

Phenomenology of photon colliders: Higgs pair production and coupling measurements

Marten Berger (University of Hamburg)

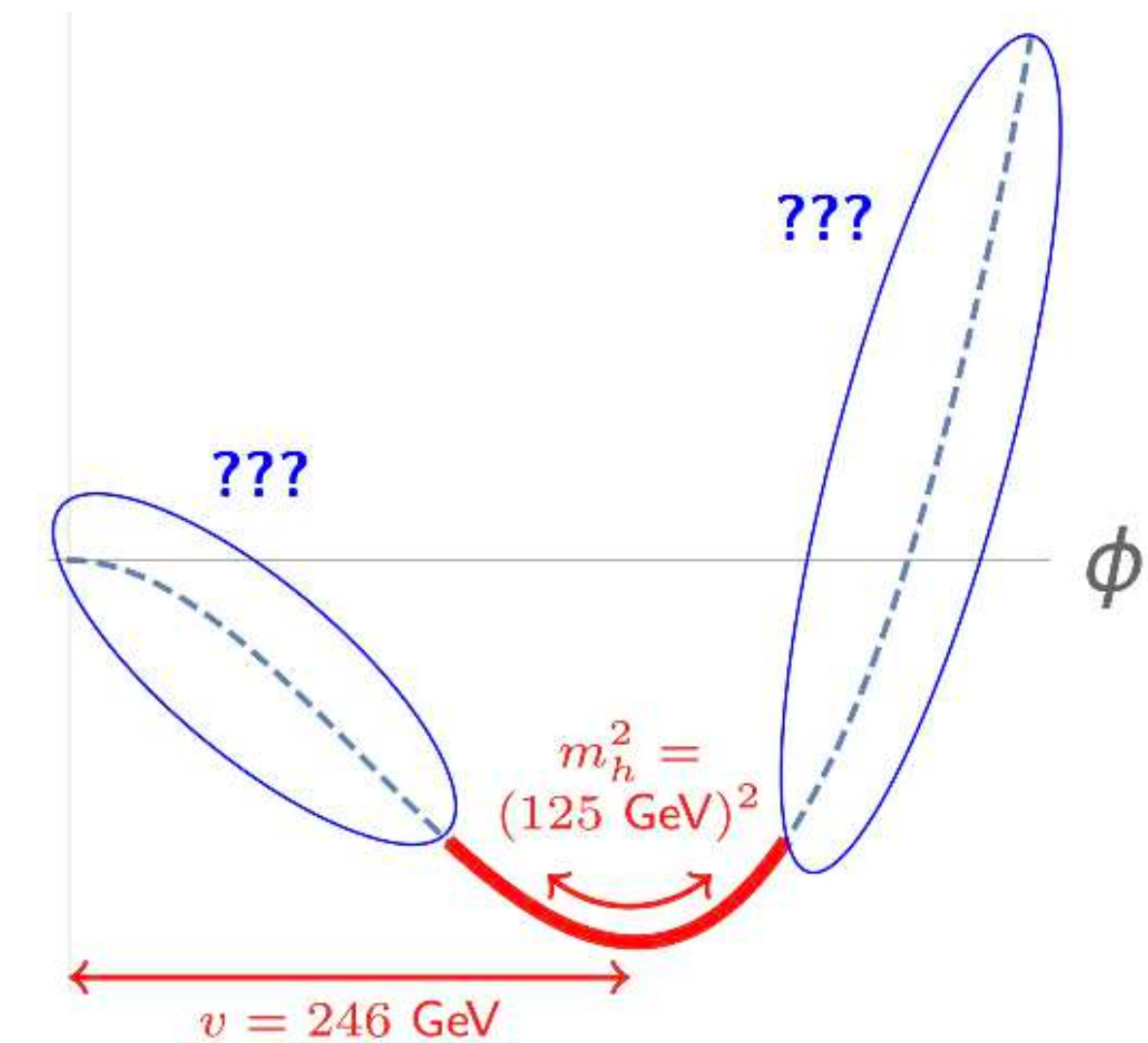
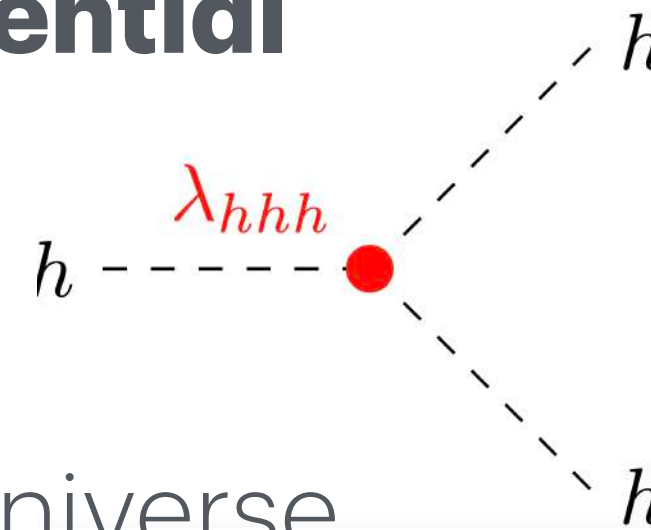
based on [2603.00264](#) (M. B., G. Moortgat-Pick) and contributions to [2503.24049](#) (H. Abramowicz, M.B. et al.)
and work in progress (M. B., J. Braathen, G. Moortgat-Pick, W. Naskar, G. Weiglein)

