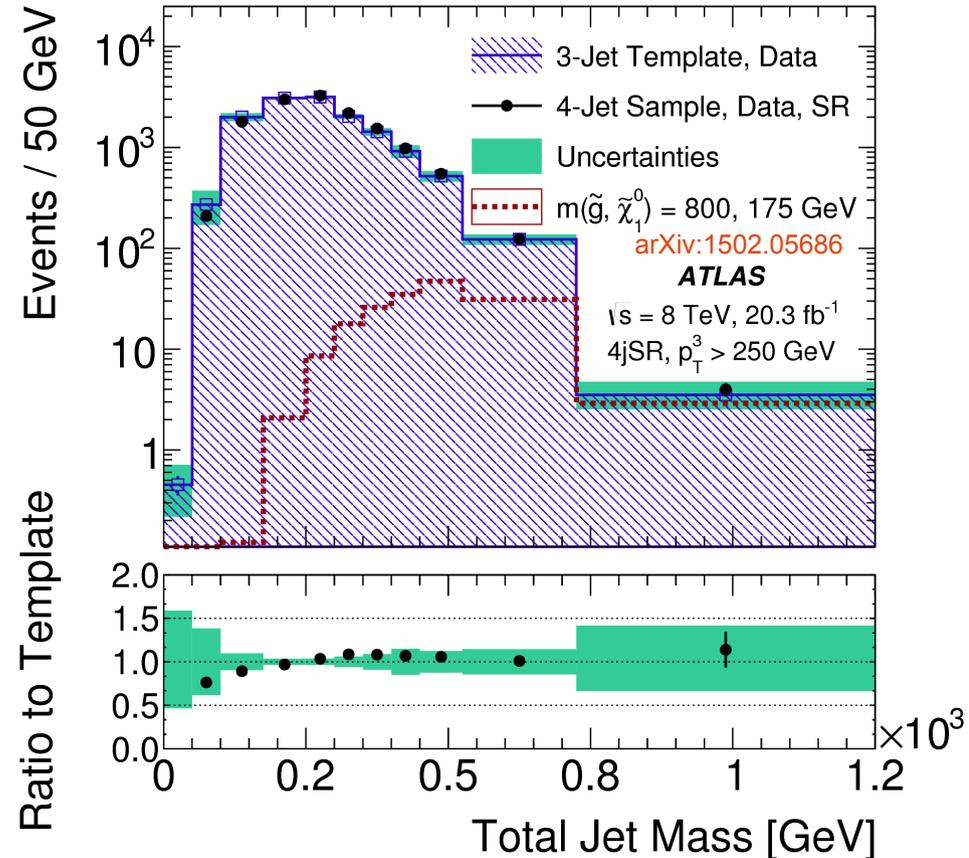
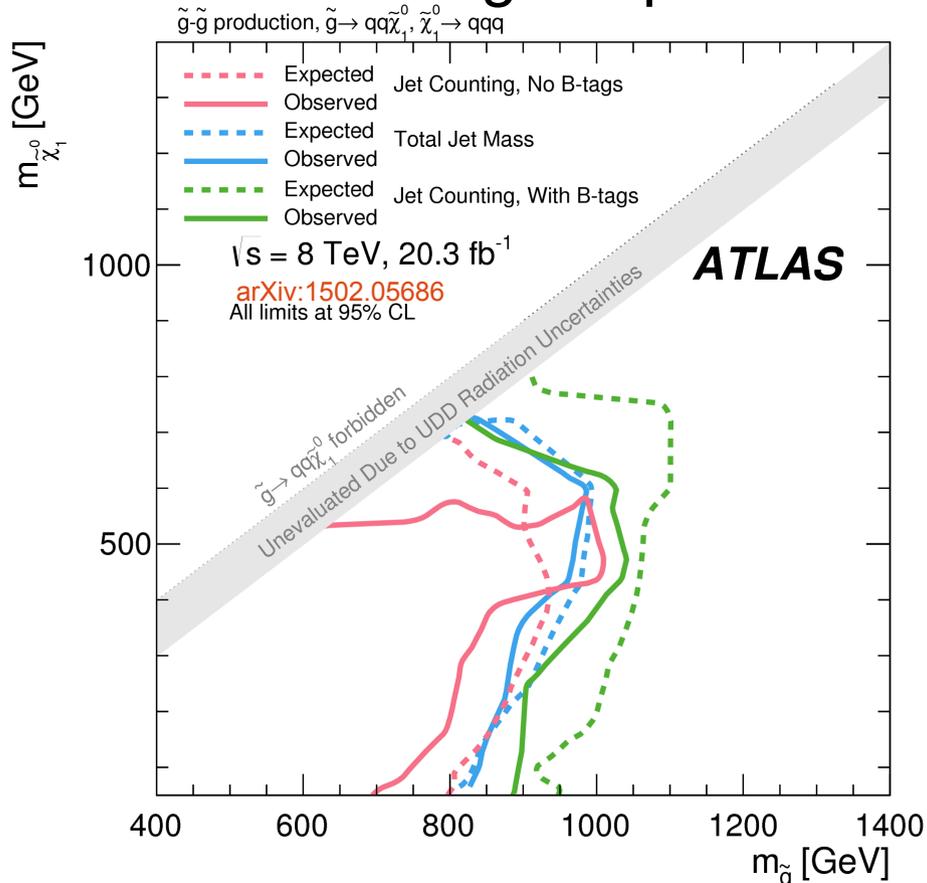
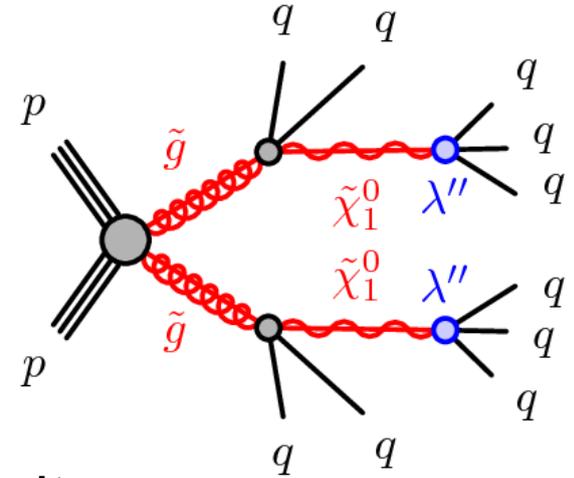


New

Glauino Search without E_T^{miss}

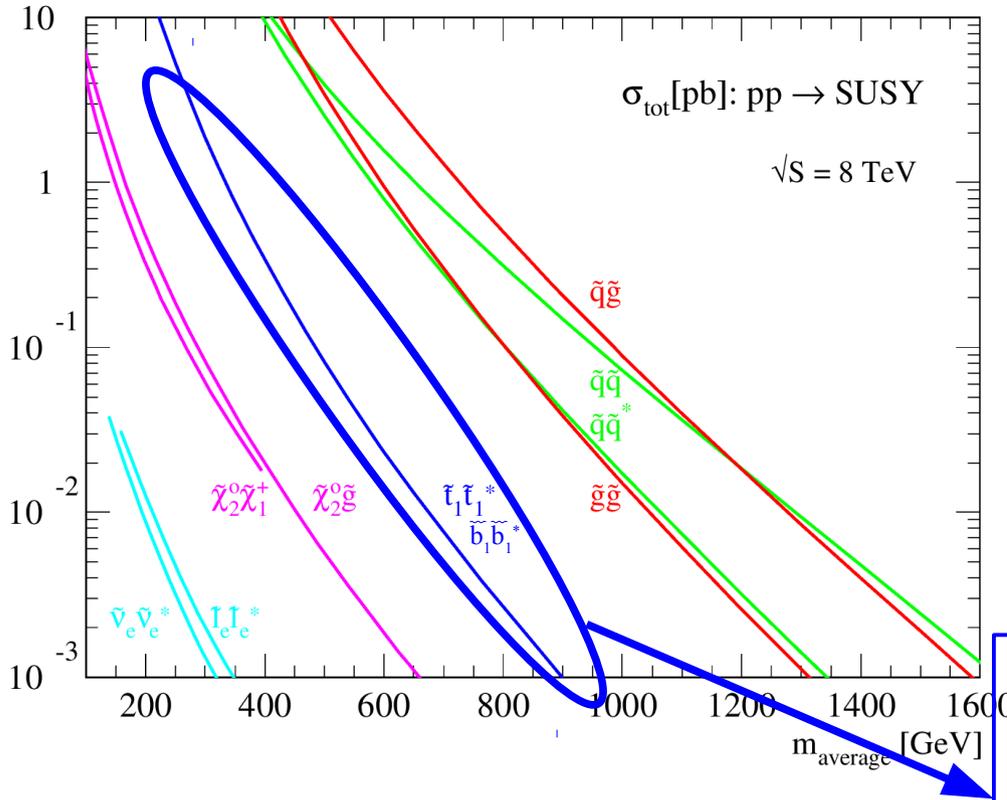
Search for E_T^{miss} less gluino decays:

- If R-parity violated can have decays to just jets
- Search uses high jet multiplicity ($\geq 6, 7$ jets) or large total (fat-)jet mass
- Dominant background from SM multi-jets estimated using templates from lower jet multiplicity



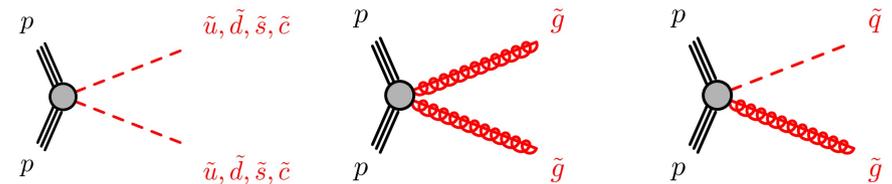
LHC SUSY Searches

At LHC can search for production for almost all SUSY particles, but with different sensitivity as production cross-sections vary



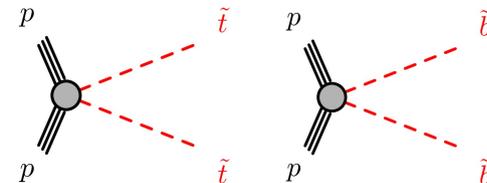
Production of squarks/gluinos:

- Strong coupling gives large sensitivity at the LHC
- Primary signature is jets and missing energy



Prod. of 3rd generation squarks:

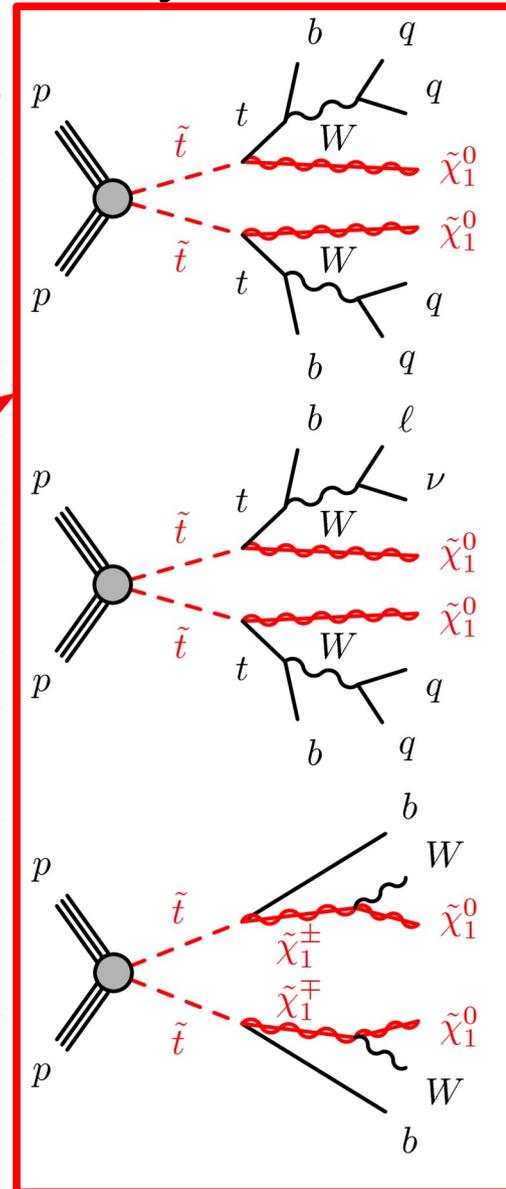
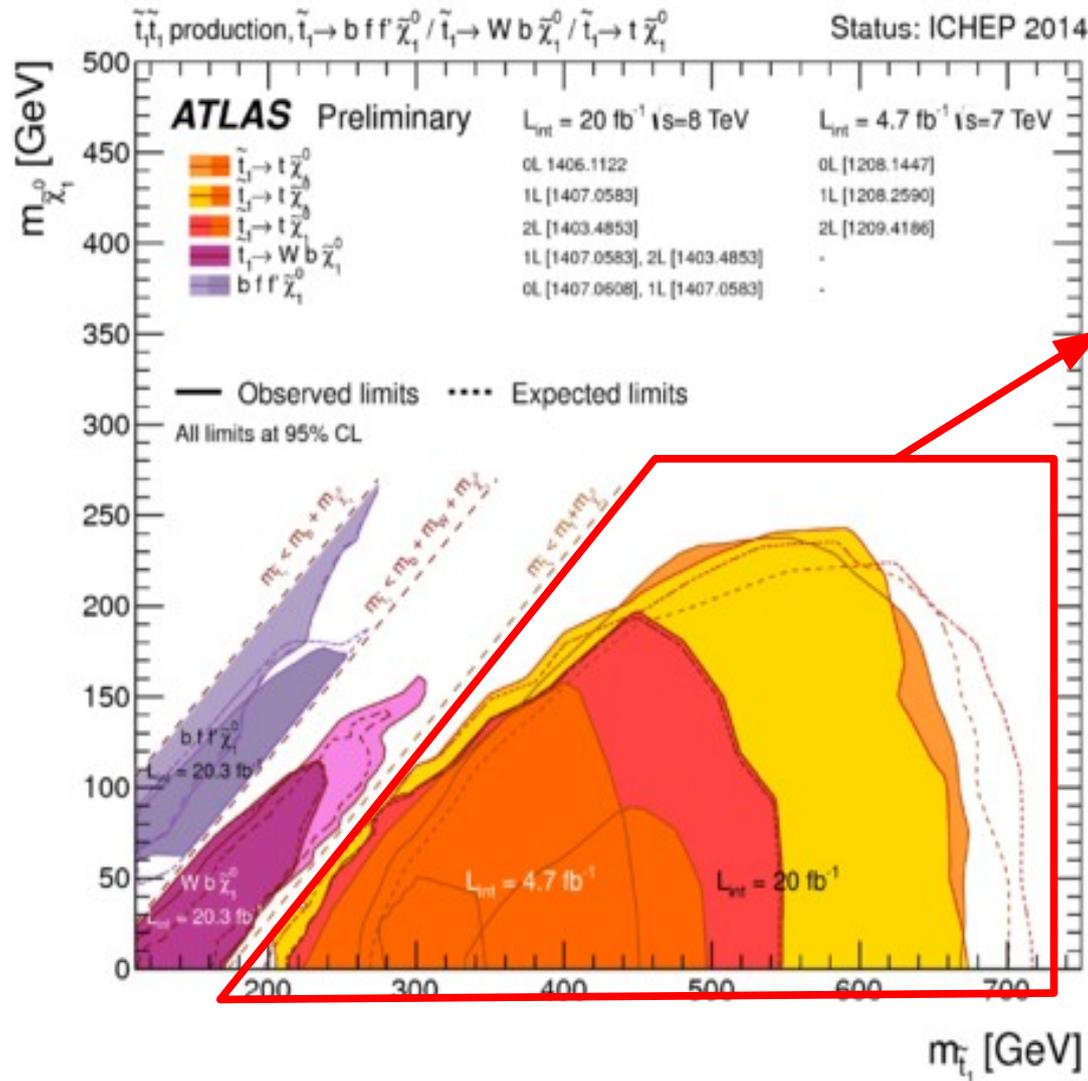
- Lower cross section,
- Decays chains frequently has top and bottom quarks
- Of great interest to “Natural SUSY”



3rd Generation Squark Searches

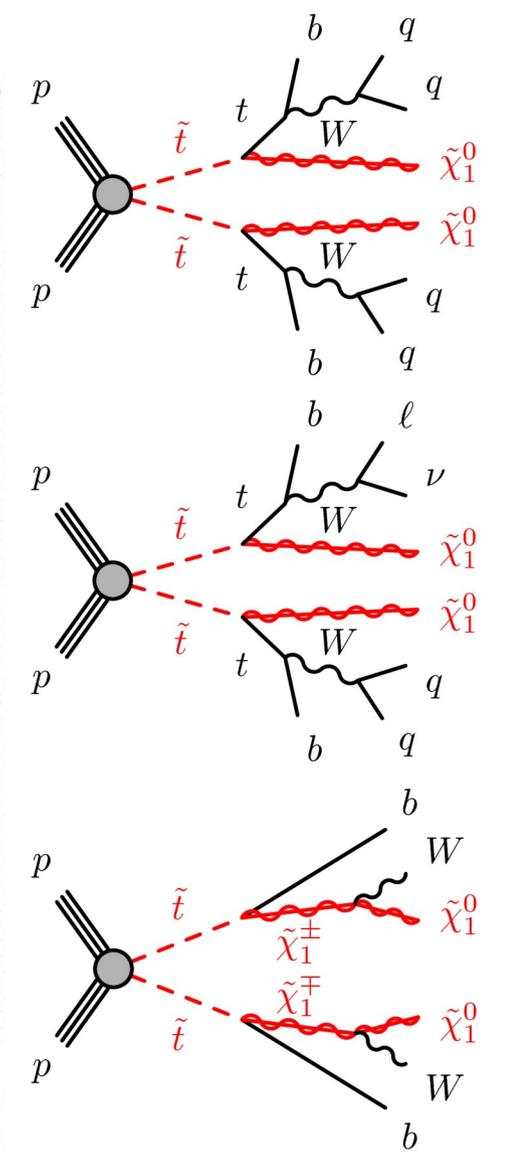
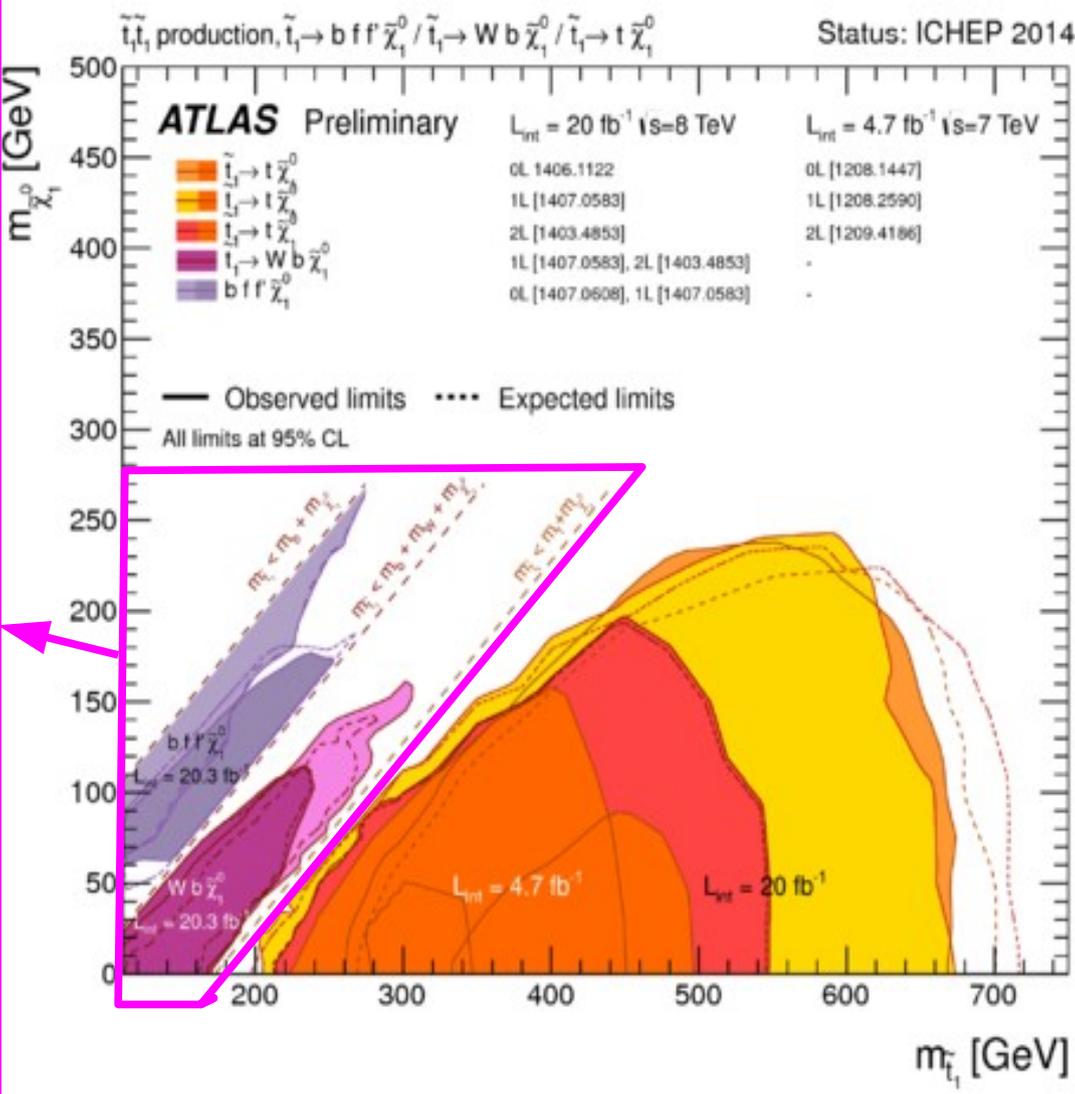
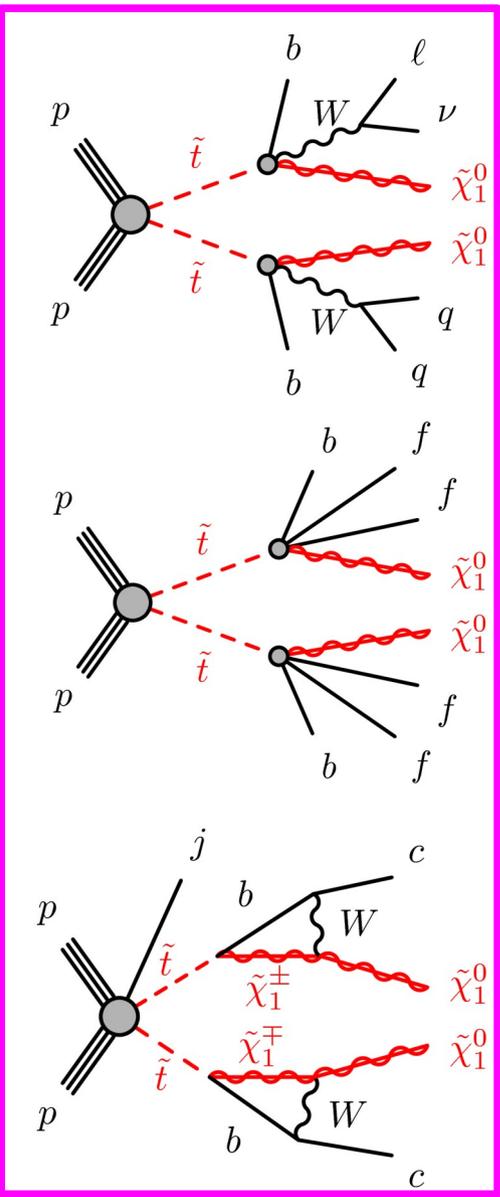
In natural SUSY at least one stop light ($m < 1$ TeV)

Have been a major focus of SUSY searches in last 3 years



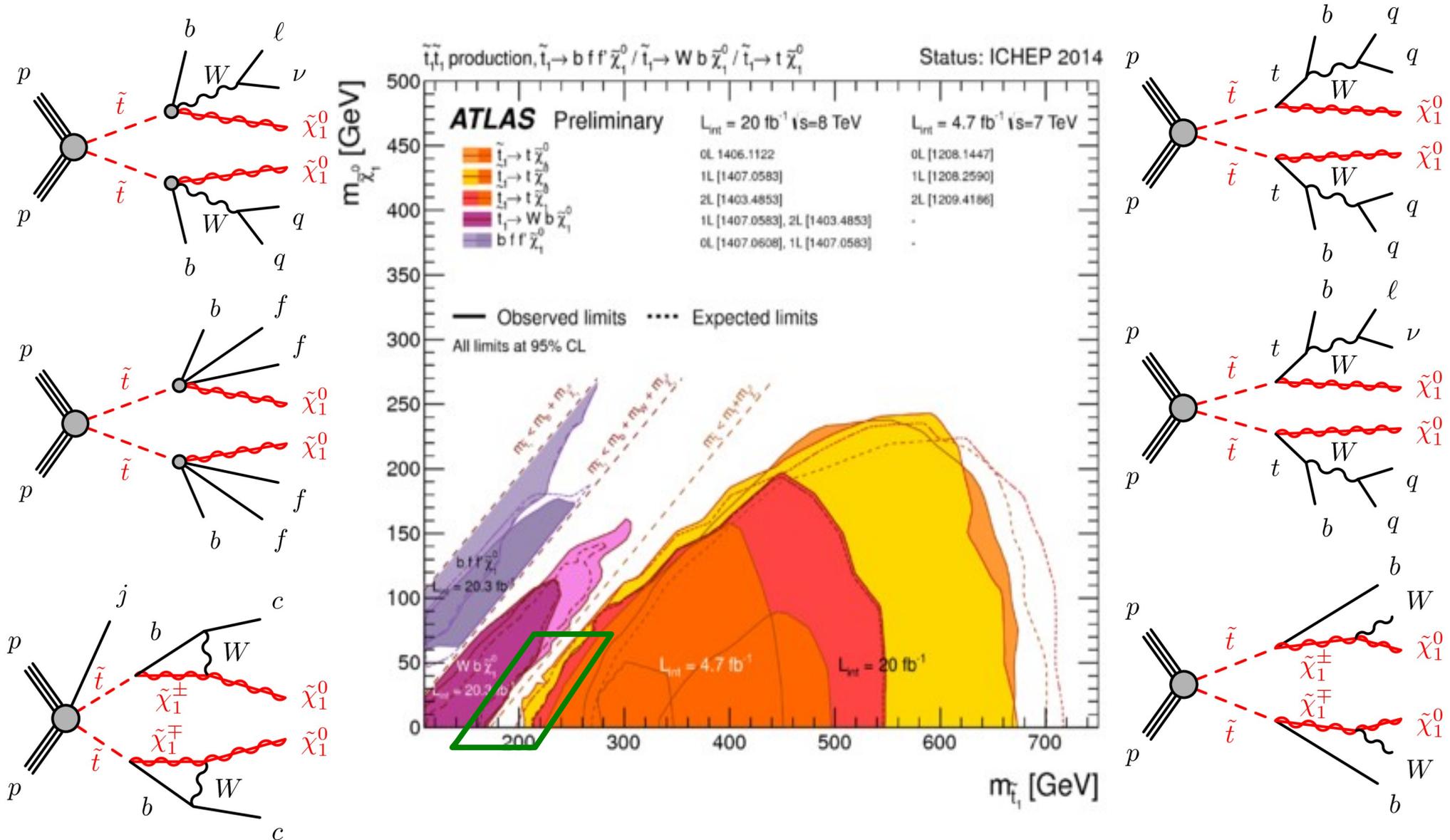
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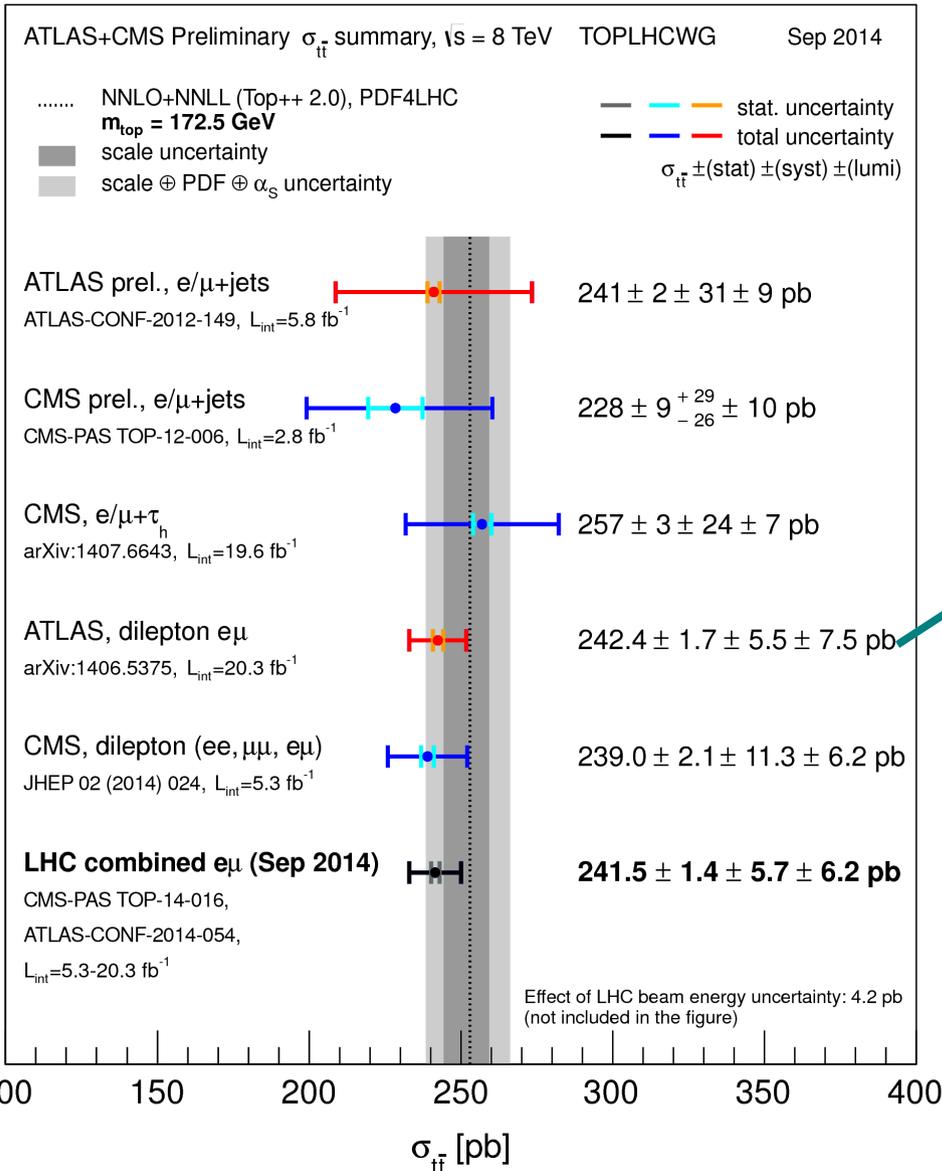
3rd Generation Squark Searches

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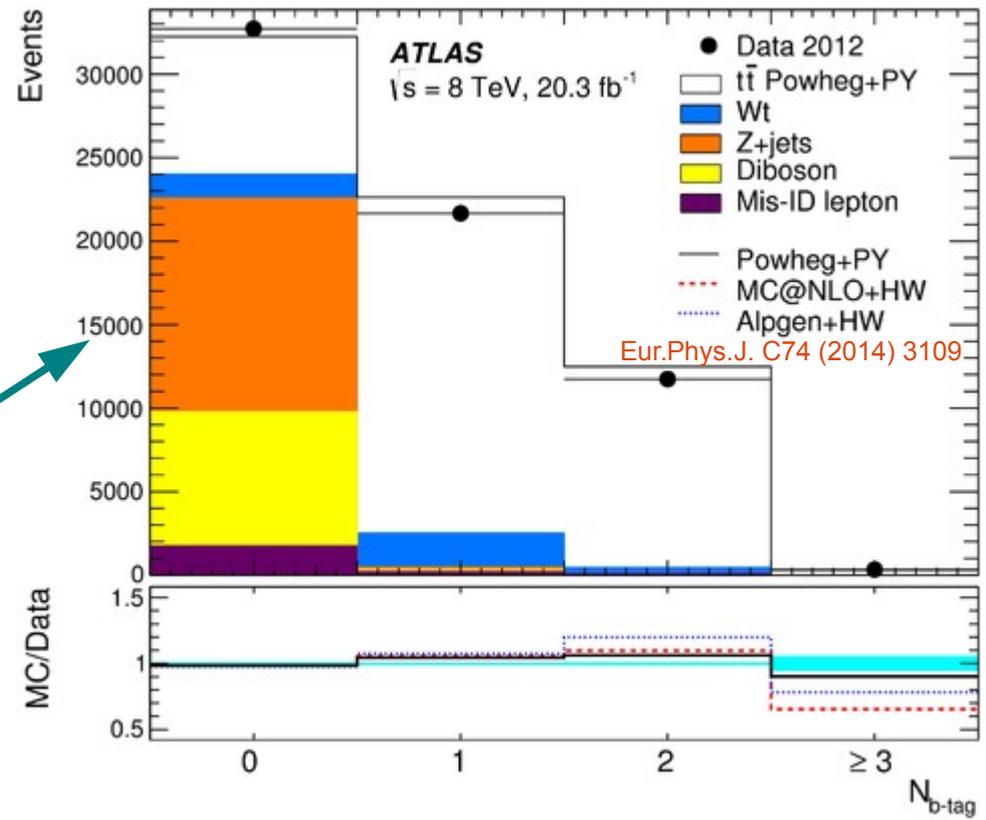


Stop Limit from $t\bar{t}$ cross section

- If $m_{\tilde{t}} \sim m_{\text{top}}$, then it is reconstructed as top, increasing cross section
- Precise NNLO+NNLL SM prediction: $177.3^{+11.5}_{-12.0}$ pb (7 TeV)
 $252.9^{+15.3}_{-16.3}$ pb (8 TeV)



ATLAS $t\bar{t}$ dilepton cross section measurement



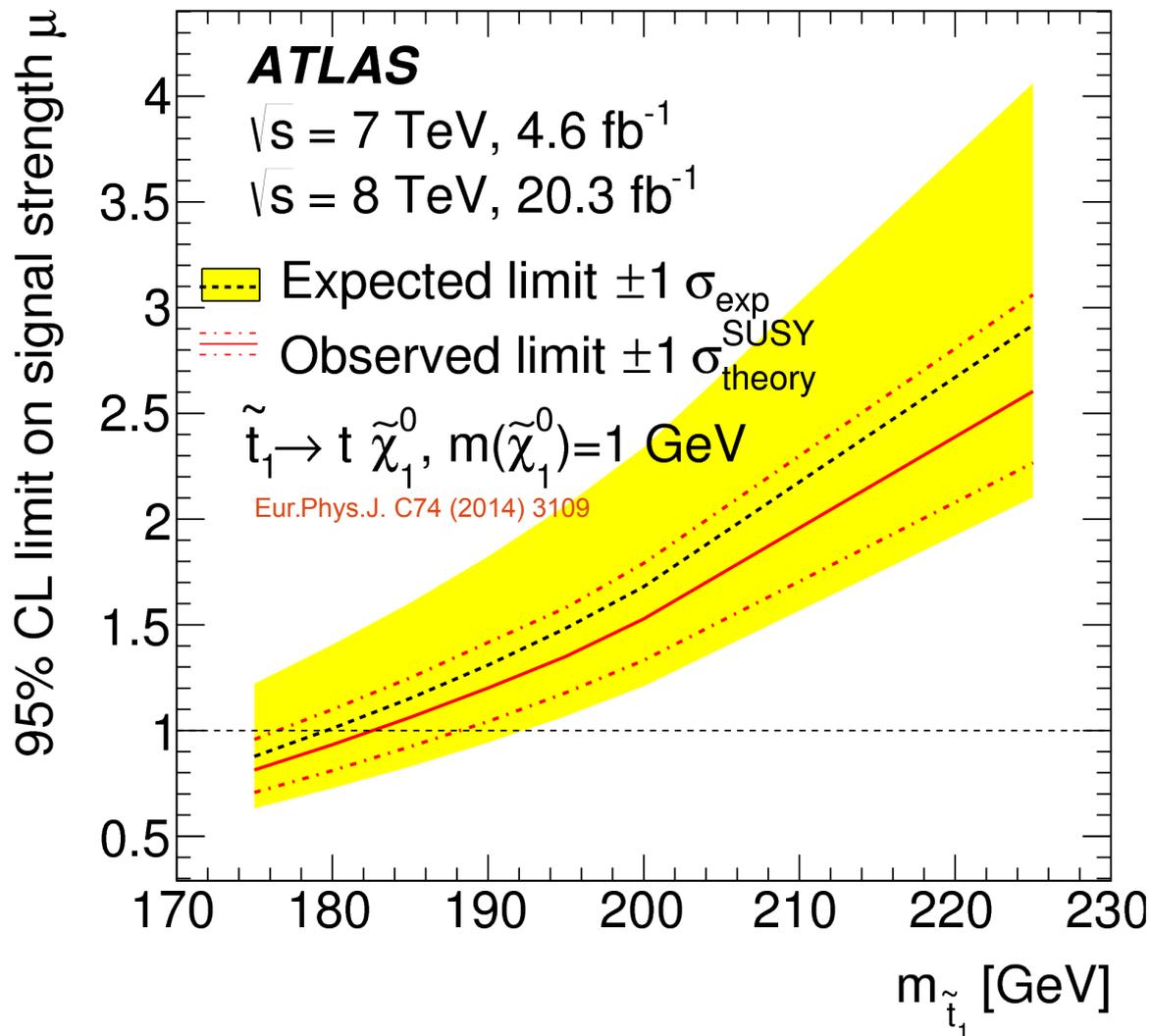
$$\sigma_{t\bar{t}} = 182.9 \pm 3.1 \pm 4.2 \pm 3.6 \pm 3.3 \text{ pb} (\sqrt{s} = 7 \text{ TeV})$$

$$\sigma_{t\bar{t}} = 242.4 \pm 1.7 \pm 5.5 \pm 7.5 \pm 4.2 \text{ pb} (\sqrt{s} = 8 \text{ TeV}),$$

Stop Limit from $t\bar{t}$ cross section

$m_{\tilde{t}} \sim m_{\text{top}}$ excluded

- $\sigma_{t\bar{t}} \sim 40$ pb at $m_{\tilde{t}} = 175$ GeV,
- $\sigma_{t\bar{t}} \sim 20$ pb at $m_{\tilde{t}} = 200$ GeV
- Selection efficiency is very similar to $t\bar{t}$ for right-handed \tilde{t}
- Exclude light stop from $m_{\tilde{t}}$ to 183 (177) GeV before(after) accounting for 15% uncertainty on $\sigma_{t\bar{t}}$
- Limit assumes $\text{BF}(\tilde{t} \rightarrow t\tilde{\chi}_1^0) = 100\%$
- Weakens to 175 GeV if stop is left-handed
- Little dependence on the neutralino mass



New

Stop Limit from $t\bar{t}$ Spin Correlation

Can also use difference in spin to put limit on light stop

$t\bar{t}$ spin correlation measurement

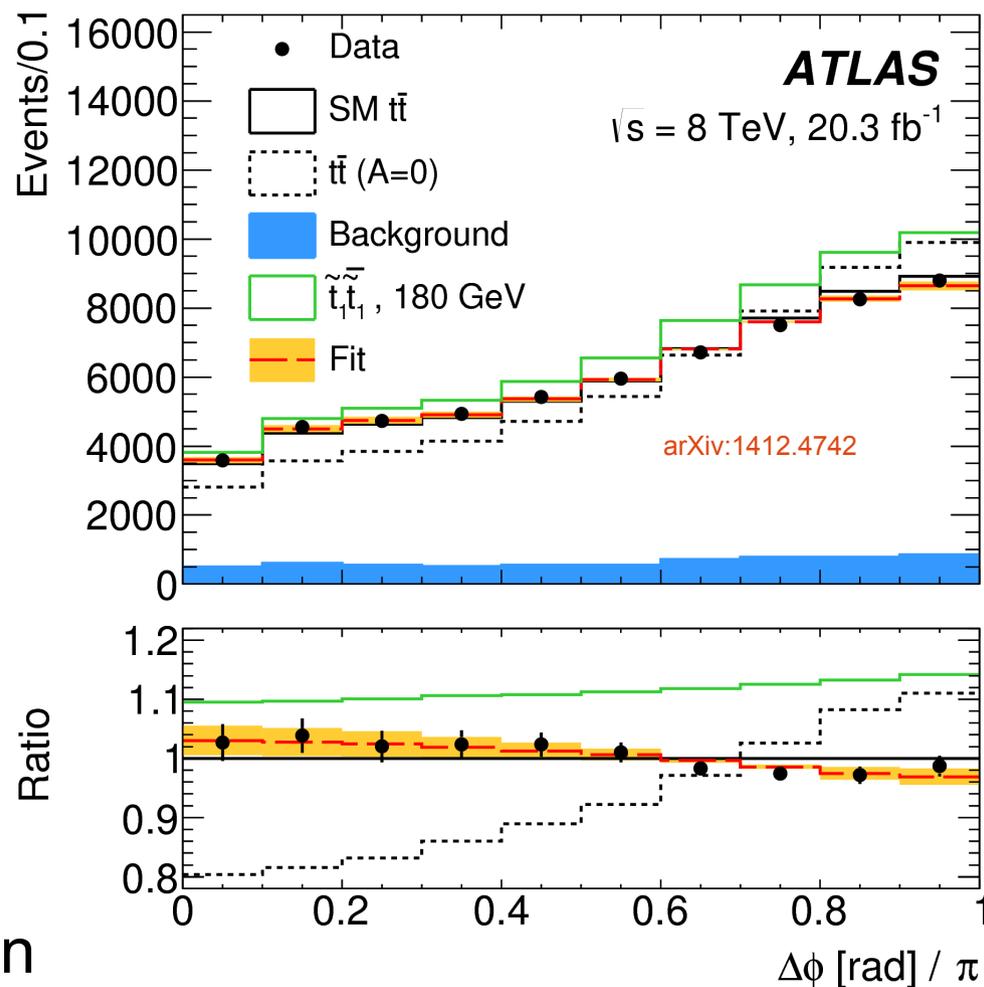
- Opposite-sign dileptons ($e\bar{e}$, $e\bar{\mu}$, $\mu\bar{\mu}$) at 8 TeV
- At least two jets, ≥ 1 b-tagged
- Further kinematic selections on E_T^{miss} and $m(\text{ll})$ to ensure clean top decay samples
- Measure spin-correlation

$$A = \frac{N_{\uparrow\uparrow} + N_{\down\down} - N_{\uparrow\down} - N_{\down\uparrow}}{N_{\uparrow\uparrow} + N_{\down\down} + N_{\uparrow\down} + N_{\down\uparrow}}$$

using azimuthal $\Delta\phi(i,j)$ distribution

- Consistent with SM prediction:

$$A_{\text{helicity}} = 0.38 \pm 0.04 \quad A_{\text{SM}} = 0.318 \pm 0.005$$

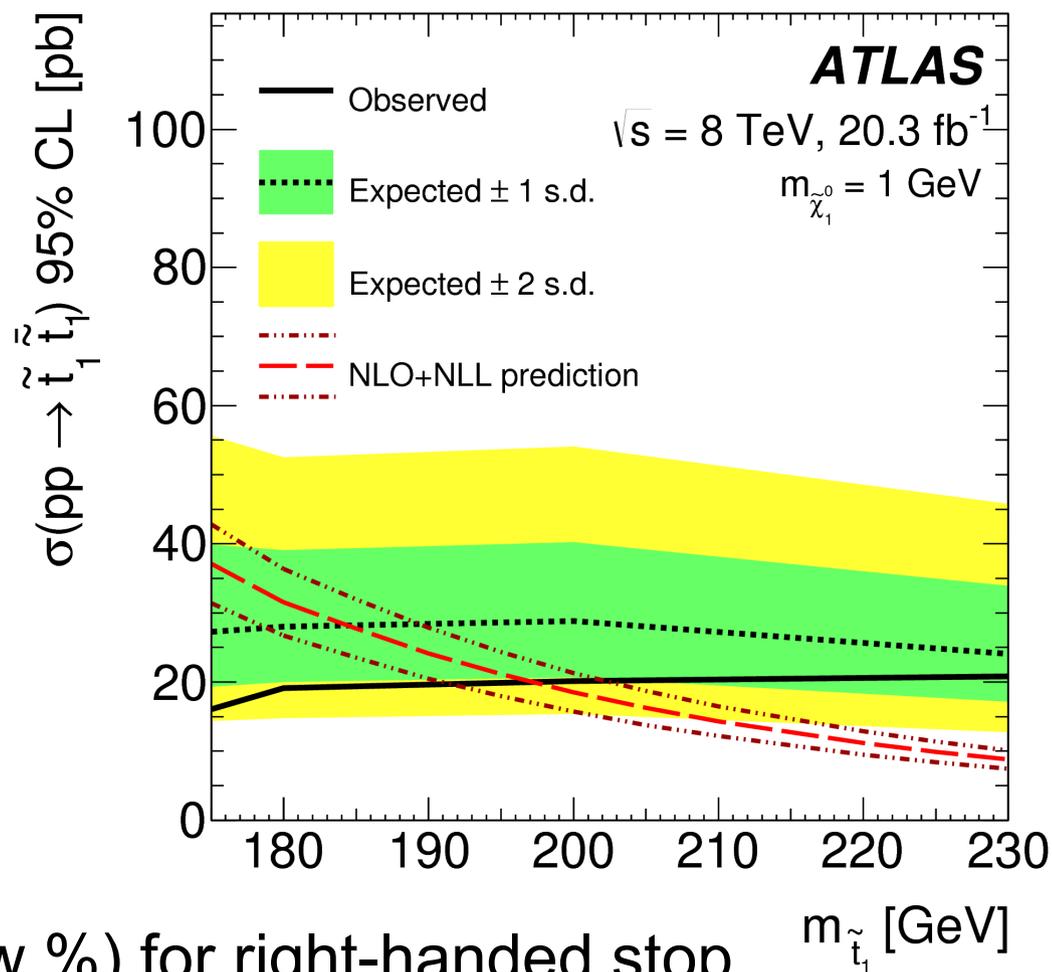


New

Stop Limit from $t\bar{t}$ Spin Correlation

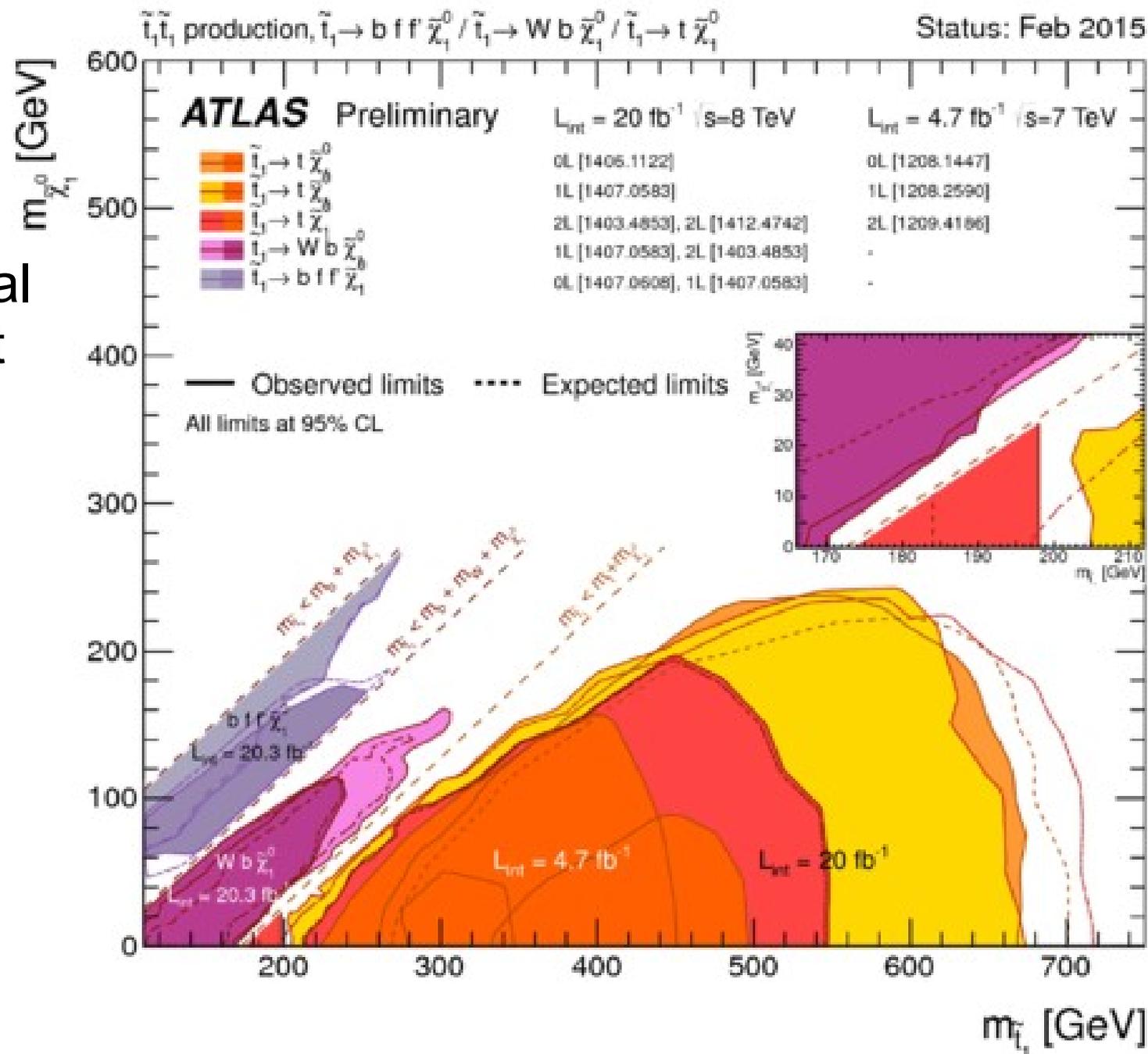
$m_{\tilde{t}_1} \sim m_{\text{top}}$ further excluded