ICEPP Seminar, The University of Tokyo, Japan

9 September 2026

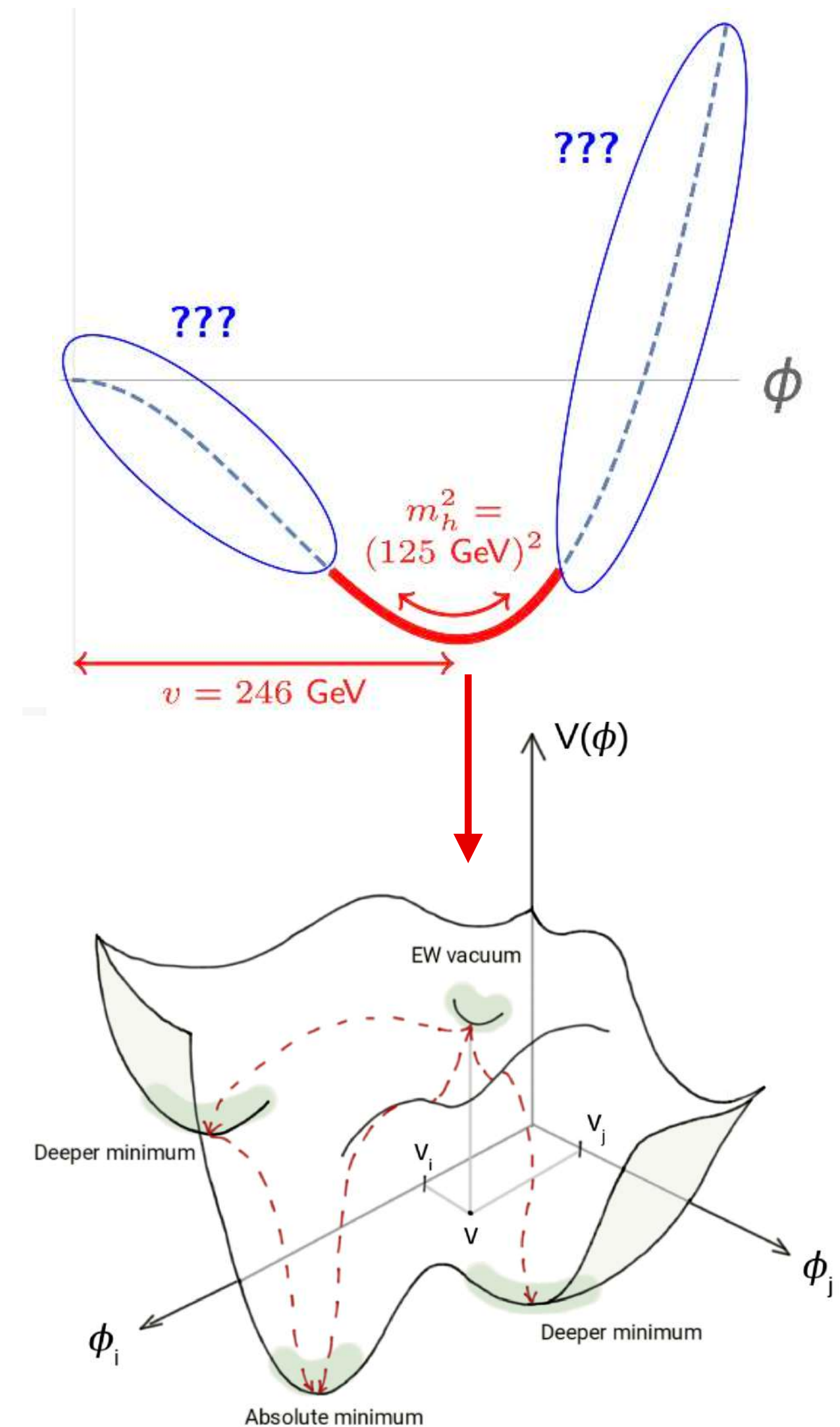
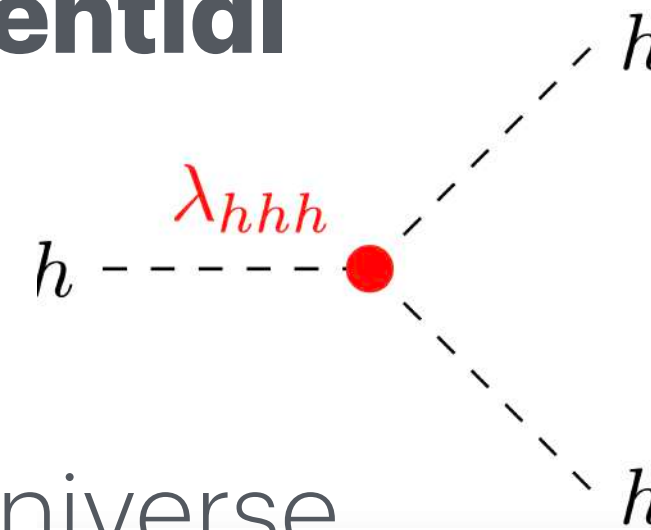
Motivation - Higgs pair production