- Simultaneous fit to overall $\sigma_{t\bar{t}}$ and $\Delta\phi(i,j)$ distribution for additional scalar top contribution assuming SM cross section for top
- Exclude light stop from $m_{\tilde{t}_1}$ to 197 (191) GeV before(after) accounting for 15% uncertainty on $\sigma_{t\bar{t}}$
- Limit assumes $\text{BF}(\tilde{t} \rightarrow t\tilde{\chi}_1^0) = 100\%$ and right-handed stop
- Only slightly weaker limit (few %) for right-handed stop and heavier neutralino



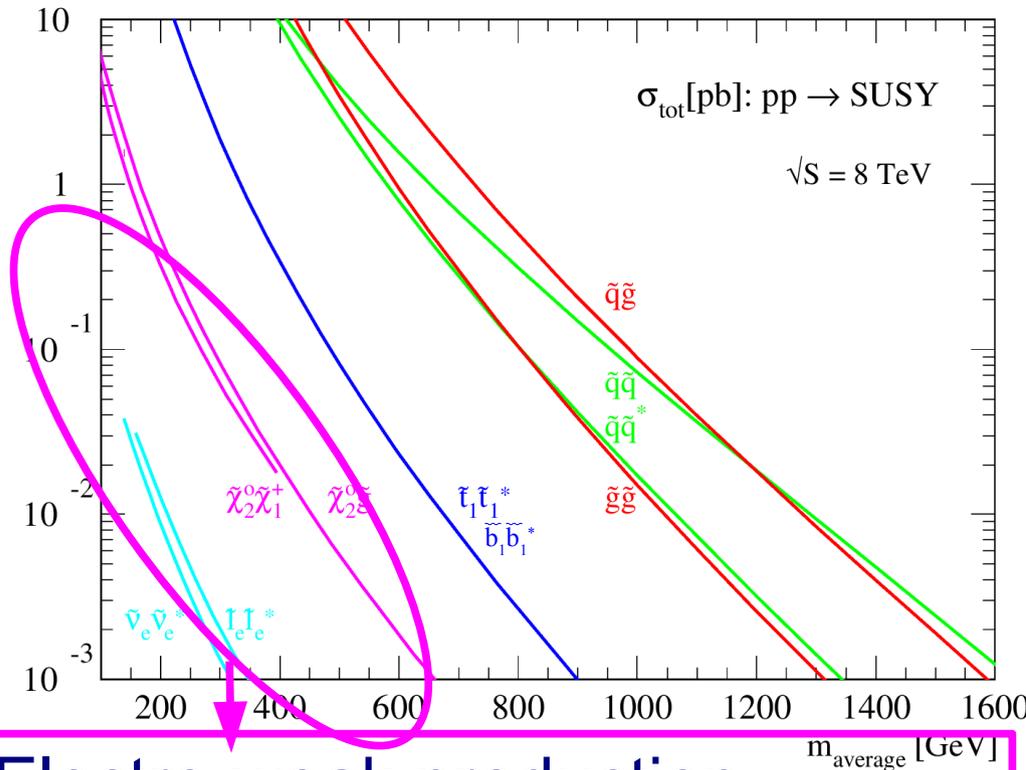
Current Stop Limit Summary

With the additional measurements at top starting to close some of the holes for light stops



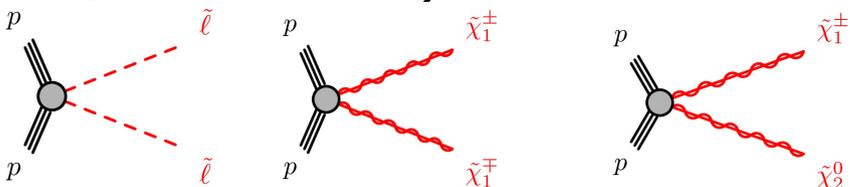
LHC SUSY Searches

At LHC can search for production for almost all SUSY particles, but with different sensitivity as production cross-sections vary



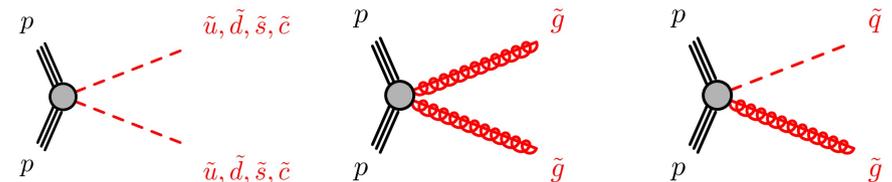
Electro-weak production

- Slepton, neutralino and charginos
- Detectable mostly through leptonic decays



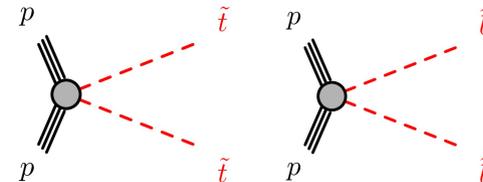
Production of squarks/gluinos:

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- Primary signature is jets and missing energy



Prod. of 3rd generation squarks:

- Lower cross section,
- Decays chains frequently has top and bottom quarks
- Of great interest to “Natural SUSY”

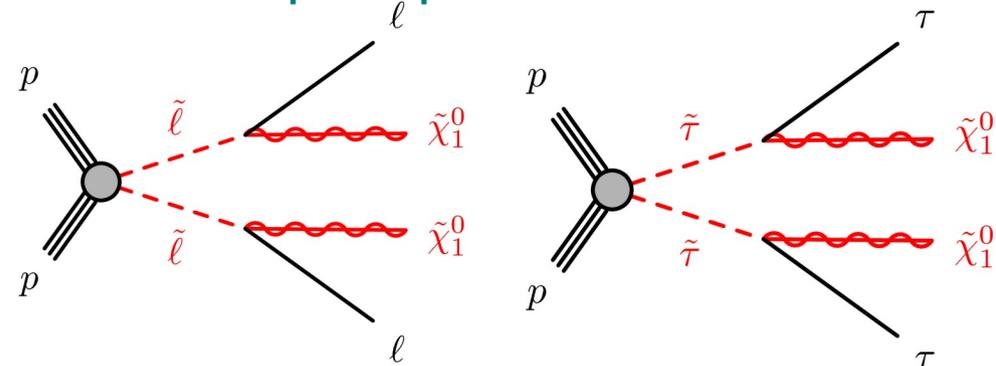


EW Production Search Program

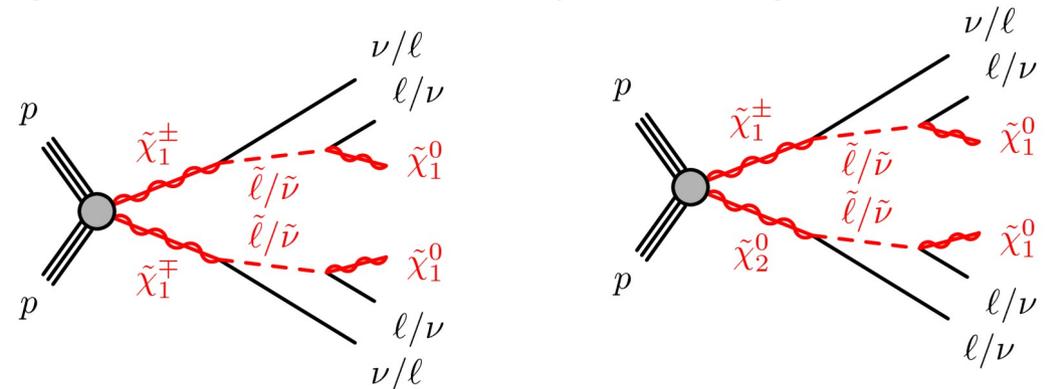
Comprehensive program

- Search for chargino, neutralino and slepton pair production
- Primarily using leptonic final states either through direct decays to leptons or through W/Z decays
- Dedicated searches with τ 's

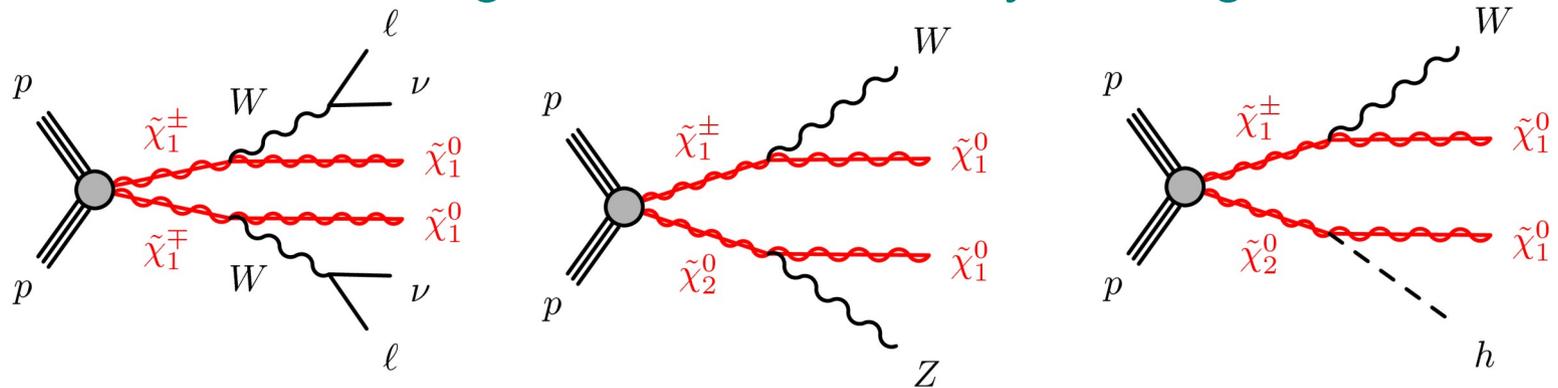
Direct slepton production



Chargino/neutralino decays through sleptons



Direct chargino/neutralino decays through W/Z/h's

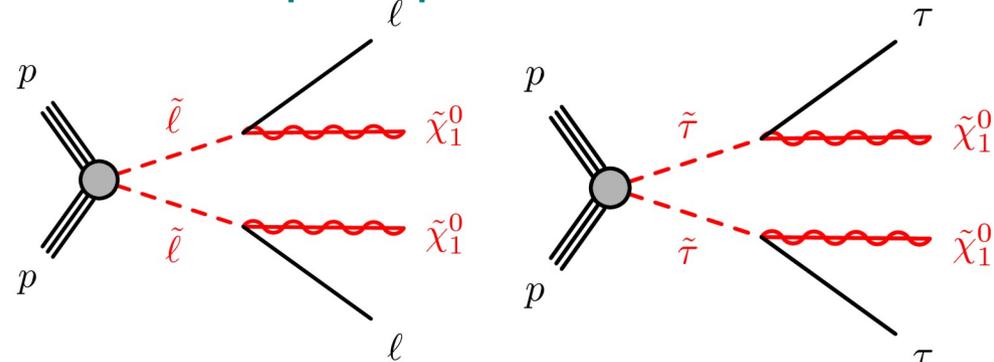


EW Production Search Program

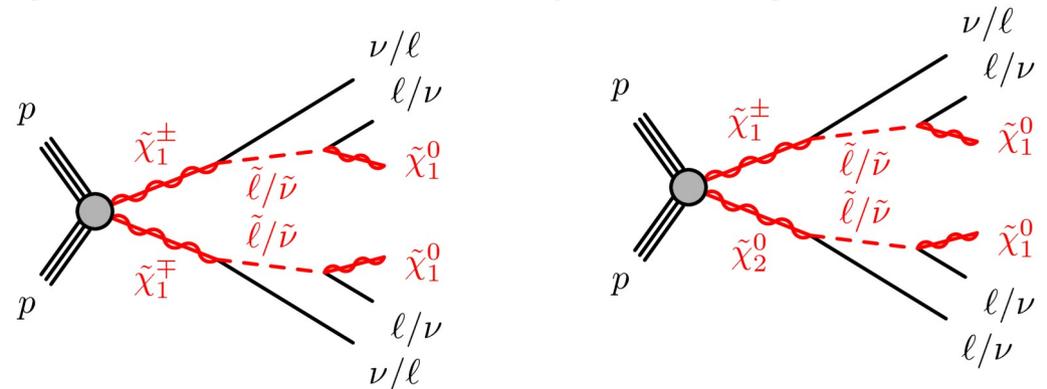
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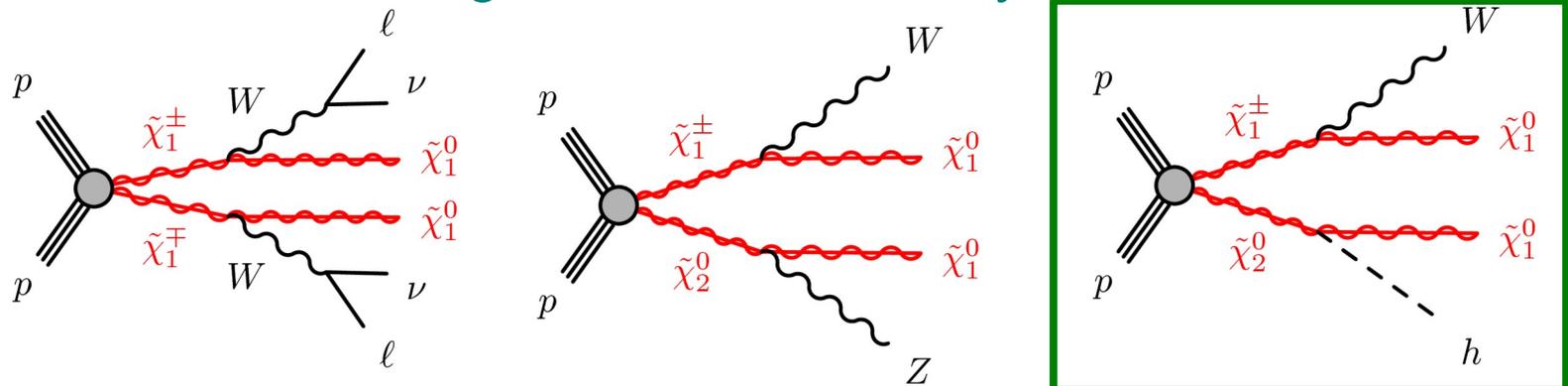
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Direct chargino/neutralino decays through W/Z/h's

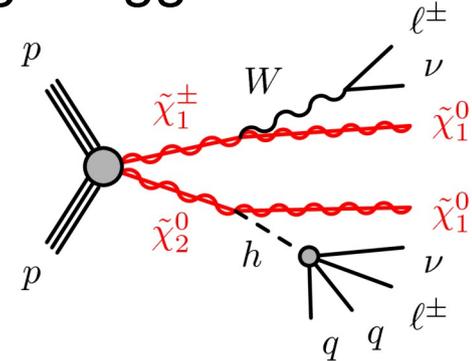
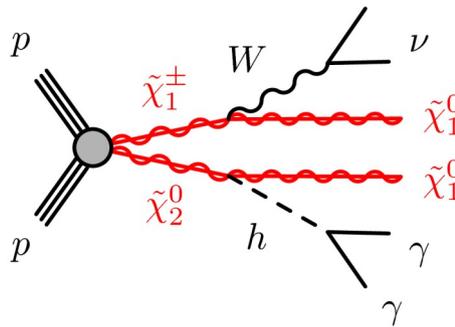
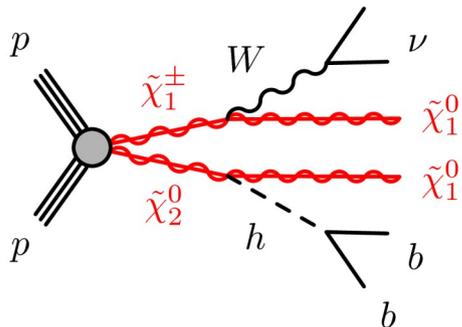


New

Search for EW Prod. With Higgs Decay

If $\tilde{\chi}_2^0$ is of wino-type and heavy enough, it will decay through Higgs emission

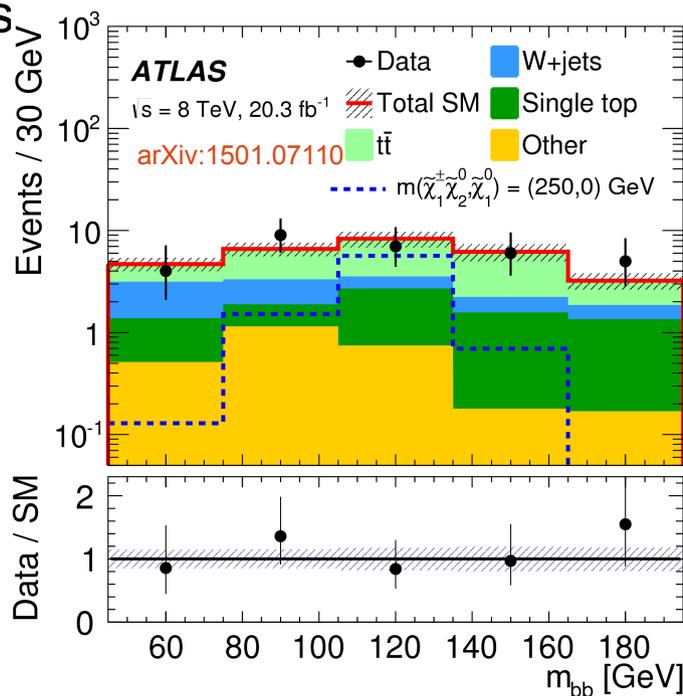
Searched for in three decay modes in association with $\tilde{\chi}_1^+$ decay



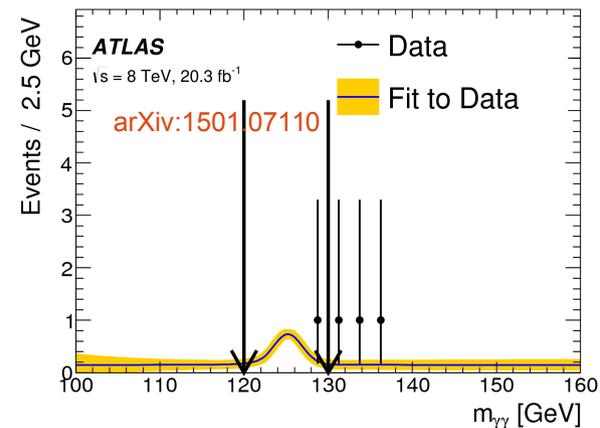
Search Strategy:

- 2 signal regions in $b\bar{b}$ using lepton- E_T^{miss} transverse mass and $m_{\text{CT},b\bar{b}}$ to suppress W +jet and $t\bar{t}$ backgrounds
- 2 signal regions in $\gamma\gamma$ using W + γ transverse mass and $\Delta\phi(W,\gamma)$ to suppress SM Higgs production
- 6 same-sign lepton signal regions split by flavor and jet multiplicity
- Combine all signal regions for maximal sensitivity

$m_{b\bar{b}}$ in $h \rightarrow b\bar{b}$ channel

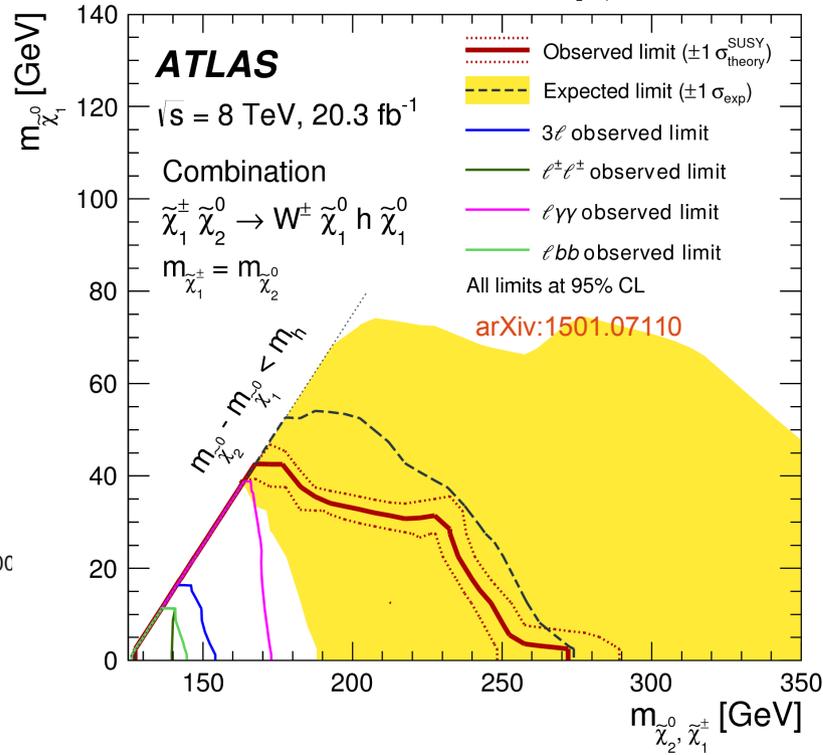
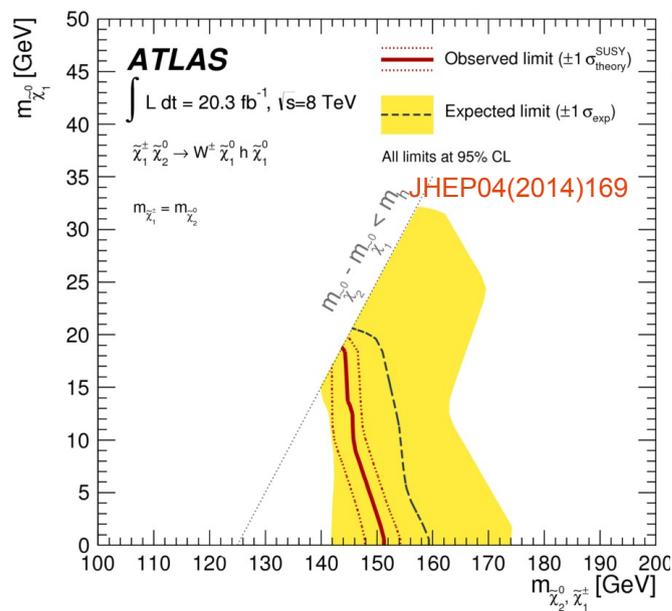
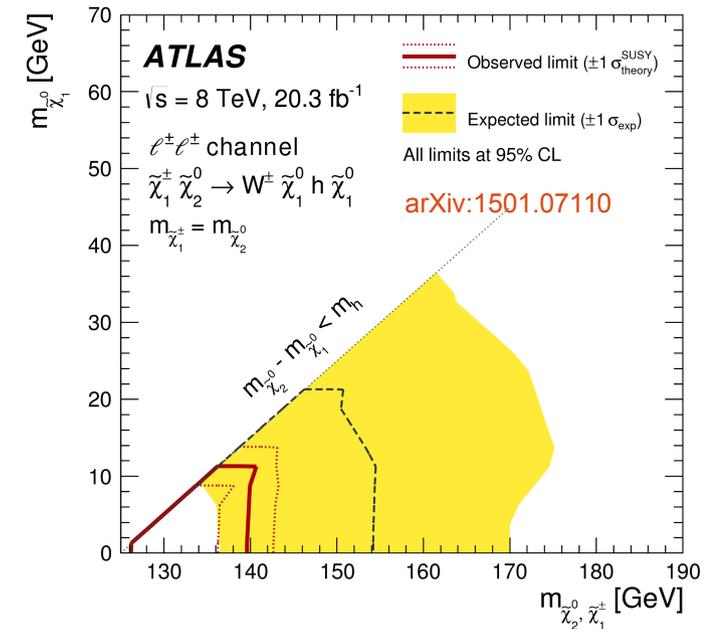
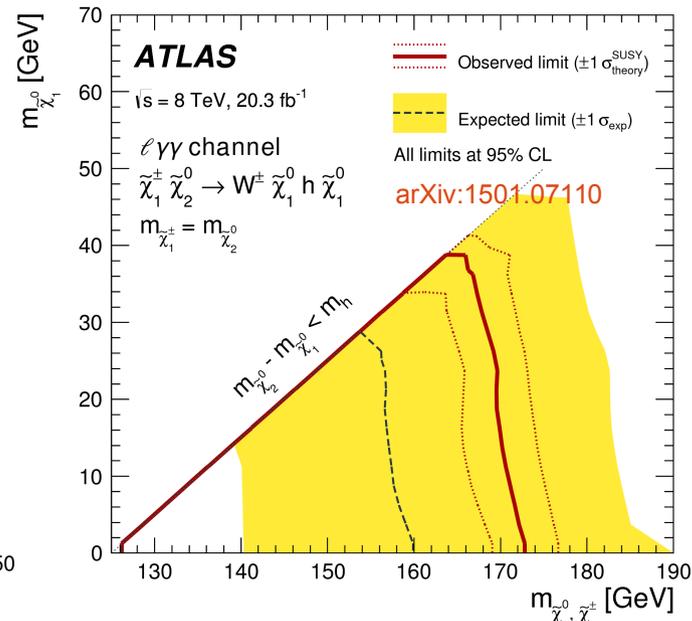
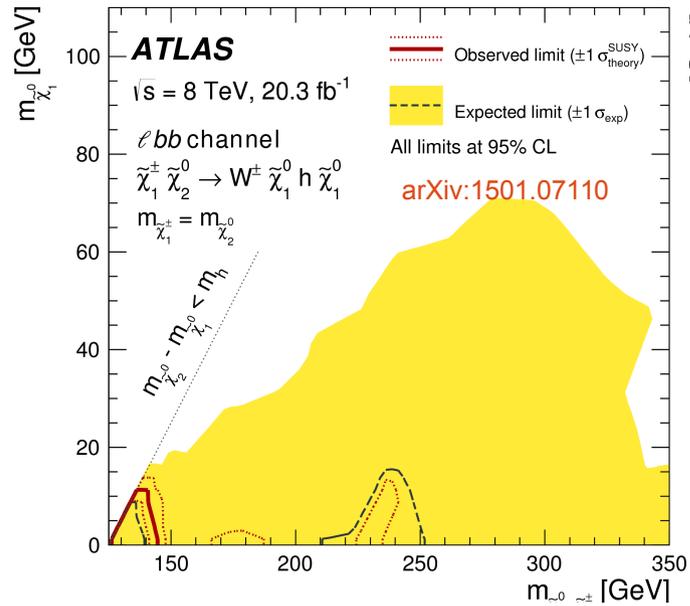


$m_{\gamma\gamma}$ in $h \rightarrow \gamma\gamma$ channel



New

Search for EW Prod. With Higgs Decay

Combined Result

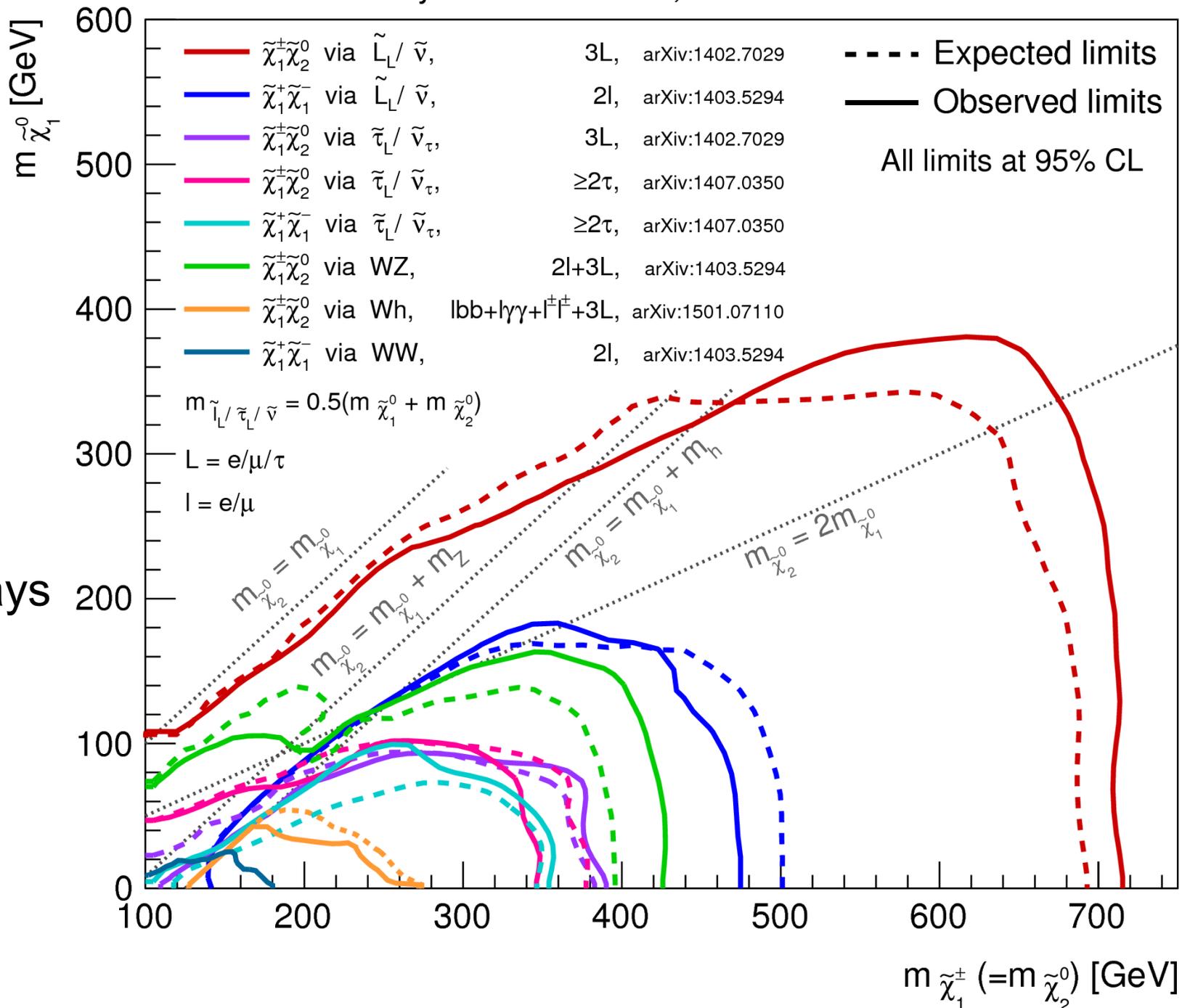
- Combine with older 3-lepton search
- Significant extension of search reach (up to 250 GeV)

Status of EW Production Searches

ATLAS Preliminary 20.3 fb⁻¹, √s=8 TeV Status: Feb 2015

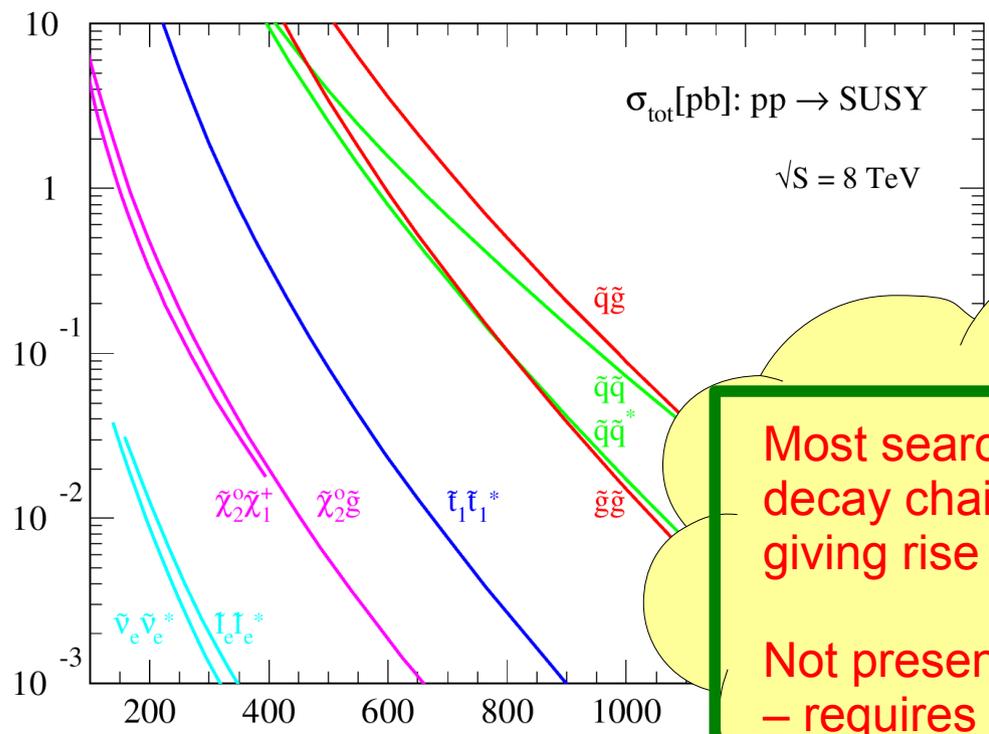
Rather strong limits in case of intermediate sleptons

Still relatively weak limits for direct chargino and neutralino decays



LHC SUSY Searches

At LHC can search for production for almost all SUSY particles, but with different sensitivity as production cross-sections vary

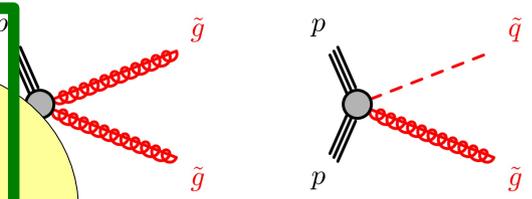


Production of squarks/gluinos:

- Strong coupling gives large sensitivity at the LHC
- Primary signature is jets and missing energy

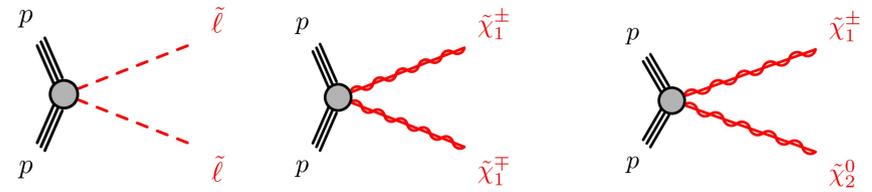
Most searches assume decay chain to stable LSP giving rise to missing energy

Not present in all SUSY models – requires dedicated searches to cover this



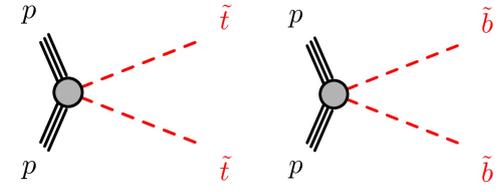
Electro-weak production

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- Detectable mostly through leptonic decays



Production squarks:

- Production cross-section, frequently has top and bottom quarks
- Of great interest to “Natural SUSY”



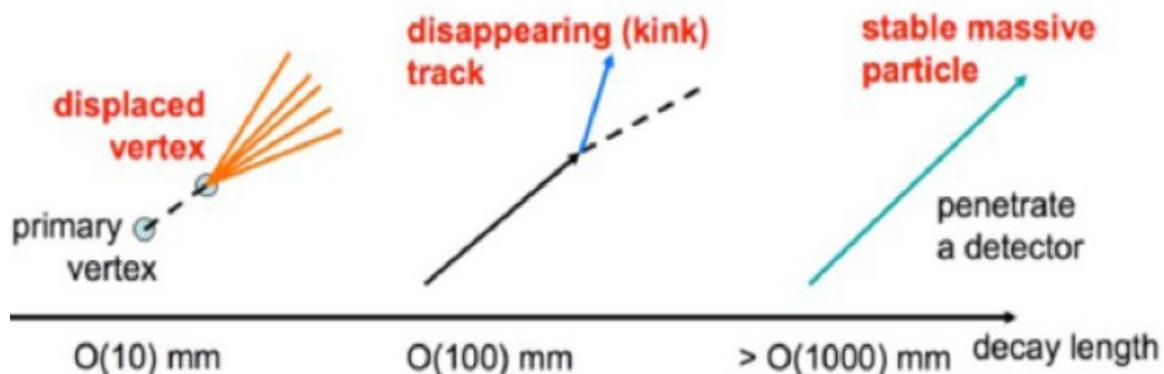
Long-lived SUSY Particles

Long-lived SUSY particles are predicted in wide variety of models: Hidden Sectors, RPV violating decays, Split SUSY, AMSB, GMSB,...

Possible Signatures

- Displaced vertices
- Disappearing tracks
- Non-pointing and delayed photons
- Stopped R-hadrons
- Stable massive particles

Observable signatures depend on lifetime



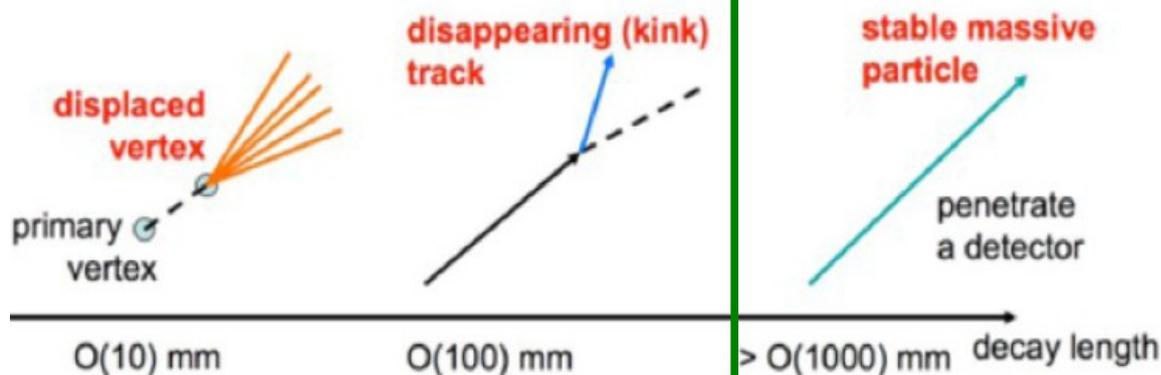
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Stable Massive Particle Search

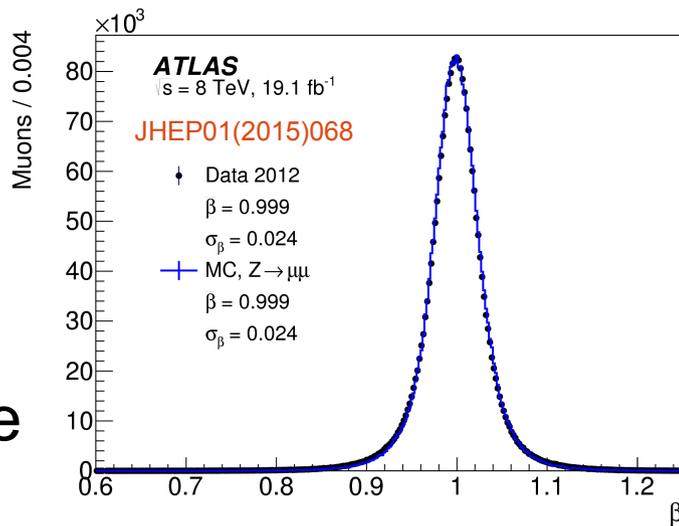
Massive charged particles leave several distinct signatures:

- Energy loss ($\beta\gamma$) measured by pixel detector
- Long time-of-flight measured by calorimeters and muon system
- With momentum measurement get mass estimate: $m = p/(\beta\gamma)$

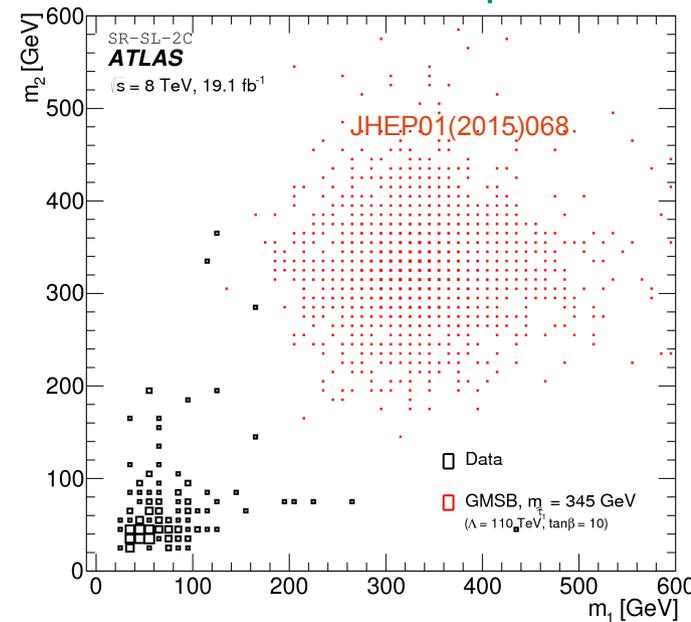
Search strategy:

- Reconstruct particles with $0.2 < \beta < 0.95$
- Determine mass from β measured in muon and calorimeters
- Confirm selection using pixel energy loss
- Either require two loose heavy particles or one tightly selected
- Optimize for different long-lived particle production scenarios

Combined β resolution



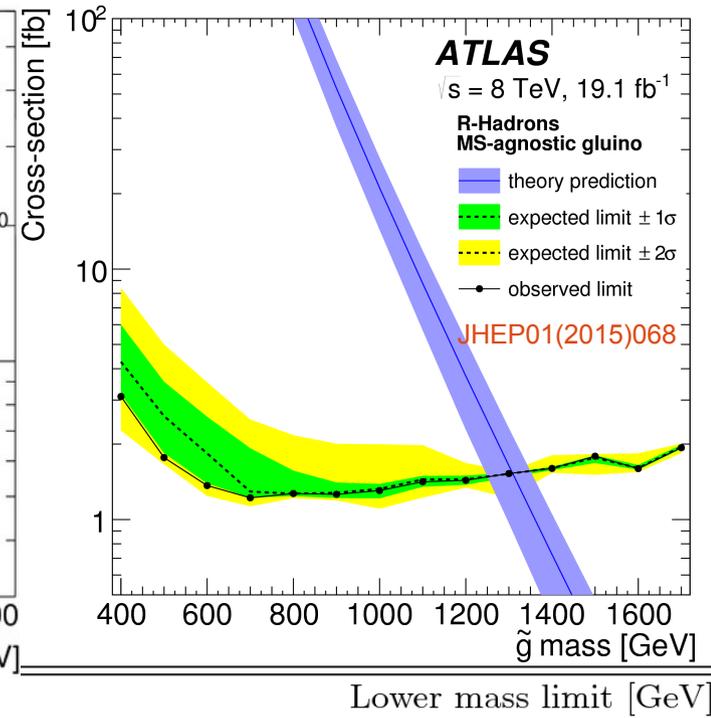
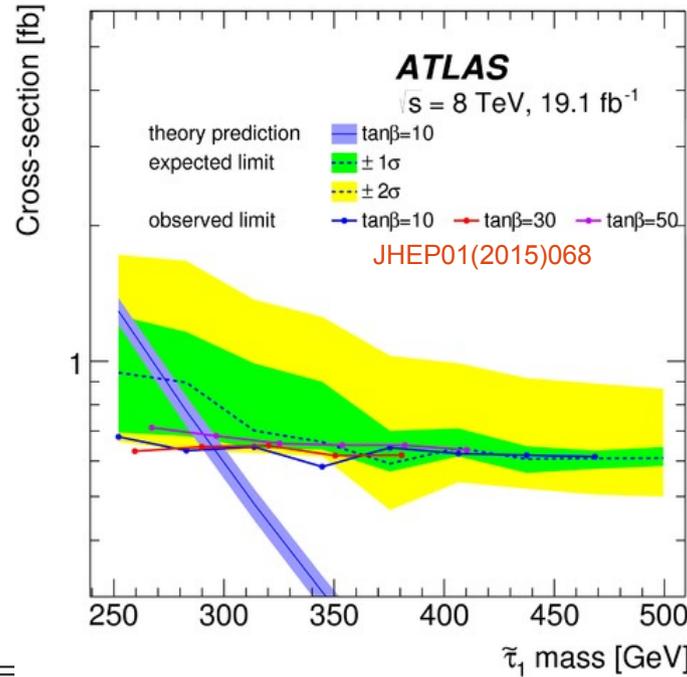
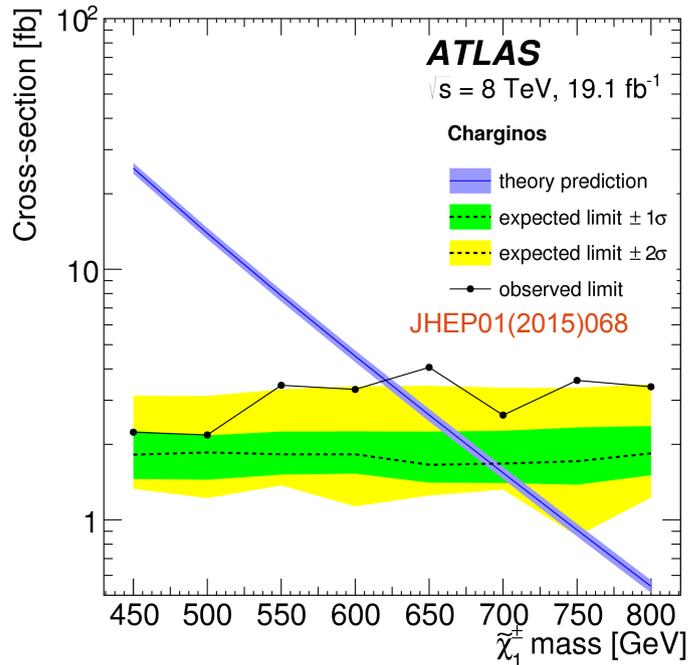
Reconstructed mass distributions for two particles



New

Stable Massive Particle Search

Interpreted in many different models with long-lived particles



Search

GMSB sleptons

- $\tan \beta = 10, 20, 30, 40, 50$
- direct $\tilde{\ell}$ production ($m_{\tilde{\ell}} - m_{\tilde{\tau}_1} = 2.7-93$ GeV)
- direct $\tilde{\tau}_1$ production
- $\tilde{\chi}_1^0 \tilde{\chi}_1^\pm$ decaying to stable $\tilde{\tau}_1$

440, 440, 430, 410, 385

377-335

289

537

LeptoSUSY

- \tilde{q}, \tilde{g}

1500, 1360

Charginos

- $\tilde{\chi}_1^\pm$

620

 R -hadrons

- $\tilde{g}, \tilde{b}, \tilde{t}$ (full-detector)
- $\tilde{g}, \tilde{b}, \tilde{t}$ (MS-agnostic)