- **4 July 2012** Higgs discovery was announced
 - Is it the SM Higgs boson?
- Very little is known about the **Higgs potential**
- $V = \frac{1}{2}m_h^2 h^2 + v\lambda_{hhh}h^3 + \lambda_{hhhh}h^4 + \dots$
- Holds strong implications for the early universe
 - Strong first-order electroweak phase transition
 - 3rd Sakharov condition [Sakharov '67]



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 - 3rd Sakharov condition [Sakharov '67]
- Beyond-the-Standard-Model theories can offer multiple field directions and minima

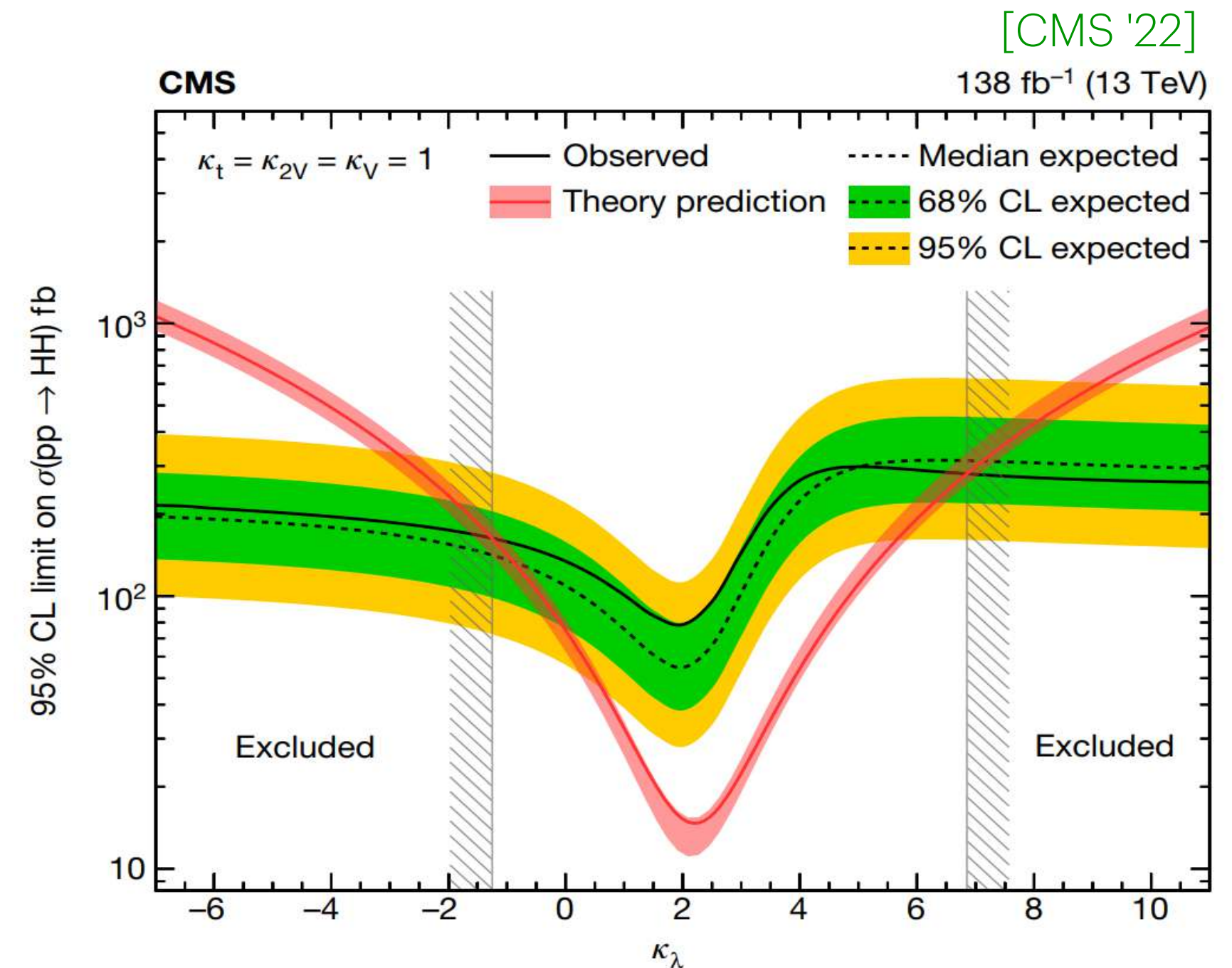
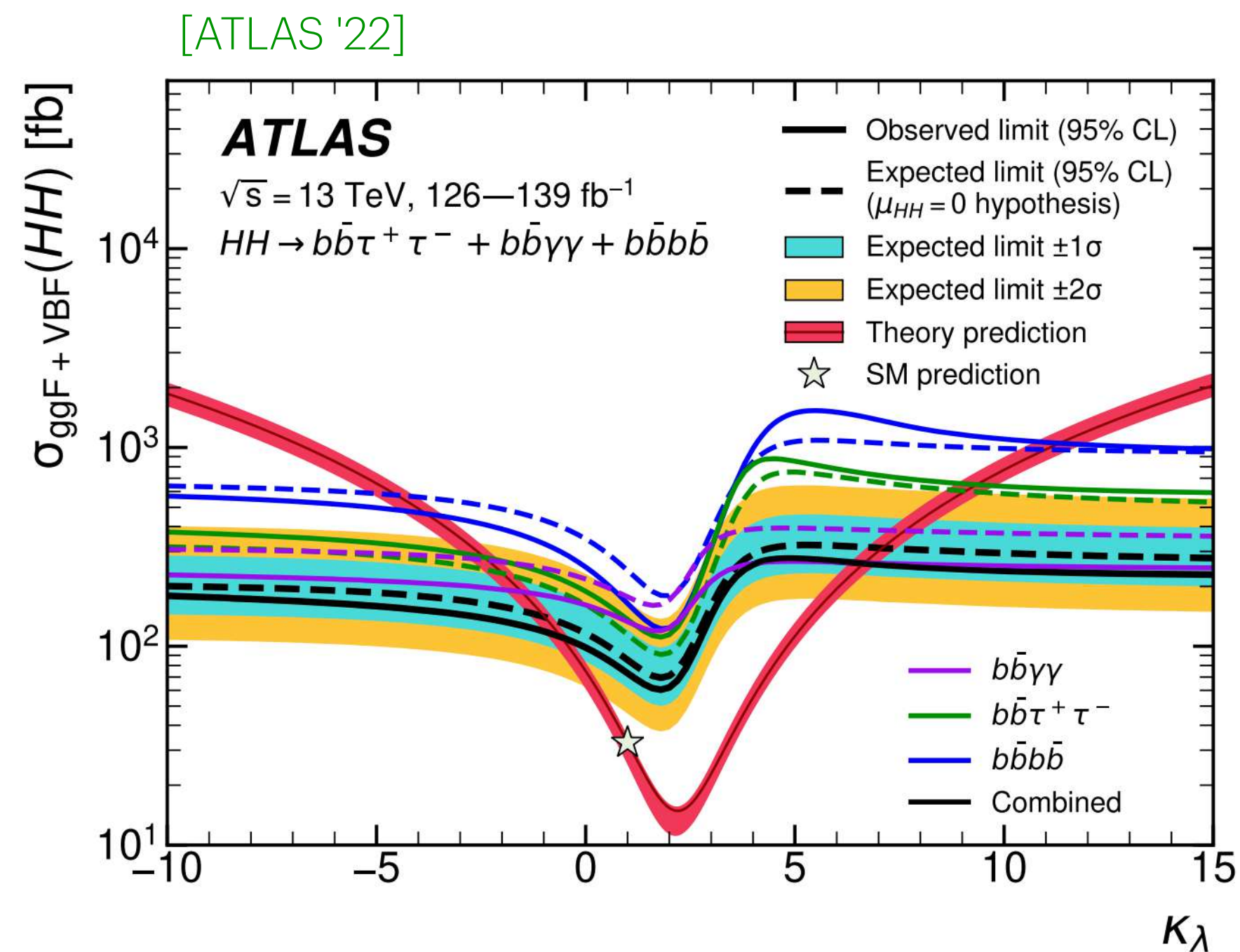