1270, 845 and 900

1260, 835 and 870

Often limits are stronger than in case of prompt decays

Model	e, μ, τ, γ	Jets	E_T^{miss}	$\int \mathcal{L} dt [\text{fb}^{-1}]$	Mass limit	Reference		
Inclusive Searches	MSUGRA/CMSSM	0	2-6 jets	Yes	20.3	\tilde{q}, \tilde{g} 1.7 TeV	$m(\tilde{q})=m(\tilde{g})$	1405.7875
	$\tilde{q}\tilde{q}, \tilde{q} \rightarrow q\tilde{\chi}_1^0$	0	2-6 jets	Yes	20.3	\tilde{q} 850 GeV	$m(\tilde{\chi}_1^0)=0 \text{ GeV}, m(1^{\text{st}} \text{ gen. } \tilde{q})=m(2^{\text{nd}} \text{ gen. } \tilde{q})$	1405.7875
	$\tilde{q}\tilde{q}\gamma, \tilde{q} \rightarrow q\tilde{\chi}_1^0$ (compressed)	1 γ	0-1 jet	Yes	20.3	\tilde{q} 250 GeV	$m(\tilde{q})-m(\tilde{\chi}_1^0) = m(c)$	1411.1559
	$\tilde{g}\tilde{g}, \tilde{g} \rightarrow q\tilde{q}\tilde{\chi}_1^0$	0	2-6 jets	Yes	20.3	\tilde{g} 1.33 TeV	$m(\tilde{\chi}_1^0)=0 \text{ GeV}$	1405.7875
	$\tilde{g}\tilde{g}, \tilde{g} \rightarrow q\tilde{q}\tilde{\chi}_1^\pm \rightarrow q\tilde{q}W^\pm\tilde{\chi}_1^0$	1 e, μ	3-6 jets	Yes	20	\tilde{g} 1.2 TeV	$m(\tilde{\chi}_1^0) < 300 \text{ GeV}, m(\tilde{\chi}^\pm)=0.5(m(\tilde{\chi}_1^0)+m(\tilde{g}))$	1501.03555
	$\tilde{g}\tilde{g}, \tilde{g} \rightarrow q\tilde{q}(\ell\ell/\ell\nu/\nu\nu)\tilde{\chi}_1^0$	2 e, μ	0-3 jets	-	20	\tilde{g} 1.32 TeV	$m(\tilde{\chi}_1^0)=0 \text{ GeV}$	1501.03555
	GMSB ($\tilde{\ell}$ NLSP)	1-2 $\tau + 0-1 \ell$	0-2 jets	Yes	20.3	\tilde{g} 1.6 TeV	$\tan\beta > 20$	1407.0603
	GGM (bino NLSP)	2 γ	-	Yes	20.3	\tilde{g} 1.28 TeV	$m(\tilde{\chi}_1^0) > 50 \text{ GeV}$	ATLAS-CONF-2014-001
	GGM (wino NLSP)	1 $e, \mu + \gamma$	-	Yes	4.8	\tilde{g} 619 GeV	$m(\tilde{\chi}_1^0) > 50 \text{ GeV}$	ATLAS-CONF-2012-144
	GGM (higgsino-bino NLSP)	γ	1 b	Yes	4.8	\tilde{g} 900 GeV	$m(\tilde{\chi}_1^0) > 220 \text{ GeV}$	1211.1167
GGM (higgsino NLSP)	2 e, μ (Z)	0-3 jets	Yes	5.8	\tilde{g} 690 GeV	$m(\text{NLSP}) > 200 \text{ GeV}$	ATLAS-CONF-2012-152	
Gravitino LSP	0	mono-jet	Yes	20.3	$\tilde{r}^{1/2}$ scale 865 GeV	$m(\tilde{G}) > 1.8 \times 10^{-4} \text{ eV}, m(\tilde{g})=m(\tilde{q})=1.5 \text{ TeV}$	1502.01518	
3^{rd} gen. \tilde{g} med.	$\tilde{g} \rightarrow b\tilde{b}\tilde{\chi}_1^0$	0	3 b	Yes	20.1	\tilde{g} 1.25 TeV	$m(\tilde{\chi}_1^0) < 400 \text{ GeV}$	1407.0600
	$\tilde{g} \rightarrow t\tilde{t}\tilde{\chi}_1^0$	0	7-10 jets	Yes	20.3	\tilde{g} 1.1 TeV	$m(\tilde{\chi}_1^0) < 350 \text{ GeV}$	1308.1841
	$\tilde{g} \rightarrow t\tilde{b}\tilde{\chi}_1^0$	0-1 e, μ	3 b	Yes	20.1	\tilde{g} 1.34 TeV	$m(\tilde{\chi}_1^0) < 400 \text{ GeV}$	1407.0600
	$\tilde{g} \rightarrow b\tilde{\chi}_1^\pm$	0-1 e, μ	3 b	Yes	20.1	\tilde{g} 1.3 TeV	$m(\tilde{\chi}_1^0) < 300 \text{ GeV}$	1407.0600
3^{rd} gen. squarks direct production	$\tilde{b}_1\tilde{b}_1, \tilde{b}_1 \rightarrow b\tilde{\chi}_1^0$	0	2 b	Yes	20.1	\tilde{b}_1 100-620 GeV	$m(\tilde{\chi}_1^0) < 90 \text{ GeV}$	1308.2631
	$\tilde{b}_1\tilde{b}_1, \tilde{b}_1 \rightarrow t\tilde{\chi}_1^\pm$	2 e, μ (SS)	0-3 b	Yes	20.3	\tilde{b}_1 275-440 GeV	$m(\tilde{\chi}_1^\pm) = 2m(\tilde{\chi}_1^0)$	1404.2500
	$\tilde{t}_1\tilde{t}_1, \tilde{t}_1 \rightarrow b\tilde{\chi}_1^\pm$	1-2 e, μ	1-2 b	Yes	4.7	\tilde{t}_1 110-167 GeV, 230-460 GeV	$m(\tilde{\chi}_1^\pm) = 2m(\tilde{\chi}_1^0), m(\tilde{\chi}_1^\pm)=55 \text{ GeV}$	1209.2102, 1407.0583
	$\tilde{t}_1\tilde{t}_1, \tilde{t}_1 \rightarrow Wb\tilde{\chi}_1^0$ or $t\tilde{\chi}_1^0$	2 e, μ	0-2 jets	Yes	20.3	\tilde{t}_1 90-191 GeV, 215-530 GeV	$m(\tilde{\chi}_1^0)=1 \text{ GeV}$	1403.4853, 1412.4742
	$\tilde{t}_1\tilde{t}_1, \tilde{t}_1 \rightarrow t\tilde{\chi}_1^0$	0-1 e, μ	1-2 b	Yes	20	\tilde{t}_1 210-640 GeV	$m(\tilde{\chi}_1^0)=1 \text{ GeV}$	1407.0583, 1406.1122
	$\tilde{t}_1\tilde{t}_1, \tilde{t}_1 \rightarrow c\tilde{\chi}_1^0$	0	mono-jet/ c -tag	Yes	20.3	\tilde{t}_1 90-240 GeV	$m(\tilde{t}_1)-m(\tilde{\chi}_1^0) < 85 \text{ GeV}$	1407.0608
	$\tilde{t}_1\tilde{t}_1$ (natural GMSB)	2 e, μ (Z)	1 b	Yes	20.3	\tilde{t}_1 150-580 GeV	$m(\tilde{\chi}_1^0) > 150 \text{ GeV}$	1403.5222
$\tilde{t}_2\tilde{t}_2, \tilde{t}_2 \rightarrow \tilde{t}_1 + Z$	3 e, μ (Z)	1 b	Yes	20.3	\tilde{t}_2 290-600 GeV	$m(\tilde{\chi}_1^0) < 200 \text{ GeV}$	1403.5222	
EW direct	$\tilde{\ell}_{L,R}\tilde{\ell}_{L,R}, \tilde{\ell} \rightarrow \ell\tilde{\chi}_1^0$	2 e, μ	0	Yes	20.3	$\tilde{\ell}$ 90-325 GeV	$m(\tilde{\chi}_1^0)=0 \text{ GeV}$	1403.5294
	$\tilde{\chi}_1^\pm\tilde{\chi}_1^\pm, \tilde{\chi}_1^\pm \rightarrow \tilde{\ell}\nu(\ell\bar{\nu})$	2 e, μ	0	Yes	20.3	$\tilde{\chi}_1^\pm$ 140-465 GeV	$m(\tilde{\chi}_1^0)=0 \text{ GeV}, m(\tilde{\ell}, \bar{\nu})=0.5(m(\tilde{\chi}_1^\pm)+m(\tilde{\chi}_1^0))$	1403.5294
	$\tilde{\chi}_1^\pm\tilde{\chi}_1^\pm, \tilde{\chi}_1^\pm \rightarrow \tilde{\tau}\nu(\tau\bar{\nu})$	2 τ	-	Yes	20.3	$\tilde{\chi}_1^\pm$ 100-350 GeV	$m(\tilde{\chi}_1^0)=0 \text{ GeV}, m(\tilde{\tau}, \bar{\nu})=0.5(m(\tilde{\chi}_1^\pm)+m(\tilde{\chi}_1^0))$	1407.0350
	$\tilde{\chi}_1^\pm\tilde{\chi}_2^0 \rightarrow \tilde{\ell}_i\nu\tilde{\ell}_i\ell(\ell\nu\nu), \tilde{\ell}\tilde{\nu}\tilde{\ell}_i\ell(\ell\nu\nu)$	3 e, μ	0	Yes	20.3	$\tilde{\chi}_1^\pm, \tilde{\chi}_2^0$ 700 GeV	$m(\tilde{\chi}_1^\pm)=m(\tilde{\chi}_2^0), m(\tilde{\chi}_1^0)=0, m(\tilde{\ell}, \bar{\nu})=0.5(m(\tilde{\chi}_1^\pm)+m(\tilde{\chi}_1^0))$	1402.7029
	$\tilde{\chi}_1^\pm\tilde{\chi}_2^0 \rightarrow W\tilde{\chi}_1^0 Z\tilde{\chi}_1^0$	2-3 e, μ	0-2 jets	Yes	20.3	$\tilde{\chi}_1^\pm, \tilde{\chi}_2^0$ 420 GeV	$m(\tilde{\chi}_1^\pm)=m(\tilde{\chi}_2^0), m(\tilde{\chi}_1^0)=0$, sleptons decoupled	1403.5294, 1402.7029
	$\tilde{\chi}_1^\pm\tilde{\chi}_2^0 \rightarrow W\tilde{\chi}_1^0 h\tilde{\chi}_1^0, h \rightarrow b\tilde{b}/WW/\tau\tau/\gamma\gamma$	e, μ, γ	0-2 b	Yes	20.3	$\tilde{\chi}_1^\pm, \tilde{\chi}_2^0$ 250 GeV	$m(\tilde{\chi}_1^\pm)=m(\tilde{\chi}_2^0), m(\tilde{\chi}_1^0)=0$, sleptons decoupled	1501.07110
	$\tilde{\chi}_{2,3}^0, \tilde{\chi}_{2,3}^0 \rightarrow \tilde{\ell}_R\ell$	4 e, μ	0	Yes	20.3	$\tilde{\chi}_{2,3}^0$ 620 GeV	$m(\tilde{\chi}_2^0)=m(\tilde{\chi}_3^0), m(\tilde{\chi}_1^0)=0, m(\tilde{\ell}, \bar{\nu})=0.5(m(\tilde{\chi}_2^0)+m(\tilde{\chi}_1^0))$	1405.5086
Long-lived particles	Direct $\tilde{\chi}_1^\pm\tilde{\chi}_1^\mp$ prod., long-lived $\tilde{\chi}_1^\pm$	Disapp. trk	1 jet	Yes	20.3	$\tilde{\chi}_1^\pm$ 270 GeV	$m(\tilde{\chi}_1^\pm)-m(\tilde{\chi}_1^0)=160 \text{ MeV}, \tau(\tilde{\chi}_1^\pm)=0.2 \text{ ns}$	1310.3675
	Stable, stopped \tilde{g} R-hadron	0	1-5 jets	Yes	27.9	\tilde{g} 832 GeV	$m(\tilde{\chi}_1^0)=100 \text{ GeV}, 10 \mu\text{s} < \tau(\tilde{g}) < 1000 \text{ s}$	1310.6584
	Stable \tilde{g} R-hadron	trk	-	-	19.1	\tilde{g} 1.27 TeV		1411.6795
	GMSB, stable $\tilde{\tau}, \tilde{\chi}_1^0 \rightarrow \tilde{\tau}(\tilde{\mu}, \tilde{\nu}) + \tau(e, \mu)$	1-2 μ	-	-	19.1	$\tilde{\chi}_1^0$ 537 GeV	$10 < \tan\beta < 50$	1411.6795
	GMSB, $\tilde{\chi}_1^0 \rightarrow \gamma\tilde{G}$, long-lived $\tilde{\chi}_1^0$	2 γ	-	Yes	20.3	$\tilde{\chi}_1^0$ 435 GeV	$2 < \tau(\tilde{\chi}_1^0) < 3 \text{ ns}$, SPS8 model	1409.5542
	$\tilde{q}\tilde{q}, \tilde{\chi}_1^0 \rightarrow q\tilde{q}$ (RPV)	1 μ , displ. vtx	-	-	20.3	\tilde{q} 1.0 TeV	$1.5 < c\tau < 156 \text{ mm}, \text{BR}(\mu)=1, m(\tilde{\chi}_1^0)=108 \text{ GeV}$	ATLAS-CONF-2013-092
RPV	LFV $pp \rightarrow \tilde{\nu}_\tau + X, \tilde{\nu}_\tau \rightarrow e + \mu$	2 e, μ	-	-	4.6	$\tilde{\nu}_\tau$ 1.61 TeV	$\lambda'_{311}=0.10, \lambda_{132}=0.05$	1212.1272
	LFV $pp \rightarrow \tilde{\nu}_\tau + X, \tilde{\nu}_\tau \rightarrow e(\mu) + \tau$	1 $e, \mu + \tau$	-	-	4.6	$\tilde{\nu}_\tau$ 1.1 TeV	$\lambda'_{311}=0.10, \lambda_{1(2)33}=0.05$	1212.1272
	Bilinear RPV CMSSM	2 e, μ (SS)	0-3 b	Yes	20.3	\tilde{q}, \tilde{g} 1.35 TeV	$m(\tilde{q})=m(\tilde{g}), c\tau_{LS} < 1 \text{ mm}$	1404.2500
	$\tilde{\chi}_1^\pm\tilde{\chi}_1^\mp, \tilde{\chi}_1^\pm \rightarrow W\tilde{\chi}_1^0, \tilde{\chi}_1^0 \rightarrow ee\tilde{\nu}_\mu, e\mu\tilde{\nu}_e$	4 e, μ	-	Yes	20.3	$\tilde{\chi}_1^\pm$ 750 GeV	$m(\tilde{\chi}_1^0) > 0.2 \times m(\tilde{\chi}_1^\pm), \lambda_{121} \neq 0$	1405.5086
	$\tilde{\chi}_1^\pm\tilde{\chi}_1^\mp, \tilde{\chi}_1^\pm \rightarrow W\tilde{\chi}_1^0, \tilde{\chi}_1^0 \rightarrow \tau\tau\tilde{\nu}_e, e\tau\tilde{\nu}_\tau$	3 $e, \mu + \tau$	-	Yes	20.3	$\tilde{\chi}_1^\pm$ 450 GeV	$m(\tilde{\chi}_1^0) > 0.2 \times m(\tilde{\chi}_1^\pm), \lambda_{133} \neq 0$	1405.5086
	$\tilde{g} \rightarrow qq\tilde{q}$	0	6-7 jets	-	20.3	\tilde{g} 916 GeV	$\text{BR}(t)=\text{BR}(b)=\text{BR}(c)=0\%$	ATLAS-CONF-2013-091
$\tilde{g} \rightarrow \tilde{t}_1 t, \tilde{t}_1 \rightarrow bs$	2 e, μ (SS)	0-3 b	Yes	20.3	\tilde{g} 850 GeV		1404.250	
Other	Scalar charm, $\tilde{c} \rightarrow c\tilde{\chi}_1^0$	0	2 c	Yes	20.3	\tilde{c} 490 GeV	$m(\tilde{\chi}_1^0) < 200 \text{ GeV}$	1501.01325

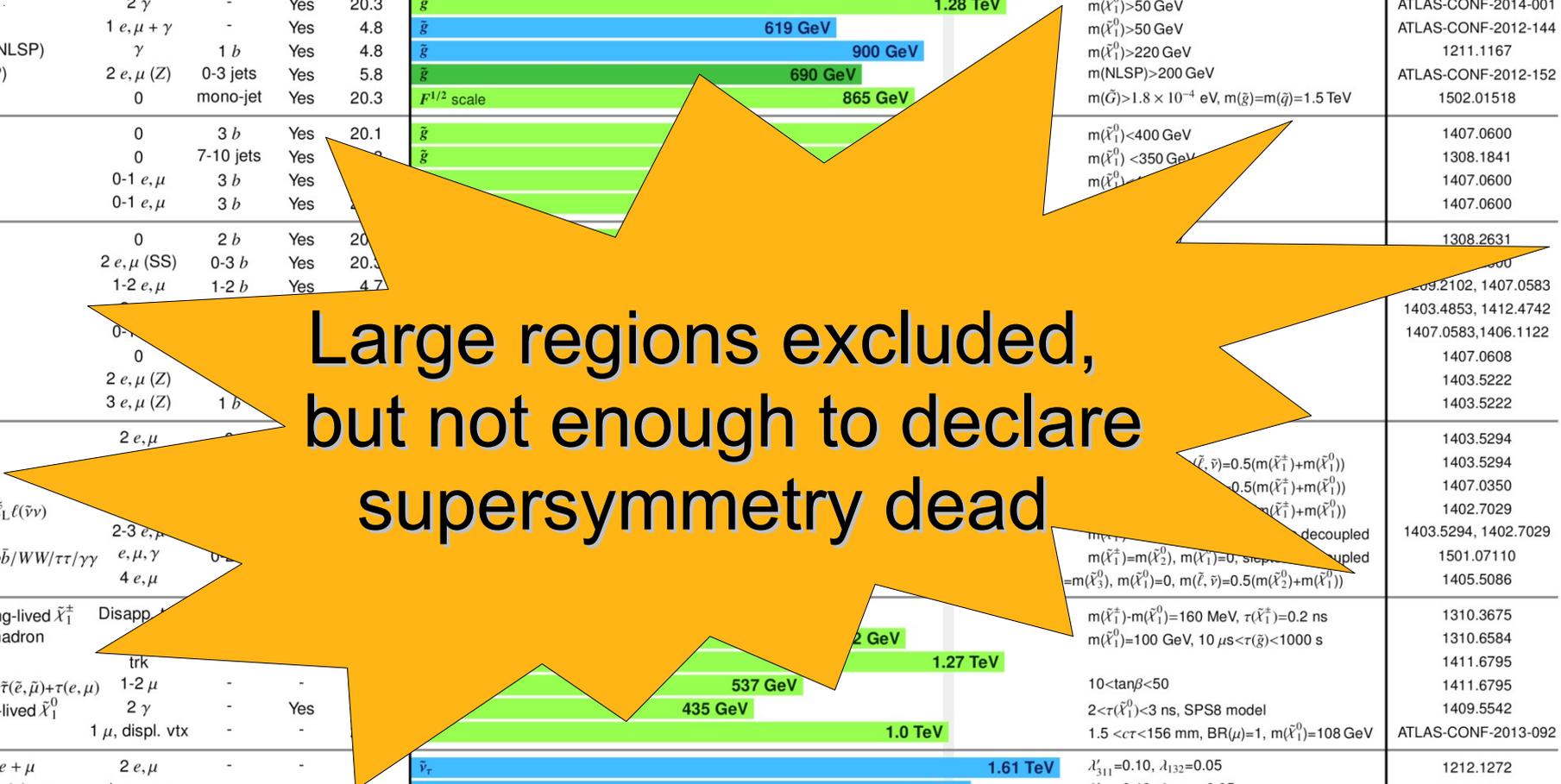
$\sqrt{s} = 7 \text{ TeV}$ full data $\sqrt{s} = 8 \text{ TeV}$ partial data $\sqrt{s} = 8 \text{ TeV}$ full data

10⁻¹ 1 Mass scale [TeV]

*Only a selection of the available mass limits on new states or phenomena is shown. All limits quoted are observed minus 1 σ theoretical signal cross section uncertainty.

Model e, μ, τ, γ Jets E_T^{miss} $\int \mathcal{L} dt [\text{fb}^{-1}]$ Mass limit Reference