Motivation - Higgs pair production

- Higgs pair production at the LHC (gluon fusion, $gg \rightarrow hh$)
- ATLAS $[-1.2 < \kappa_\lambda < 7.2]$ [ATLAS '24] and CMS $[-1.2 < \kappa_\lambda < 7.5]$ [CMS '24]
- Combined results $[-0.71 < \kappa_\lambda < 6.1]$ [ATLAS and CMS '26]

Notation:

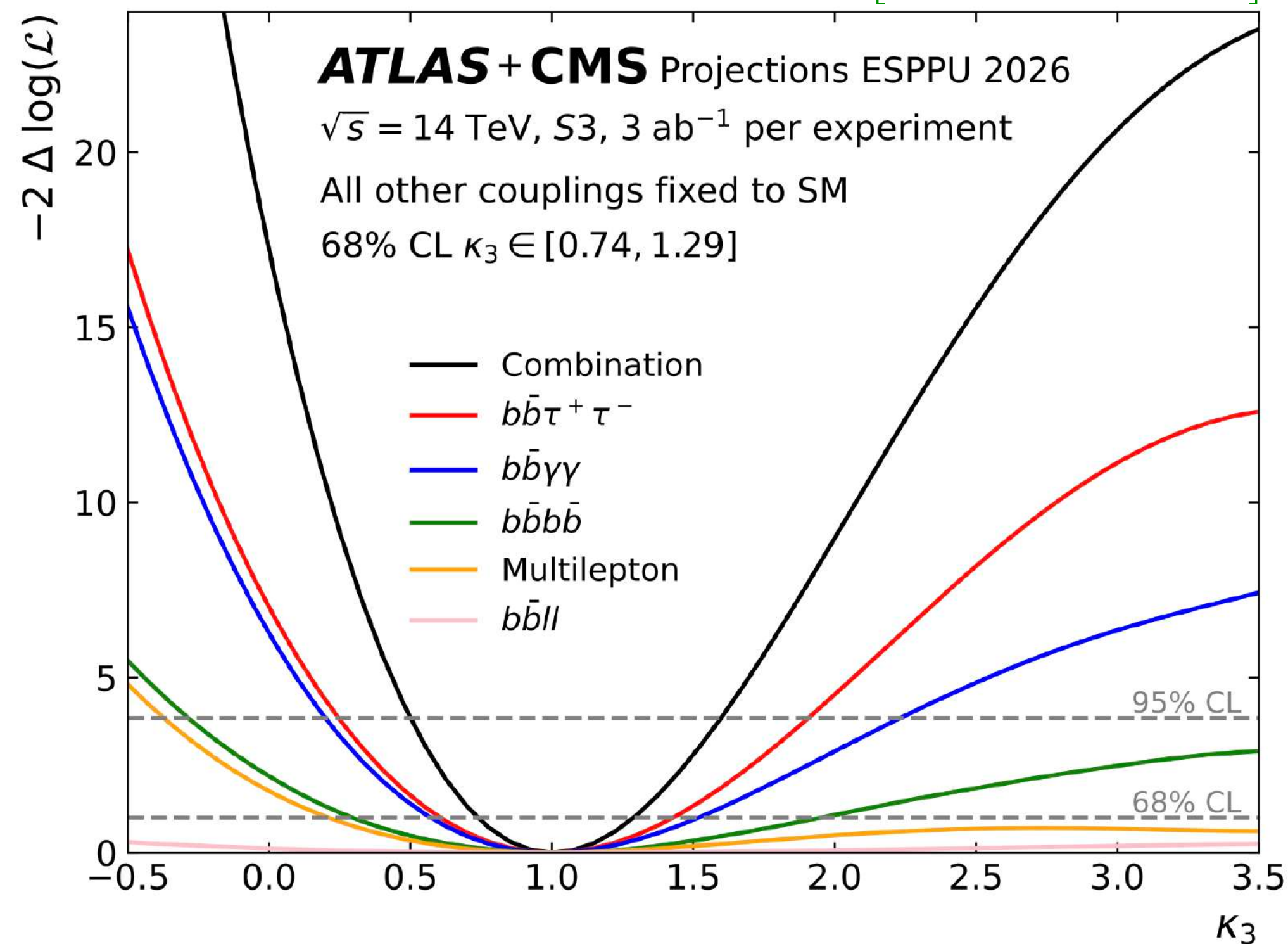
$$\kappa_\lambda = \frac{\lambda_{hhh}}{\lambda_{hhh}^{\text{SM}(0)}}$$



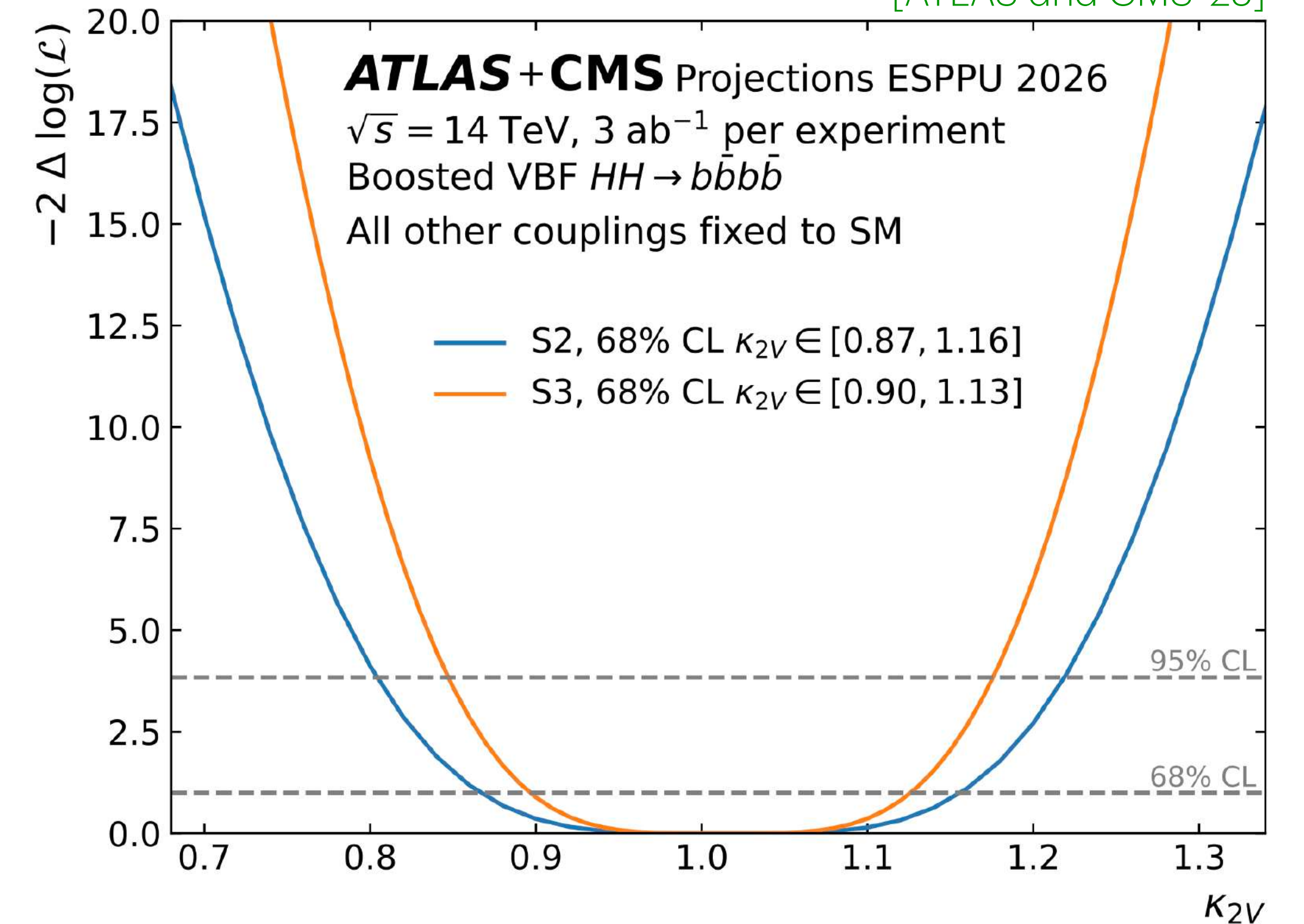
HL-LHC projected limits

- Projected sensitivity on κ_λ (here κ_3) and κ_{2V} assuming the Standard Model
- $\kappa_\lambda \in [0.53, 1.62]$ and $\kappa_{2V} \in [0.85, 1.17]$ at 95% CL

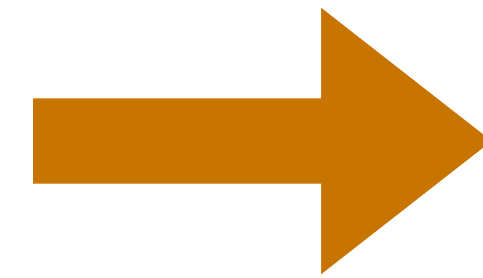
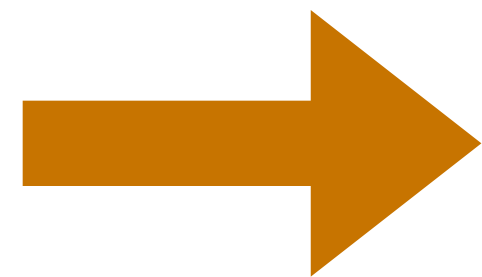
[ATLAS and CMS '25]