Model	e, μ, τ, γ	Jets	E_T^{miss}	$\int \mathcal{L} dt [\text{fb}^{-1}]$	Mass limit	Reference		
Inclusive Searches	MSUGRA/CMSSM	0	2-6 jets	Yes	20.3	\tilde{q}, \tilde{g} 1.7 TeV	$m(\tilde{q})=m(\tilde{g})$	1405.7875
	$\tilde{q}\tilde{q}, \tilde{q} \rightarrow q\tilde{\chi}_1^0$	0	2-6 jets	Yes	20.3	\tilde{q} 850 GeV	$m(\tilde{\chi}_1^0)=0$ GeV, $m(1^{st} \text{ gen. } \tilde{q})=m(2^{nd} \text{ gen. } \tilde{q})$	1405.7875
	$\tilde{q}\tilde{q}\gamma, \tilde{q} \rightarrow q\tilde{\chi}_1^0$ (compressed)	1 γ	0-1 jet	Yes	20.3	\tilde{q} 250 GeV	$m(\tilde{q})-m(\tilde{\chi}_1^0) = m(c)$	1411.1559
	$\tilde{g}\tilde{g}, \tilde{g} \rightarrow q\tilde{q}\tilde{\chi}_1^0$	0	2-6 jets	Yes	20.3	\tilde{g} 1.33 TeV	$m(\tilde{\chi}_1^0)=0$ GeV	1405.7875
	$\tilde{g}\tilde{g}, \tilde{g} \rightarrow q\tilde{q}\tilde{\chi}_1^\pm \rightarrow q\tilde{q}W^\pm\tilde{\chi}_1^0$	1 e, μ	3-6 jets	Yes	20	\tilde{g} 1.2 TeV	$m(\tilde{\chi}_1^0) < 300$ GeV, $m(\tilde{\chi}^\pm)=0.5(m(\tilde{\chi}_1^0)+m(\tilde{g}))$	1501.03555
	$\tilde{g}\tilde{g}, \tilde{g} \rightarrow q\tilde{q}(\ell\ell/\ell\nu/\nu\nu)\tilde{\chi}_1^0$	2 e, μ	0-3 jets	-	20	\tilde{g} 1.32 TeV	$m(\tilde{\chi}_1^0)=0$ GeV	1501.03555
	GMSB ($\tilde{\ell}$ NLSP)	1-2 τ + 0-1 ℓ	0-2 jets	Yes	20.3	\tilde{g} 1.6 TeV	$\tan\beta > 20$	1407.0603
	GGM (bino NLSP)	2 γ	-	Yes	20.3	\tilde{g} 1.28 TeV	$m(\tilde{\chi}_1^0) > 50$ GeV	ATLAS-CONF-2014-001
	GGM (wino NLSP)	1 $e, \mu + \gamma$	-	Yes	4.8	\tilde{g} 619 GeV	$m(\tilde{\chi}_1^0) > 50$ GeV	ATLAS-CONF-2012-144
	GGM (higgsino-bino NLSP)	γ	1 b	Yes	4.8	\tilde{g} 900 GeV	$m(\tilde{\chi}_1^0) > 220$ GeV	1211.1167
GGM (higgsino NLSP)	2 e, μ (Z)	0-3 jets	Yes	5.8	\tilde{g} 690 GeV	$m(\text{NLSP}) > 200$ GeV	ATLAS-CONF-2012-152	
Gravitino LSP	0	mono-jet	Yes	20.3	$R^{1/2}$ scale 865 GeV	$m(\tilde{G}) > 1.8 \times 10^{-4}$ eV, $m(\tilde{g})=m(\tilde{q})=1.5$ TeV	1502.01518	
3 rd gen. \tilde{g} med.	$\tilde{g} \rightarrow b\tilde{b}\tilde{\chi}_1^0$	0	3 b	Yes	20.1	\tilde{g}	$m(\tilde{\chi}_1^0) < 400$ GeV	1407.0600
	$\tilde{g} \rightarrow t\tilde{t}\tilde{\chi}_1^0$	0	7-10 jets	Yes	20.1	\tilde{g}	$m(\tilde{\chi}_1^0) < 350$ GeV	1308.1841
	$\tilde{g} \rightarrow u\tilde{u}\tilde{\chi}_1^0$	0-1 e, μ	3 b	Yes	20.1	\tilde{g}	$m(\tilde{\chi}_1^0) < 350$ GeV	1407.0600
	$\tilde{g} \rightarrow b\tilde{d}\tilde{\chi}_1^+$	0-1 e, μ	3 b	Yes	20.1	\tilde{g}	$m(\tilde{\chi}_1^0) < 350$ GeV	1407.0600
3 rd gen. squarks direct production	$\tilde{b}_1\tilde{b}_1, \tilde{b}_1 \rightarrow b\tilde{\chi}_1^0$	0	2 b	Yes	20.1	\tilde{b}_1	$m(\tilde{\chi}_1^0) < 350$ GeV	1308.2631
	$\tilde{b}_1\tilde{b}_1, \tilde{b}_1 \rightarrow t\tilde{\chi}_1^\pm$	2 e, μ (SS)	0-3 b	Yes	20.3	\tilde{b}_1	$m(\tilde{\chi}_1^0) < 350$ GeV	1407.0600
	$\tilde{t}_1\tilde{t}_1, \tilde{t}_1 \rightarrow b\tilde{\chi}_1^\pm$	1-2 e, μ	1-2 b	Yes	4.7	\tilde{t}_1	$m(\tilde{\chi}_1^0) < 350$ GeV	1409.2102, 1407.0583
	$\tilde{t}_1\tilde{t}_1, \tilde{t}_1 \rightarrow Wb\tilde{\chi}_1^0$ or $t\tilde{\chi}_1^0$	0-1 e, μ	1-2 b	Yes	4.7	\tilde{t}_1	$m(\tilde{\chi}_1^0) < 350$ GeV	1403.4853, 1412.4742
	$\tilde{t}_1\tilde{t}_1, \tilde{t}_1 \rightarrow t\tilde{\chi}_1^0$	0-1 e, μ	1-2 b	Yes	4.7	\tilde{t}_1	$m(\tilde{\chi}_1^0) < 350$ GeV	1407.0583, 1406.1122
	$\tilde{t}_1\tilde{t}_1, \tilde{t}_1 \rightarrow c\tilde{\chi}_1^0$	0	1-2 b	Yes	4.7	\tilde{t}_1	$m(\tilde{\chi}_1^0) < 350$ GeV	1407.0608
EW direct	$\tilde{t}_1\tilde{t}_1$ (natural GMSB)	2 e, μ (Z)	1 b	Yes	4.7	\tilde{t}_1	$m(\tilde{\chi}_1^0) < 350$ GeV	1403.5222
	$\tilde{t}_2\tilde{t}_2, \tilde{t}_2 \rightarrow \tilde{t}_1 + Z$	3 e, μ (Z)	1 b	Yes	4.7	\tilde{t}_2	$m(\tilde{\chi}_1^0) < 350$ GeV	1403.5222
	$\tilde{\ell}_{L,R}\tilde{\ell}_{L,R}, \tilde{\ell} \rightarrow \ell\tilde{\chi}_1^0$	2 e, μ	0	Yes	20.3	$\tilde{\ell}$	$m(\tilde{\chi}_1^0) < 350$ GeV	1403.5294
Long-lived particles	Direct $\tilde{\chi}_1^+\tilde{\chi}_1^-$ prod., long-lived $\tilde{\chi}_1^\pm$	Disapp.	-	-	20.3	$\tilde{\chi}_1^\pm$ 2 GeV	$m(\tilde{\chi}_1^\pm) > 160$ MeV, $\tau(\tilde{\chi}_1^\pm) > 0.2$ ns	1310.3675
	Stable, stopped \tilde{g} R-hadron	trk	-	-	20.3	\tilde{g} 1.27 TeV	$m(\tilde{\chi}_1^0)=100$ GeV, $10 \mu\text{s} < \tau(\tilde{g}) < 1000$ s	1310.6584
	Stable \tilde{g} R-hadron	trk	-	-	20.3	\tilde{g} 1.27 TeV	$m(\tilde{\chi}_1^0)=100$ GeV, $10 \mu\text{s} < \tau(\tilde{g}) < 1000$ s	1411.6795
	GMSB, stable $\tilde{\tau}, \tilde{\chi}_1^0 \rightarrow \tilde{\tau}(\tilde{e}, \tilde{\mu}) + \tau(e, \mu)$	1-2 μ	-	-	20.3	$\tilde{\tau}$ 537 GeV	$10 < \tan\beta < 50$	1411.6795
	GMSB, $\tilde{\chi}_1^0 \rightarrow \gamma\tilde{G}$, long-lived $\tilde{\chi}_1^0$	2 γ	-	Yes	20.3	$\tilde{\chi}_1^0$ 435 GeV	$2 < \tau(\tilde{\chi}_1^0) < 3$ ns, SPS8 model	1409.5542
	$\tilde{q}\tilde{q}, \tilde{\chi}_1^0 \rightarrow q\tilde{q}\mu$ (RPV)	1 μ , displ. vtx	-	-	20.3	\tilde{q} 1.0 TeV	$1.5 < c\tau < 156$ mm, $\text{BR}(\mu)=1$, $m(\tilde{\chi}_1^0)=108$ GeV	ATLAS-CONF-2013-092
	LFV $pp \rightarrow \tilde{\nu}_\tau + X, \tilde{\nu}_\tau \rightarrow e + \mu$	2 e, μ	-	-	4.6	$\tilde{\nu}_\tau$ 1.61 TeV	$\lambda'_{311}=0.10, \lambda_{132}=0.05$	1212.1272
	LFV $pp \rightarrow \tilde{\nu}_\tau + X, \tilde{\nu}_\tau \rightarrow e(\mu) + \tau$	1 $e, \mu + \tau$	-	-	4.6	$\tilde{\nu}_\tau$ 1.1 TeV	$\lambda'_{311}=0.10, \lambda_{1(2)333}=0.05$	1212.1272
	Bilinear RPV CMSSM	2 e, μ (SS)	0-3 b	Yes	20.3	\tilde{q}, \tilde{g} 1.35 TeV	$m(\tilde{q})=m(\tilde{g}), c\tau_{LS} < 1$ mm	1404.2500
	$\tilde{\chi}_1^+\tilde{\chi}_1^-, \tilde{\chi}_1^+ \rightarrow W\tilde{\chi}_1^0, \tilde{\chi}_1^0 \rightarrow ee\tilde{\nu}_\mu, e\mu\tilde{\nu}_e$	4 e, μ	-	Yes	20.3	$\tilde{\chi}_1^\pm$ 750 GeV	$m(\tilde{\chi}_1^0) > 0.2 \times m(\tilde{\chi}_1^\pm), \lambda_{121} \neq 0$	1405.5086
$\tilde{\chi}_1^+\tilde{\chi}_1^-, \tilde{\chi}_1^+ \rightarrow W\tilde{\chi}_1^0, \tilde{\chi}_1^0 \rightarrow \tau\tau\tilde{\nu}_e, e\tau\tilde{\nu}_\tau$	3 $e, \mu + \tau$	-	Yes	20.3	$\tilde{\chi}_1^\pm$ 450 GeV	$m(\tilde{\chi}_1^0) > 0.2 \times m(\tilde{\chi}_1^\pm), \lambda_{133} \neq 0$	1405.5086	
RPV	$\tilde{g} \rightarrow qq\tilde{q}$	0	6-7 jets	-	20.3	\tilde{g} 916 GeV	$\text{BR}(\tau)=\text{BR}(b)=\text{BR}(c)=0\%$	ATLAS-CONF-2013-091
	$\tilde{g} \rightarrow \tilde{t}_1 t, \tilde{t}_1 \rightarrow bs$	2 e, μ (SS)	0-3 b	Yes	20.3	\tilde{g} 850 GeV		1404.250
Other	Scalar charm, $\tilde{c} \rightarrow c\tilde{\chi}_1^0$	0	2 c	Yes	20.3	\tilde{c} 490 GeV	$m(\tilde{\chi}_1^0) < 200$ GeV	1501.01325



$\sqrt{s} = 7$ TeV full data $\sqrt{s} = 8$ TeV partial data $\sqrt{s} = 8$ TeV full data

10⁻¹ 1 Mass scale [TeV]

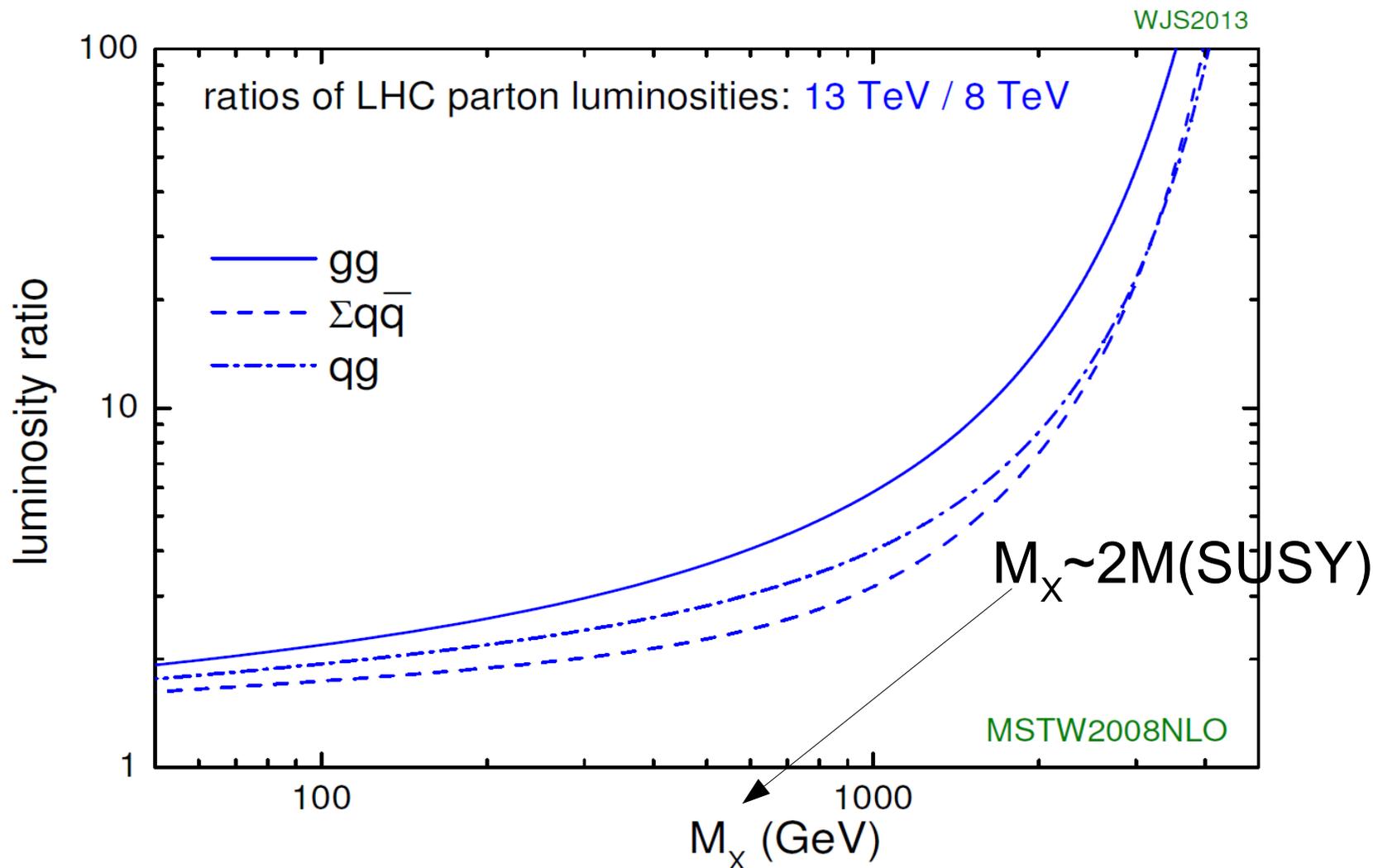
*Only a selection of the available mass limits on new states or phenomena is shown. All limits quoted are observed minus 1 σ theoretical signal cross section uncertainty.

Outlook for Supersymmetry Searches at the LHC

Future LHC Running

Two major improvements coming at the LHC

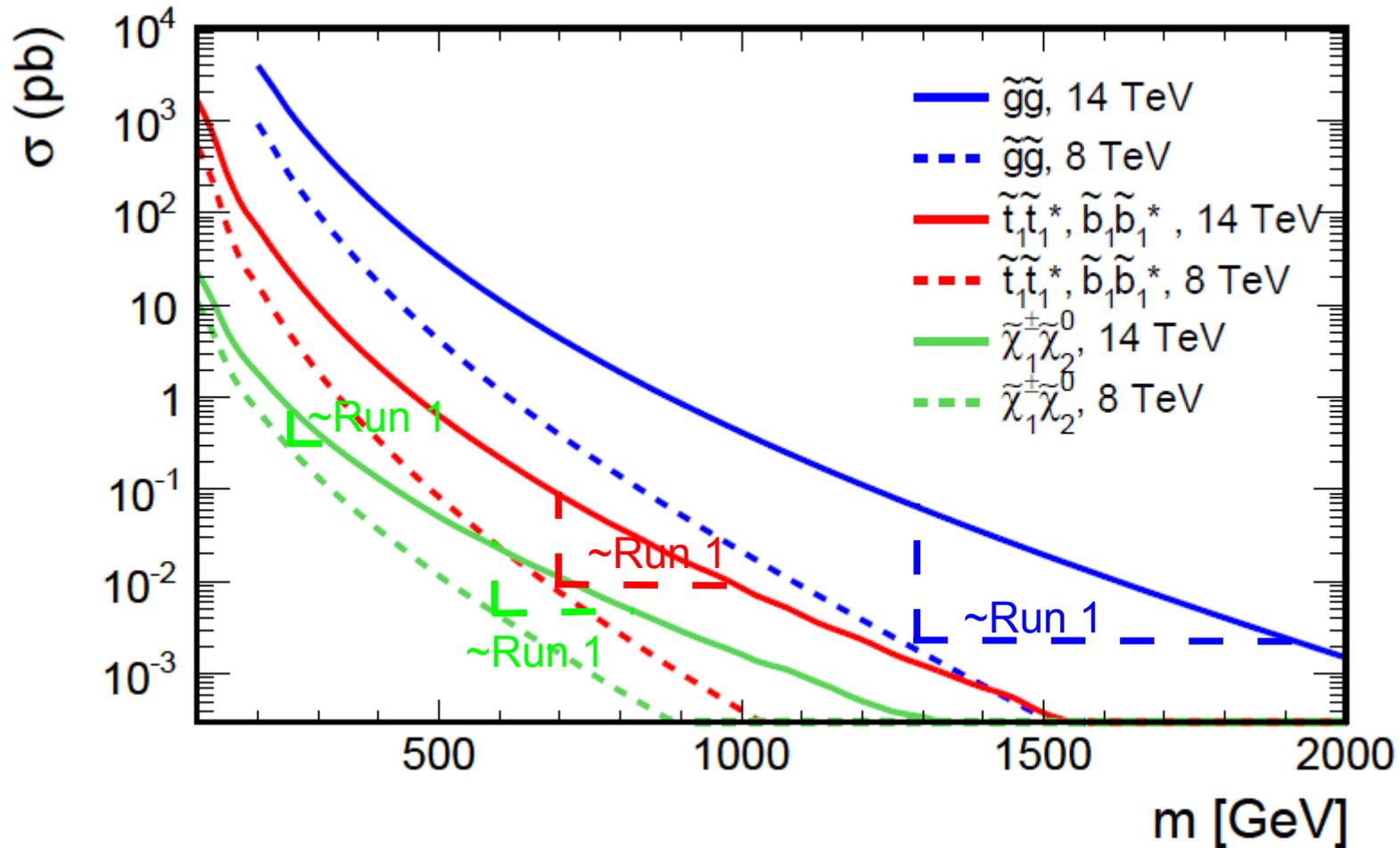
- Increase of collision energy from 8 TeV to 13 TeV this year
Will eventually increase towards 14 TeV
- Increase in integrated luminosity by O(100) over next ~20 years



Collision Energy Increase

Immediate impact of higher collision energy

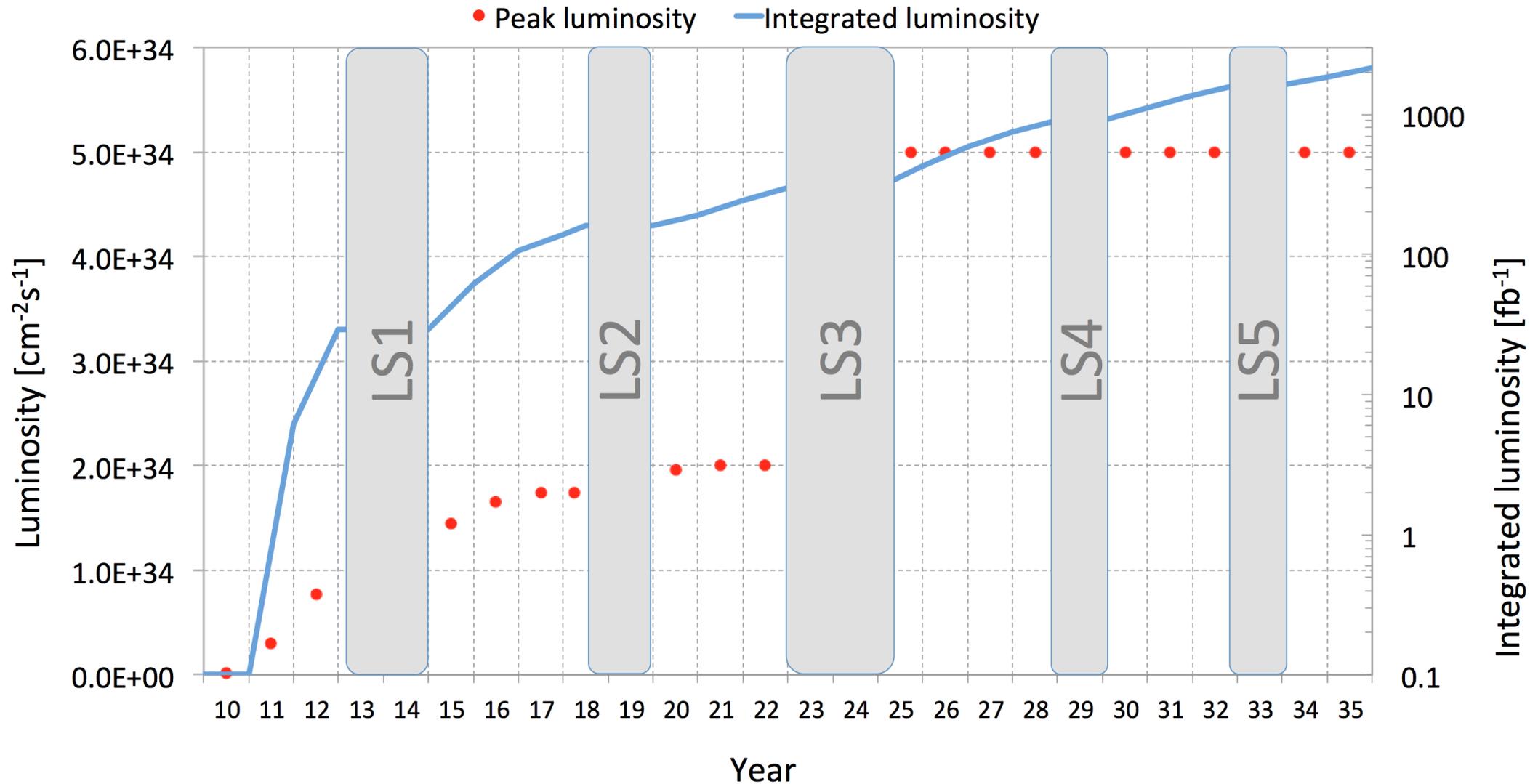
- O(50) increase in cross section for heavy gluinos
- O(10) increase in cross section for heavy stop
- Around factor two increase for EW production cross sections
- SM physics backgrounds also increase by factor 2-4



Luminosity Evolution

Expectations for the LHC

- $\sim 10 \text{ fb}^{-1}$ at 13 TeV by end of this year
- $\sim 300 \text{ fb}^{-1}$ at 13-14 TeV by 2022
- $\sim 3000 \text{ fb}^{-1}$ by 2035



Projections for SUSY Searches

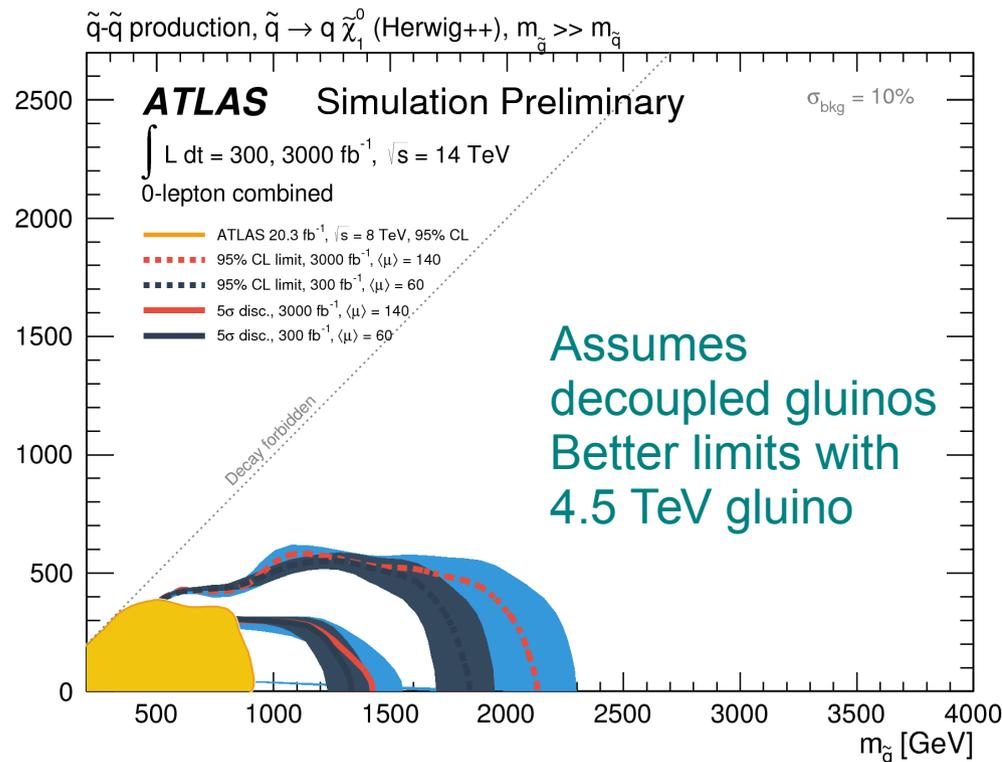
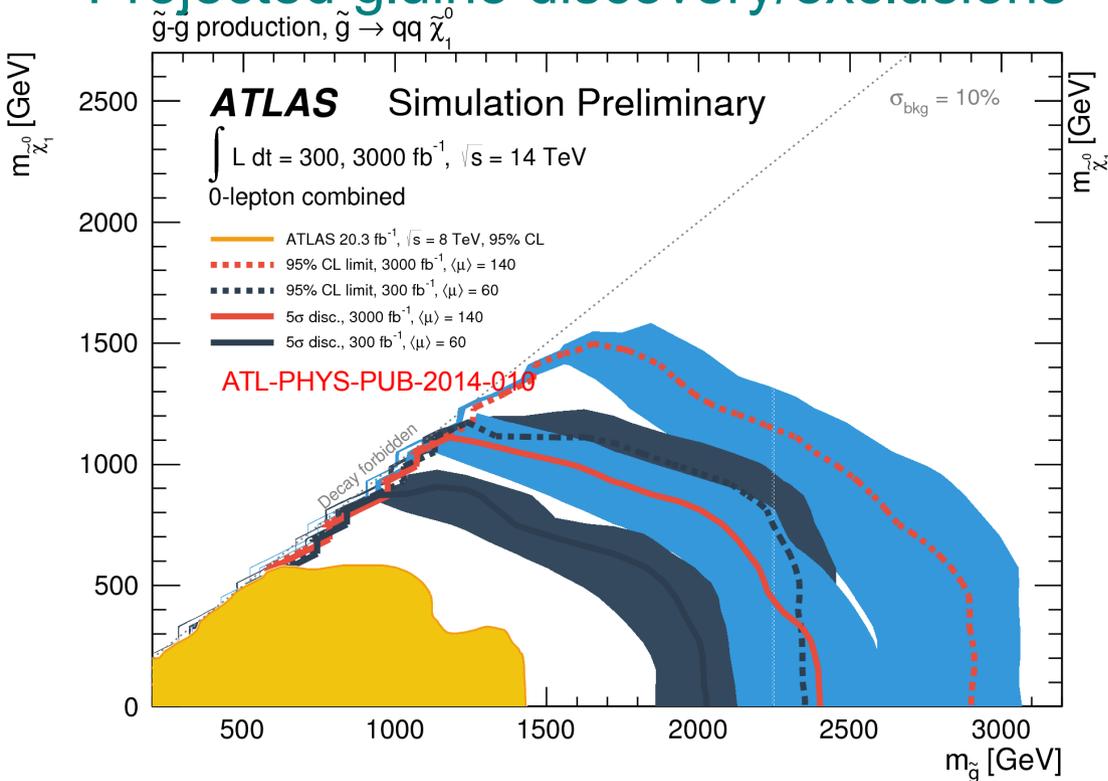
Unfortunately, no public projection for 2015 sensitivity

Longer-term sensitivity studied in good detail

- Main backgrounds simulated using parameterized detector response
- Squarks and gluino searches will reach into multi-TeV space
- If no gluinos found by end of LHC, Natural SUSY disfavored
- Already in 2015, expect that just a few fb⁻¹ enough to exceed 2012 exclusions levels, particularly for heavy gluinos

Projected gluino discovery/exclusions

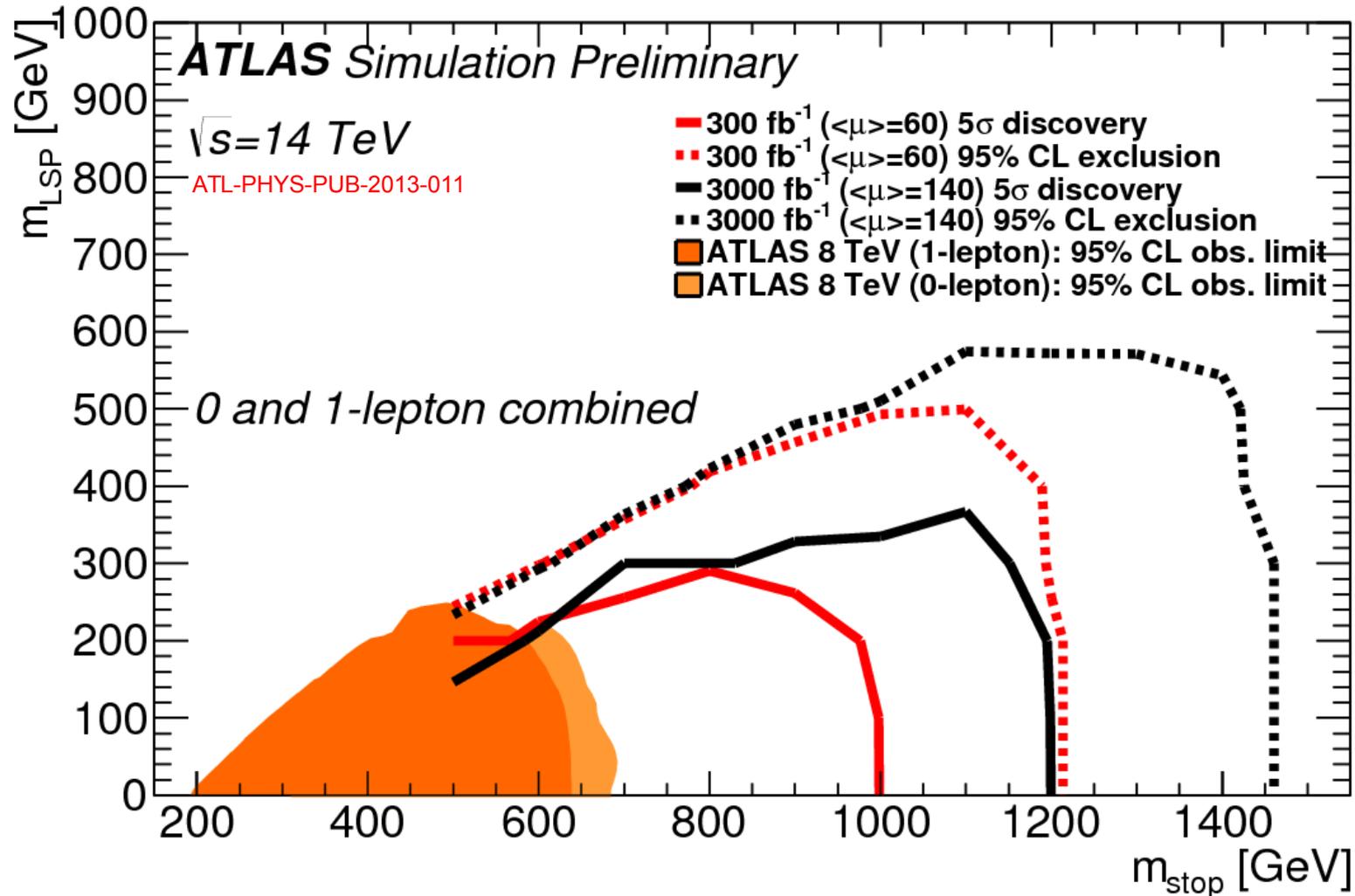
Projected squark discovery/exclusions



3rd Generation Squark Searches

Similar studies done for stop and sbottom production

- Expect to discover stop/sbottom up to 1 TeV with 300 fb^{-1}
- Sensitivity up to almost 1.5 TeV with complete dataset in the favorable decay mode $t \rightarrow t\tilde{\chi}_1^0$
- For 2015 will need $>5\text{fb}^{-1}$ to excess 2012 high mass limit

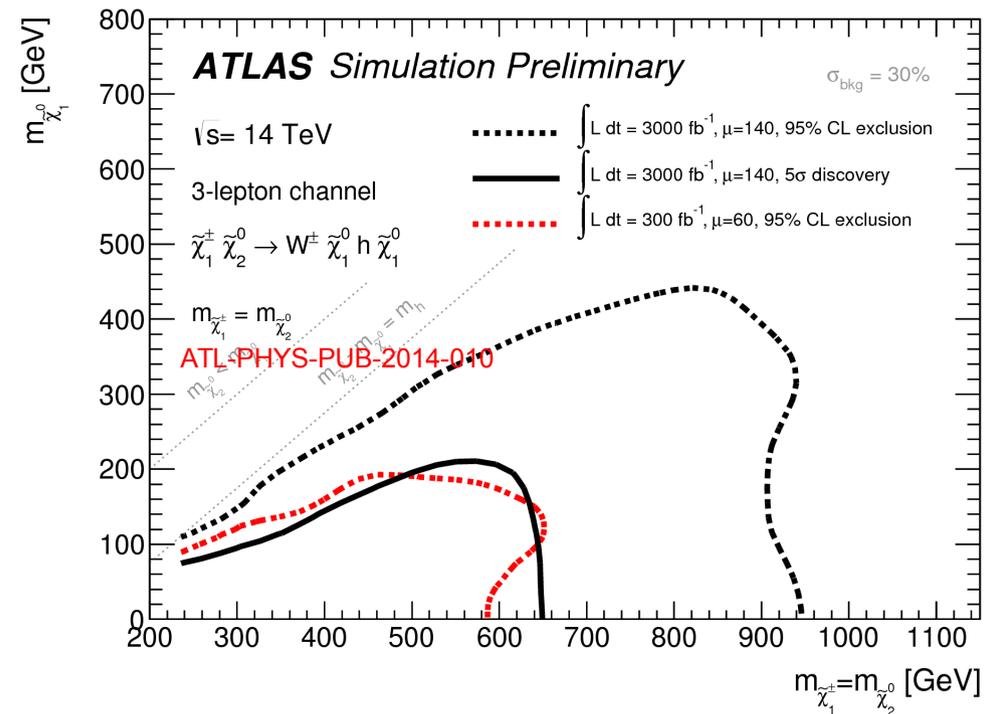
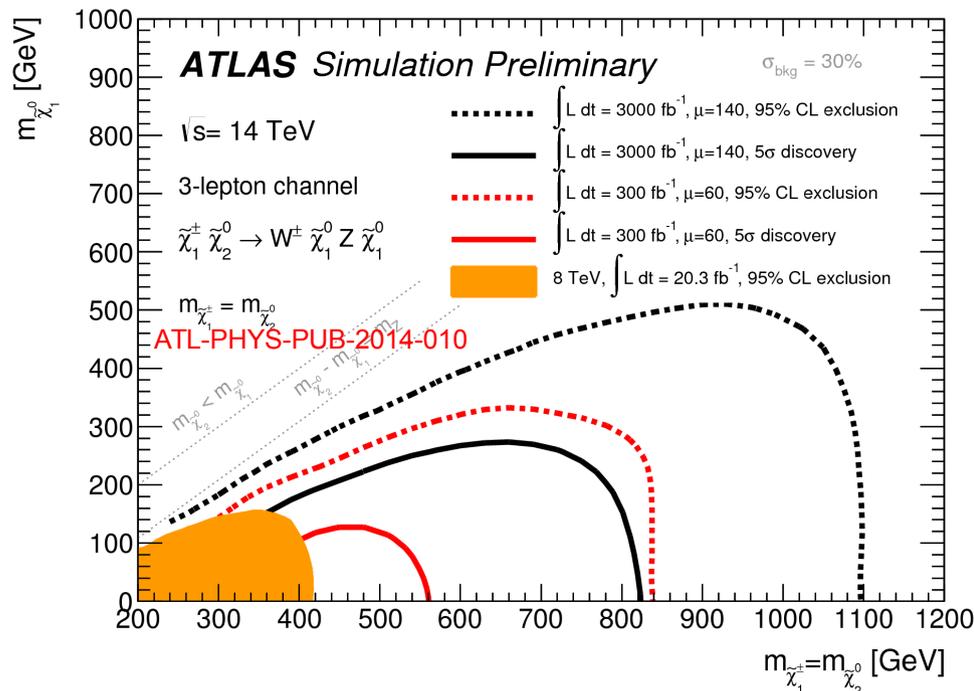


Chargino/Neutralino Searches

Searches for EW production sees little gain from energy increase
Large luminosity needed to really push sensitivity

High luminosity projections for direct chargino/neutralino decays

- Use 3-lepton channel as very low background (primarily WZ)
- Even with 300 fb^{-1} not much possibility of 5σ discovery
- With 3000 fb^{-1} , will approach 1 TeV for exclusion sensitivity, 650-800 TeV discover with a light LSP



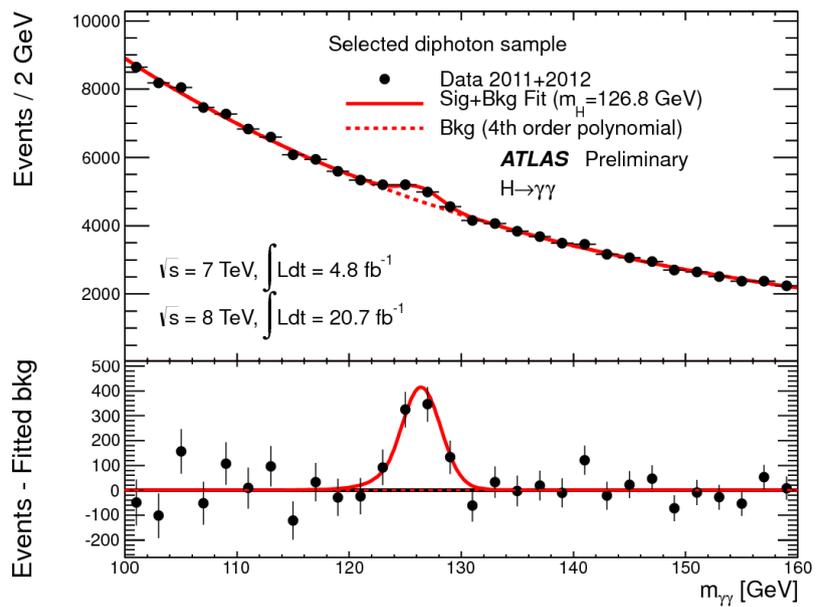
Summary

Summary

- Extensive search for SUSY at the LHC in Run-1
- Provide comprehensive coverage of SUSY detectable with current LHC luminosity and energy
- No significant signals seen, excluding a big chunk of SUSY with sparticles below 1 TeV
 - do have one signal region at 3σ significance
- Even greater sensitivity in upcoming LHC run(s)

2012: Higgs Discovery year

2015: SUSY Discovery year?

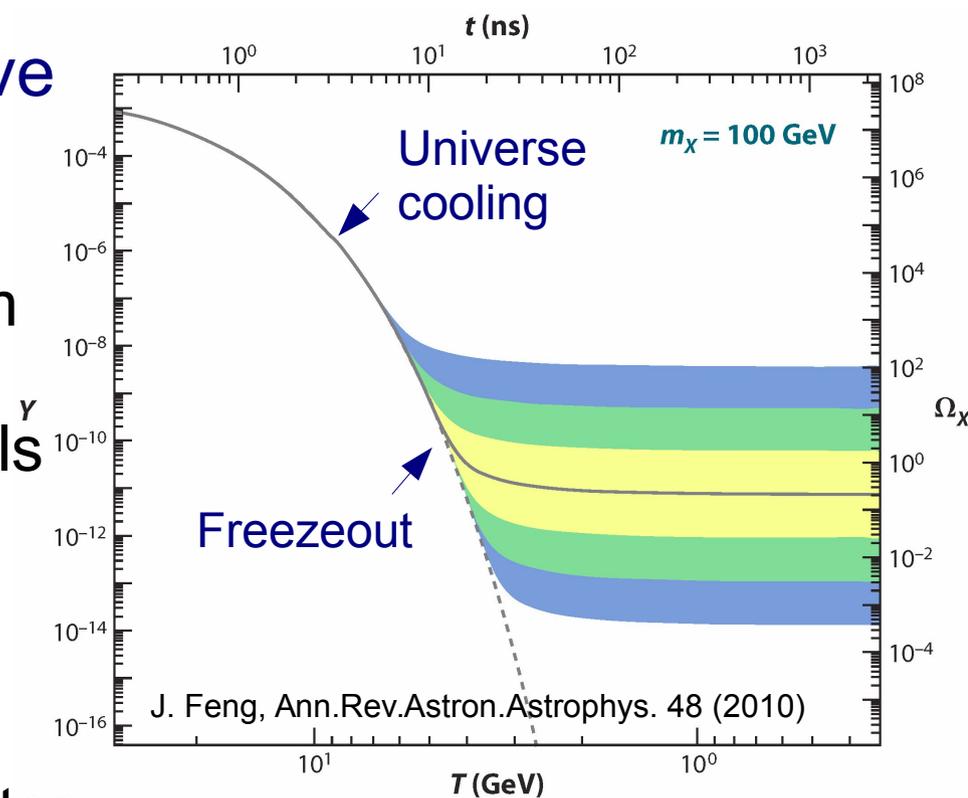


Backup

“WIMP Miracle”

A new, weakly-interacting massive particle (WIMP) is an excellent dark matter candidate

- Assume initial thermal equilibrium
 $X X \leftrightarrow SM SM$
 - Density will drop as universe cools
 $X X \rightarrow SM SM$
 - Annihilation stops as universe expands
 $X X \not\leftrightarrow SM SM$
- Leaving relic density as dark matter



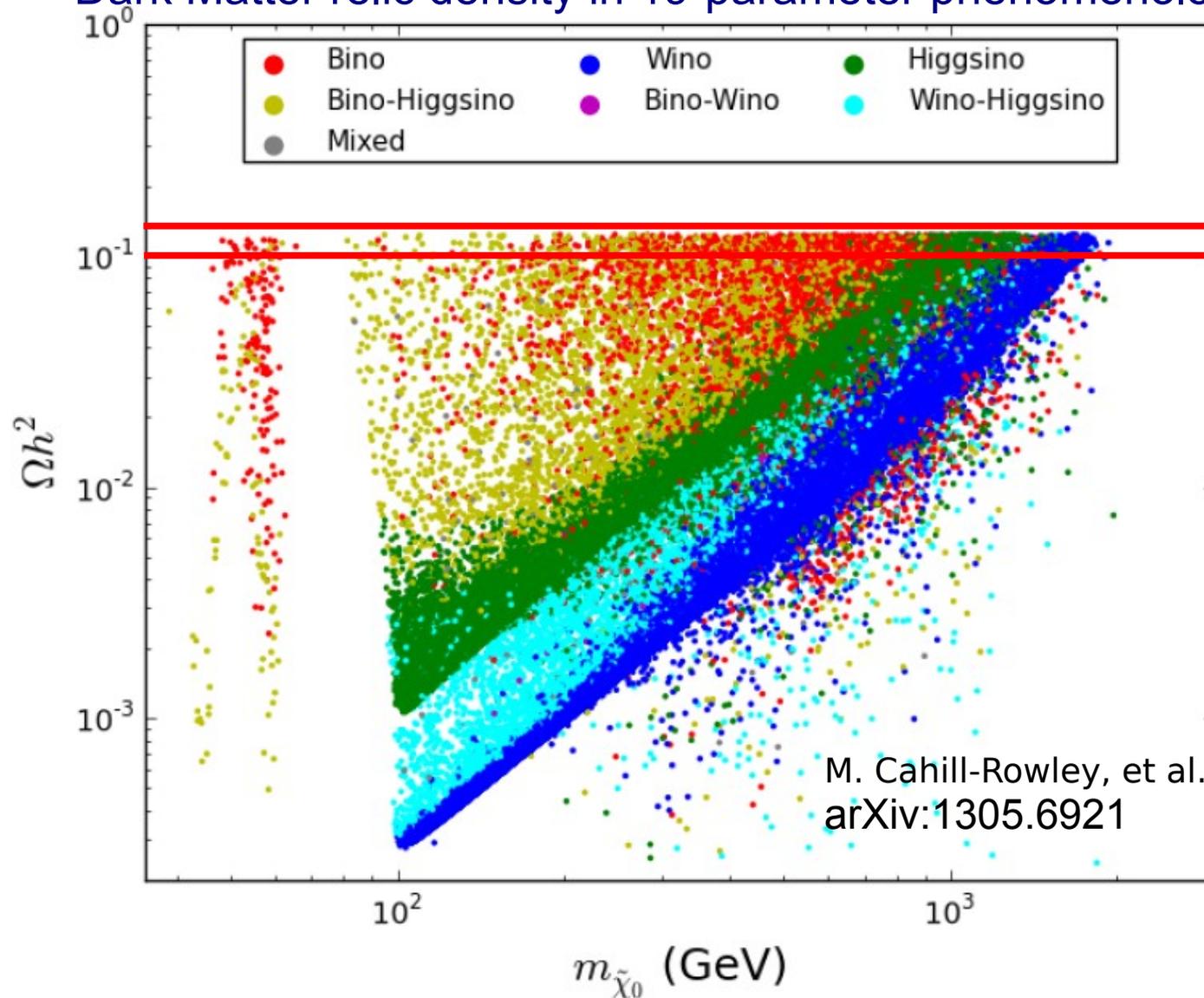
If $m_X \sim 100 \text{ GeV}$ and has weak-scale interaction, the thermal relic density $\Omega_X \sim \Omega_{DM}$

Supersymmetry provides a natural WIMP in the lightest supersymmetric particle

Supersymmetry and Dark Matter

LSP mass strongly dependent on SUSY model, but should normally be 100 GeV to 2 TeV to match Dark Matter observation

Dark Matter relic density in 19-parameter phenomenological MSSM



Observed Dark Matter density

All 19 pMSSM model parameters are varied randomly and each model subjected to indirect constraints

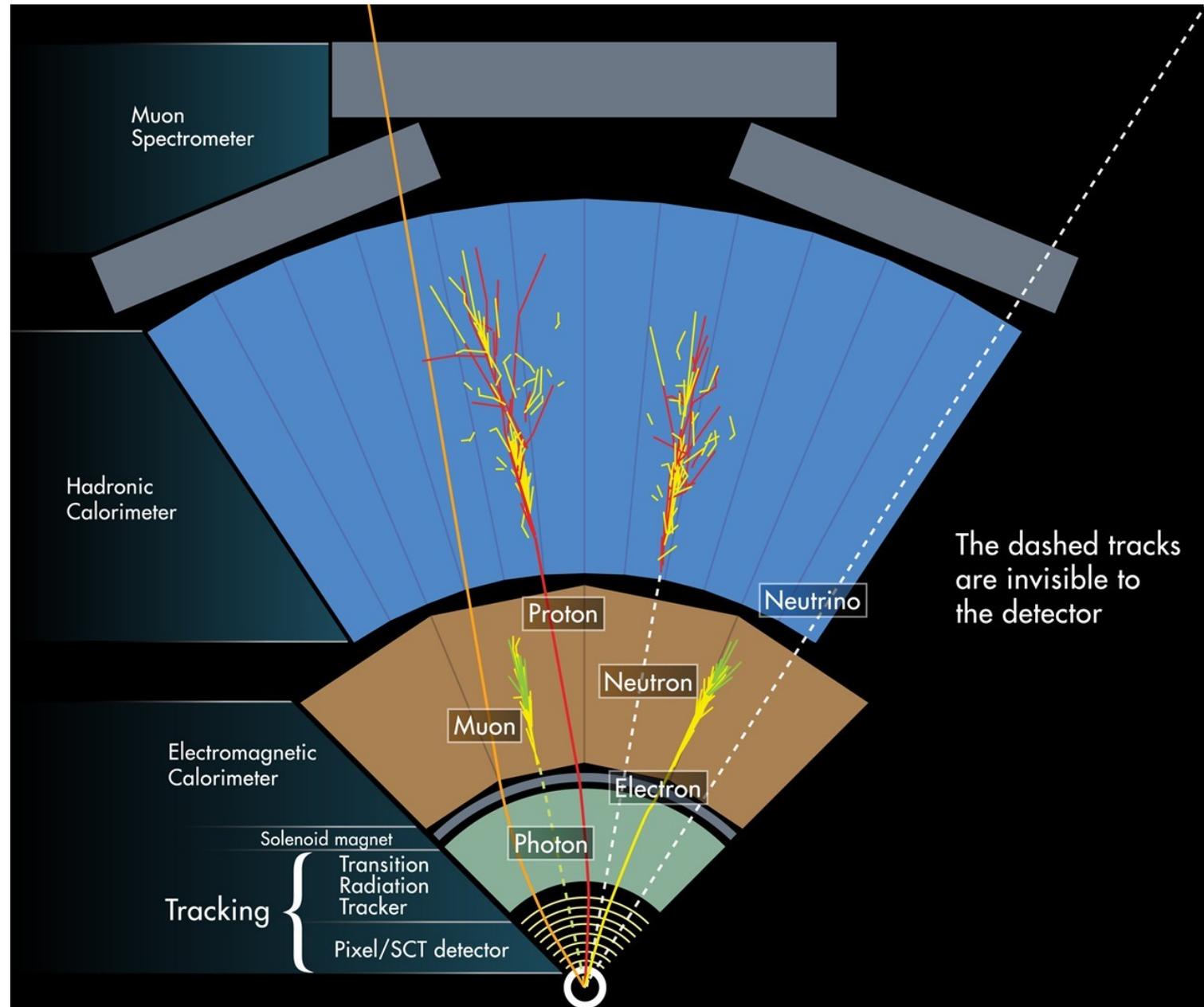
Particle Reconstruction and Identification

Schematic view of ATLAS detector

Hermetic Detector

Almost all particles fully reconstructed and identified in one or more sub-detectors

Only weakly-interacting particles pass through undetected



Pileup

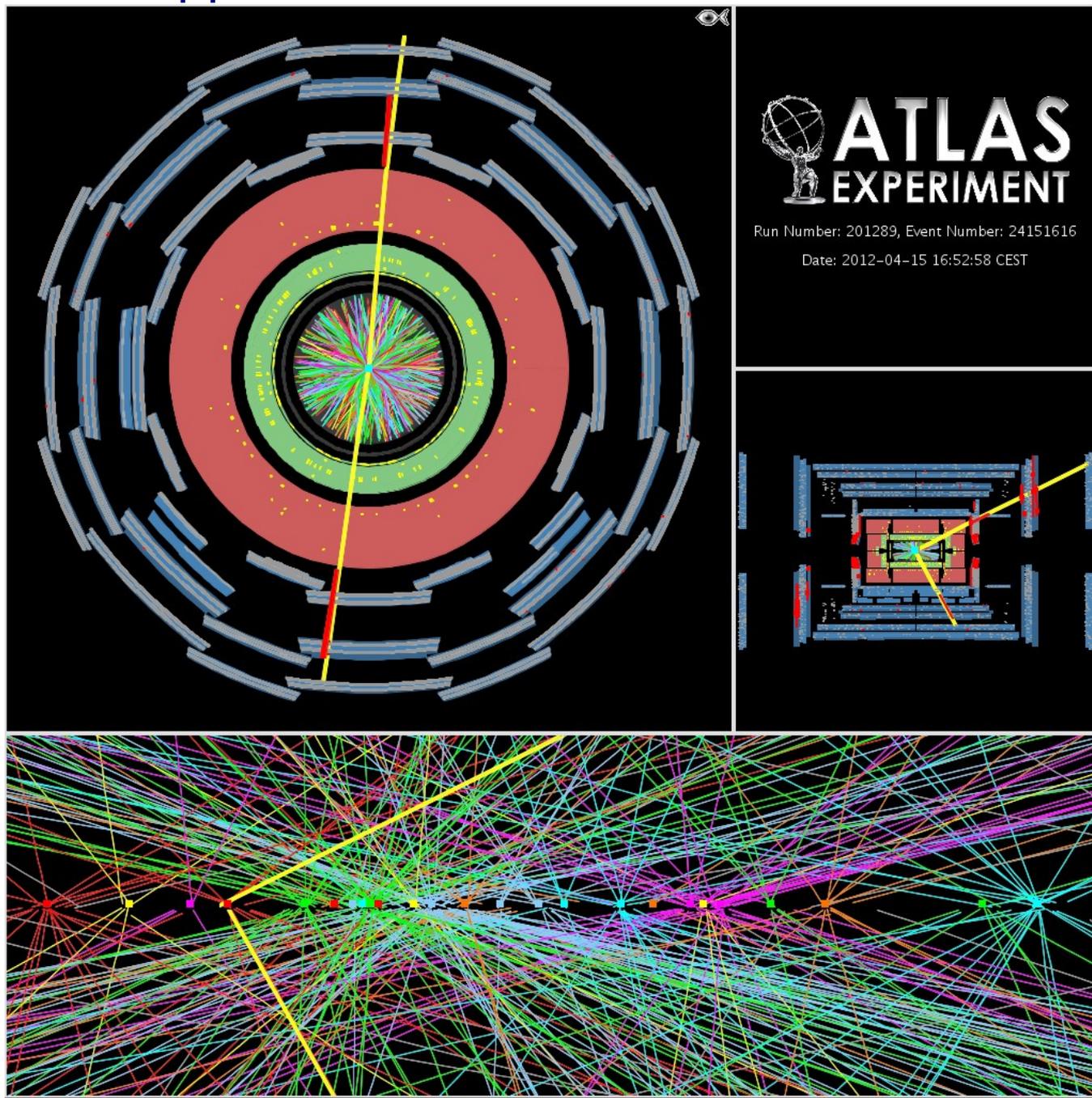
$Z \rightarrow \mu\mu$ event with 25 reconstructed vertices

High luminosity at a cost

Record luminosity achieved by having 20-40 interactions per beam crossing every 50 ns

Design was peak of 23 interaction per 25 ns

Results in degraded performance, but mostly compensated by use of smarter selection algorithms



Inclusive 0-lepton Signal Regions

Requirement	Signal Region					
	2jl	2jm	2jt	2jW	3j	4jW
$E_T^{\text{miss}} [\text{GeV}] >$	160					
$p_T(j_1) [\text{GeV}] >$	130					
$p_T(j_2) [\text{GeV}] >$	60					
$p_T(j_3) [\text{GeV}] >$	-				60	40
$p_T(j_4) [\text{GeV}] >$	-					40
$\Delta\phi(\text{jet}_{1,2,(3)}, \mathbf{E}_T^{\text{miss}})_{\text{min}} >$	0.4					
$\Delta\phi(\text{jet}_{i>3}, \mathbf{E}_T^{\text{miss}})_{\text{min}} >$	-					0.2
W candidates	-			$2(W \rightarrow j)$	-	$(W \rightarrow j) + (W \rightarrow jj)$
$E_T^{\text{miss}}/\sqrt{H_T} [\text{GeV}^{1/2}] >$	8	15		-		
$E_T^{\text{miss}}/m_{\text{eff}}(N_j) >$	-			0.25	0.3	0.35
$m_{\text{eff}}(\text{incl.}) [\text{GeV}] >$	800	1200	1600	1800	2200	1100

Requirement	Signal Region								
	4jl-	4jl	4jm	4jt	5j	6jl	6jm	6jt	6jt+
$E_T^{\text{miss}} [\text{GeV}] >$	160								
$p_T(j_1) [\text{GeV}] >$	130								
$p_T(j_2) [\text{GeV}] >$	60								
$p_T(j_3) [\text{GeV}] >$	60								
$p_T(j_4) [\text{GeV}] >$	60								
$p_T(j_5) [\text{GeV}] >$	-				60				
$p_T(j_6) [\text{GeV}] >$	-					60			
$\Delta\phi(\text{jet}_{1,2,(3)}, \mathbf{E}_T^{\text{miss}})_{\text{min}} >$	0.4								
$\Delta\phi(\text{jet}_{i>3}, \mathbf{E}_T^{\text{miss}})_{\text{min}} >$	0.2								
$E_T^{\text{miss}}/\sqrt{H_T} [\text{GeV}^{1/2}] >$	10		-						
$E_T^{\text{miss}}/m_{\text{eff}}(N_j) >$	-		0.4	0.25	0.2			0.25	0.15
$m_{\text{eff}}(\text{incl.}) [\text{GeV}] >$	700	1000	1300	2200	1200	900	1200	1500	1700

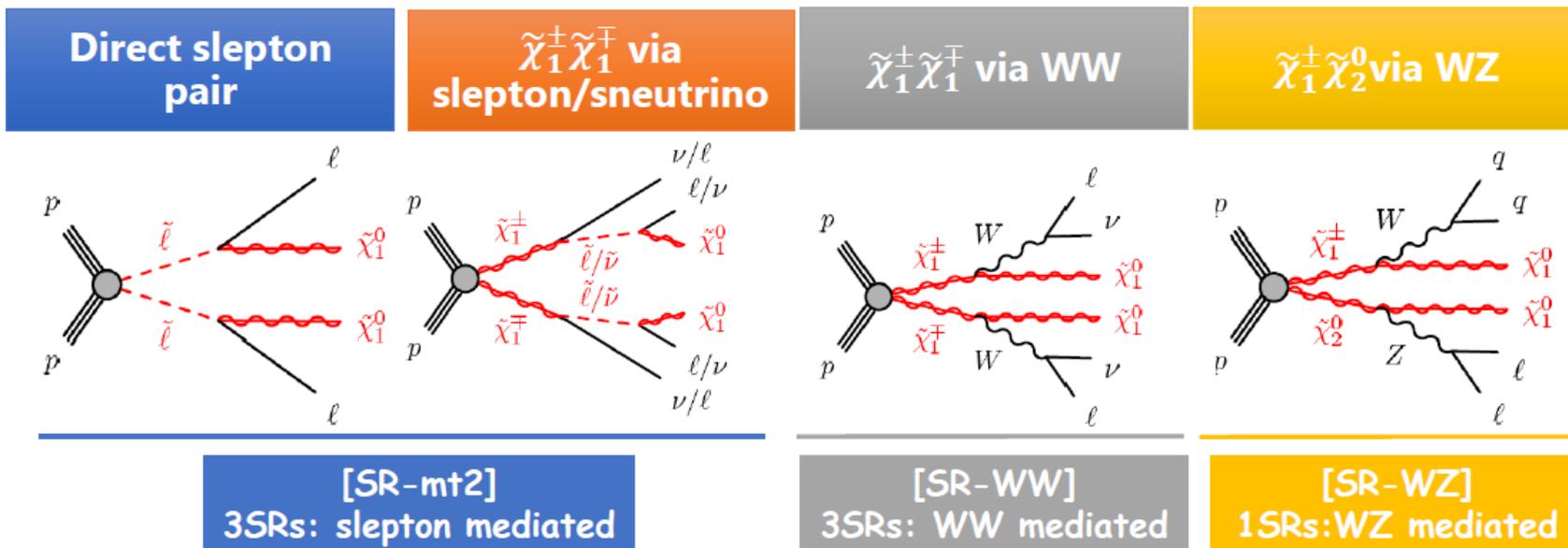
Top polarisation and Spin Correlations

- Top quark lifetime of $\sim 3 \cdot 10^{-25}$ sec is much shorter than hadronisation time
 - Top decays as a bare quark, and does not form hadrons
 - Top spin info is not 'corrupted' by QCD interactions, transferred to decay products
- Angular decay distribution: $\frac{1}{\Gamma} \frac{d\Gamma}{d\cos(\theta_i)} = (1 + \alpha_i \mathbf{P} \cos(\theta_i))/2$
 - θ_i angle between top decay product i and top polarisation P along chosen axis
 - α_i is spin analysing power: $\sim \pm 1$ for charged leptons, $-0.966 / -0.393$ for d / b quark
 - Normally use helicity basis, chose quantisation axis as top quark momentum direction in tT rest frame
- Negligible polarisation in SM, but spins of t and T are correlated

$$\frac{1}{\sigma} \frac{d\sigma}{d\cos(\theta_+) d\cos(\theta_-)} = \frac{1}{4} (1 + A \alpha_+ \alpha_- \cos(\theta_+) \cos(\theta_-)) \quad A = \frac{N_{\text{like}} - N_{\text{unlike}}}{N_{\text{like}} + N_{\text{unlike}}} = \frac{N(\uparrow\uparrow) + N(\downarrow\downarrow) - N(\uparrow\downarrow) - N(\downarrow\uparrow)}{N(\uparrow\uparrow) + N(\downarrow\downarrow) + N(\uparrow\downarrow) + N(\downarrow\uparrow)}$$

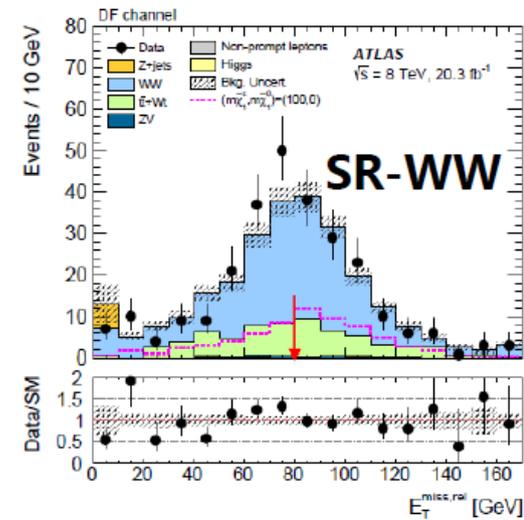
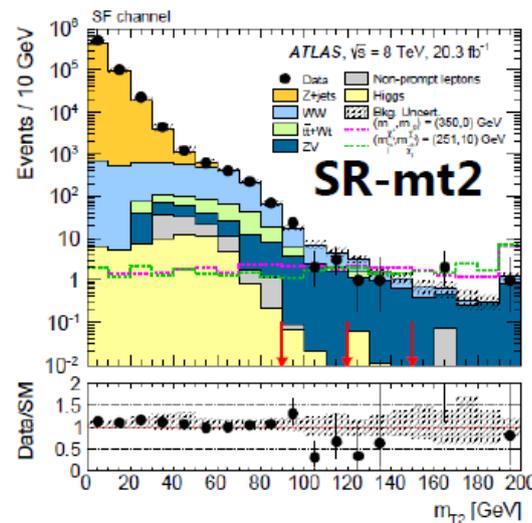
- Can be measured from dilepton $\Delta\phi_{\parallel}$, or observables involving $\cos\theta_i$
 - $\Delta\phi_{\parallel}$ is straightforward to measure precisely
 - $\cos\theta_i$ requires full event reconstruction (dilepton or l+jets events)

2-lepton EW Search



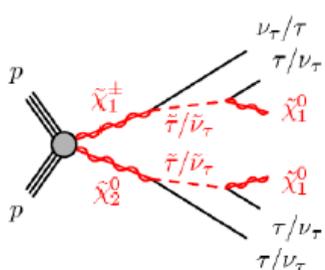
* $\tilde{\chi}_1^\pm, \tilde{\chi}_2^0$
(pure wino, mass degenerate);
 $\tilde{\chi}_1^0$ (pure bino)

- 7 SRs designed targeting different models.
 - The same flavor and different flavor are considered separately in each SR.
- Main backgrounds: top-quark(ttbar and Wt) and dibosons.
 - For SM ttbar and WW: mt2 has an upper end-point at the W mass.



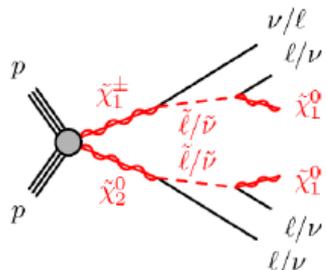
3 lepton (e/ μ / τ) EW Search

$\tilde{\chi}_1^\pm \tilde{\chi}_2^0$ via
stau/sneutrino



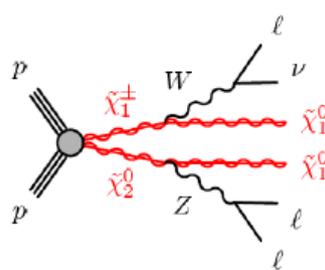
1 SR:
2 taus + light lepton

$\tilde{\chi}_1^\pm \tilde{\chi}_2^0$ via
slepton/sneutrino

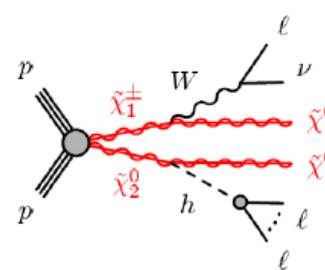


1 SR: 0 tau + light leptons
characterized by 20 independent bins

$\tilde{\chi}_1^\pm \tilde{\chi}_2^0$ via WZ



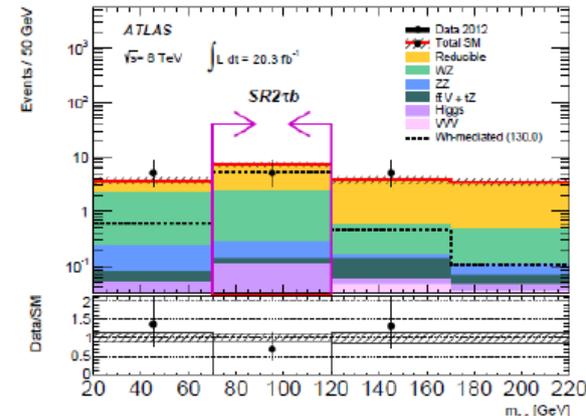
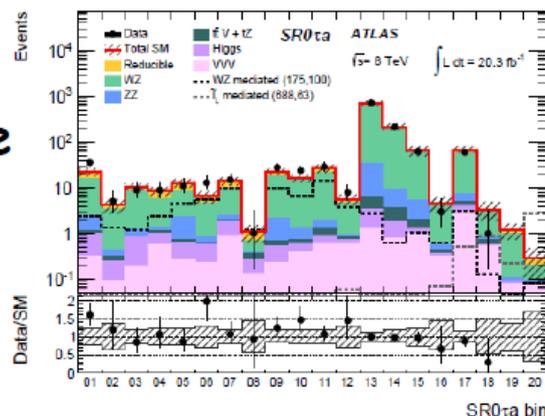
$\tilde{\chi}_1^\pm \tilde{\chi}_2^0$ via WH



3 SRs: 0,1,2 taus
+ light leptons

* $\tilde{\chi}_1^\pm, \tilde{\chi}_2^0$
(pure wino,
mass
degenerate):
 $\tilde{\chi}_1^0$ (pure bino)

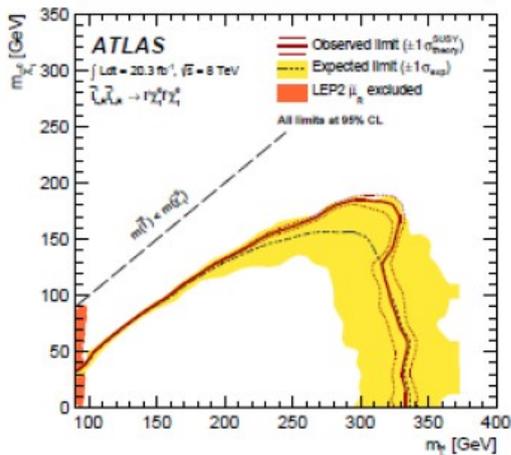
- Analysis includes up to 2 hadronic taus.
- 5 SRs are defined according to the flavor and charge of the leptons, targeting different models.
- Main backgrounds: diboson, triboson, $t\bar{t}b\bar{v}$, tZ and VH .



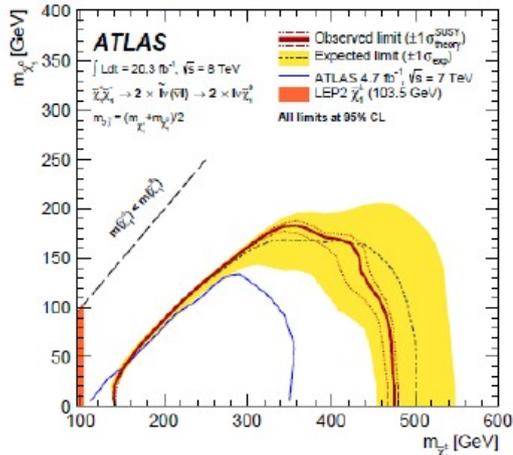
2/3 Lepton Results

2-lepton searches

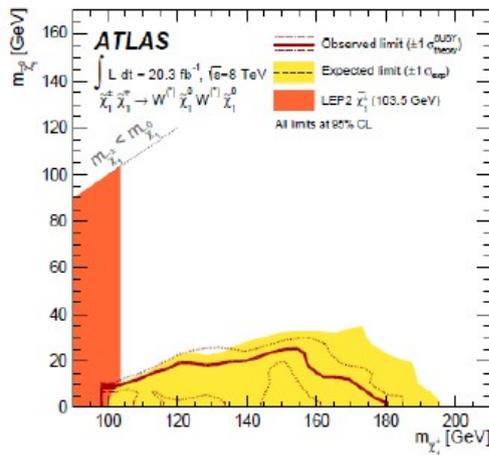
Direct slepton pair



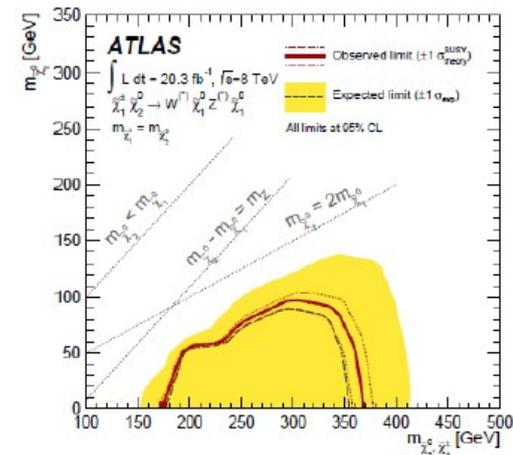
$\tilde{\chi}_1^\pm \tilde{\chi}_1^\mp$ via slepton/sneutrino



$\tilde{\chi}_1^\pm \tilde{\chi}_1^\mp$ via WW

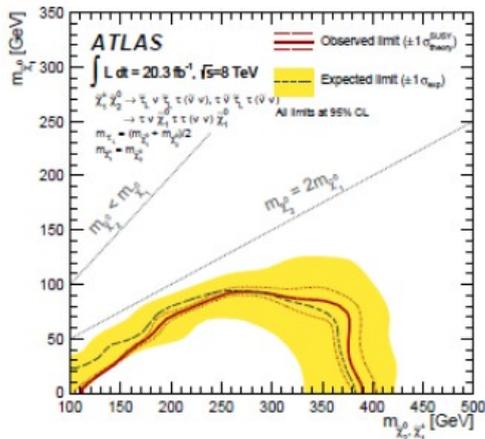


$\tilde{\chi}_1^\pm \tilde{\chi}_2^0$ via WZ

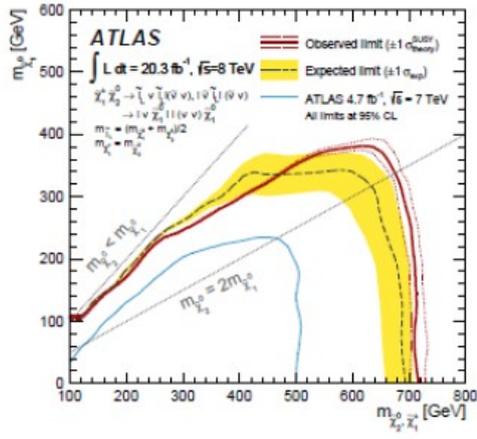


3-lepton searches

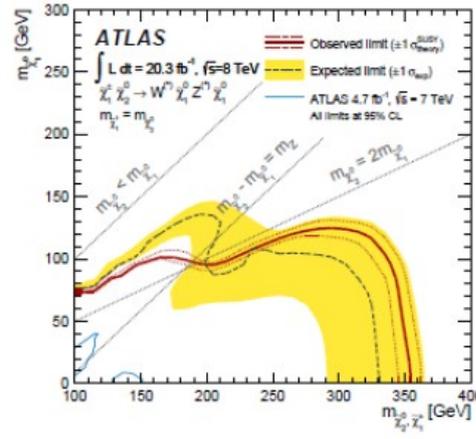
$\tilde{\chi}_1^\pm \tilde{\chi}_2^0$ via stau/sneutrino



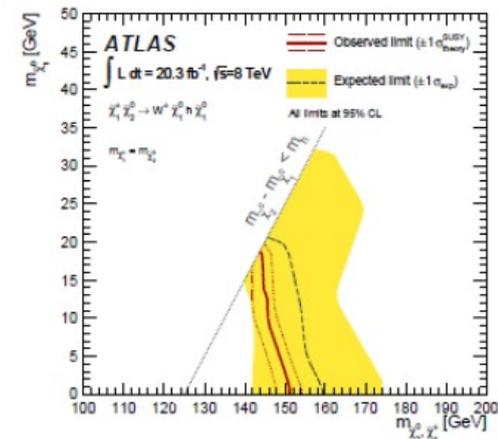
$\tilde{\chi}_1^\pm \tilde{\chi}_2^0$ via slepton/sneutrino



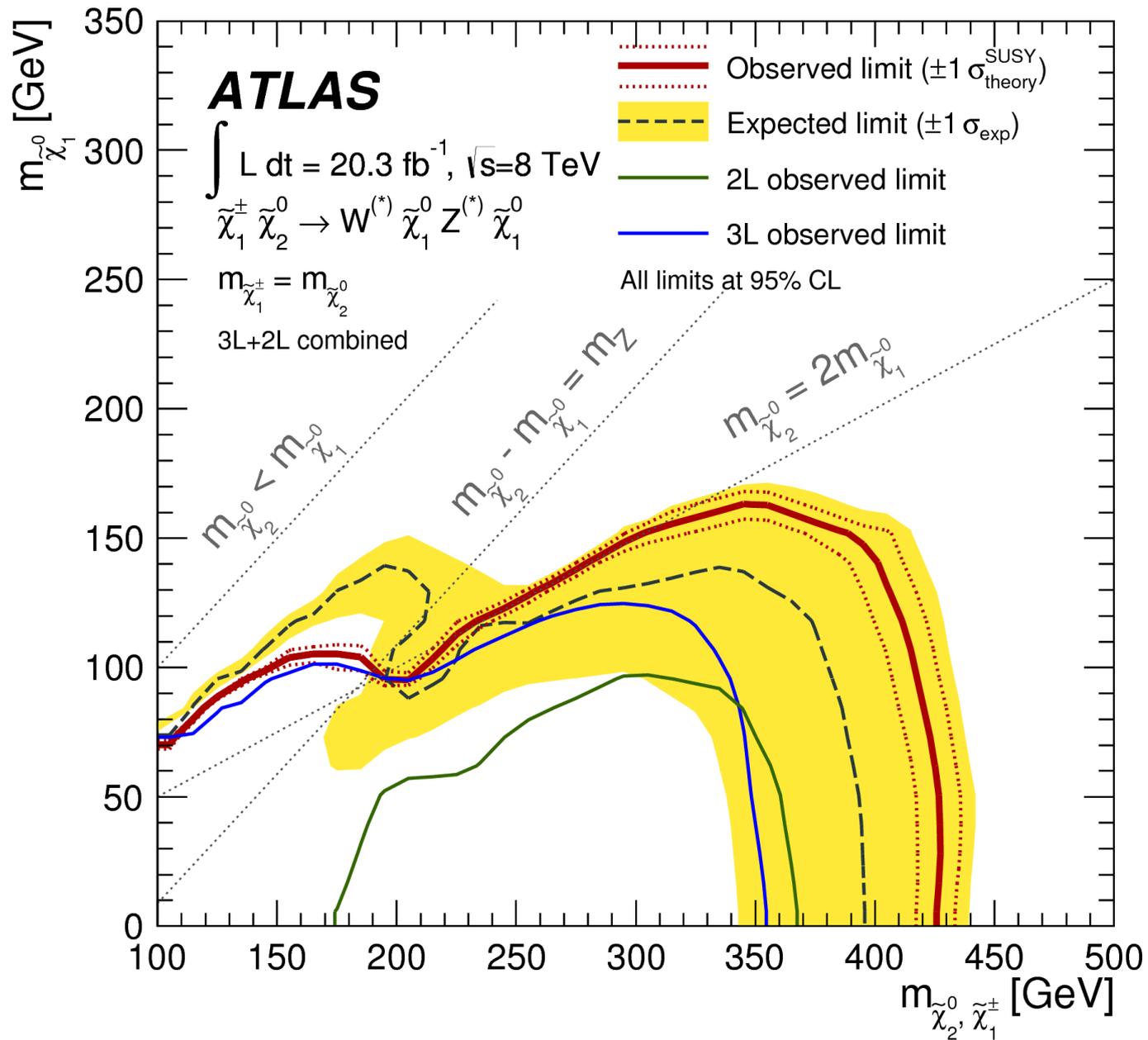
$\tilde{\chi}_1^\pm \tilde{\chi}_2^0$ via WZ



$\tilde{\chi}_1^\pm \tilde{\chi}_2^0$ via WH



Combined 2/3 Lepton Result



Long-lived

What if gluino is just a little long-lived, about 1 ns? (mini-split SUSY)
Standard jets+MET SUSY searches should still apply (up to what lifetime?)

- Leptons vetos may start to fail impact-parameter cuts (when?)
- Jets will start to be identified as b-jets (when?)
- Jets may fail cleaning cuts using track pT fraction, EM fraction (when?)

First explicit re-interpretation of prompt SUSY searches for long-lived gluinos!