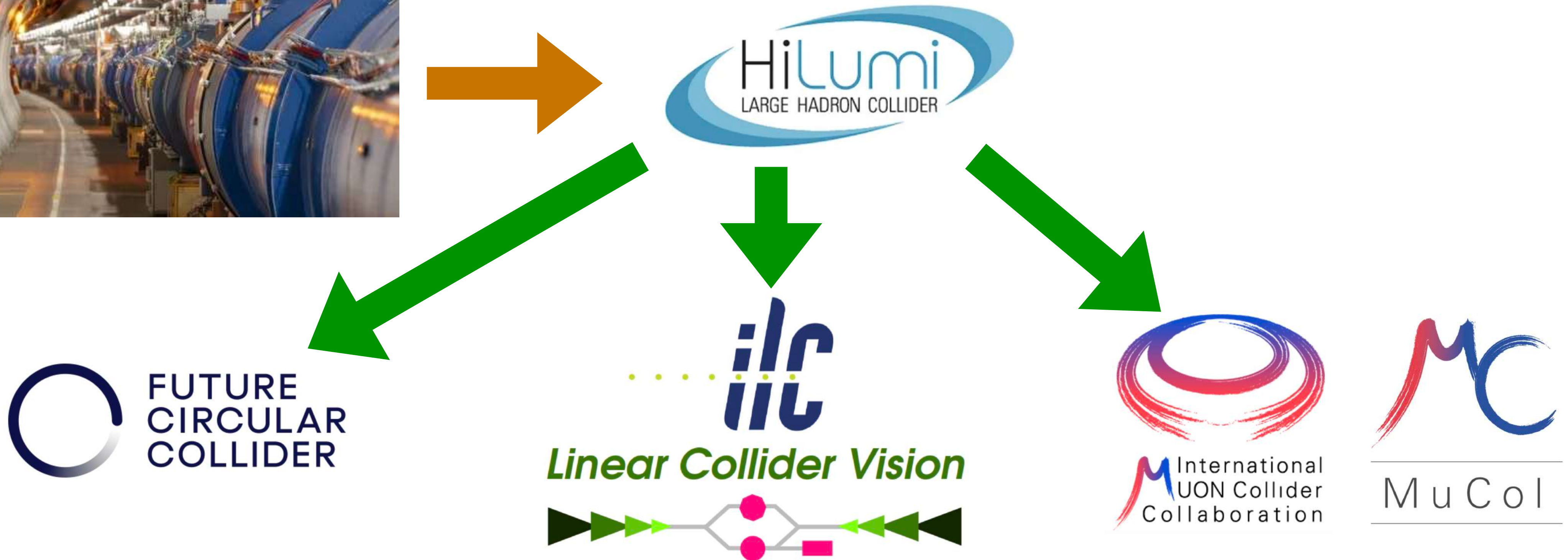
[ATLAS and CMS '25]



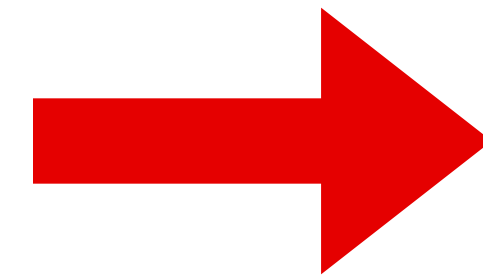
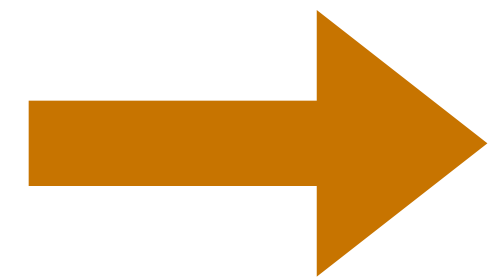
Motivation - Photon collider



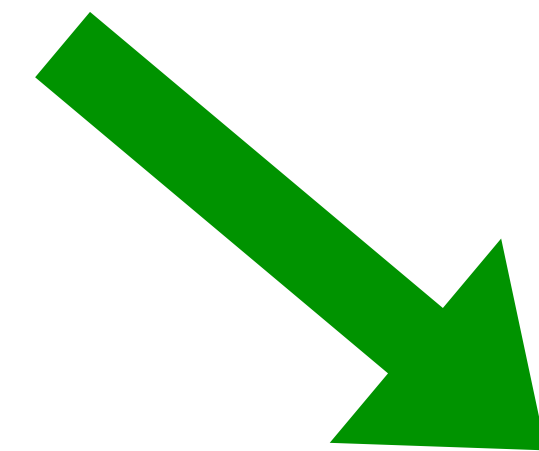
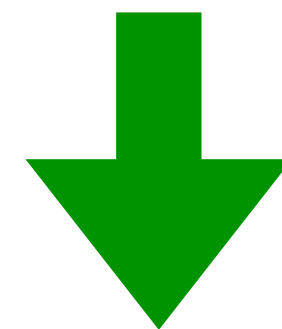
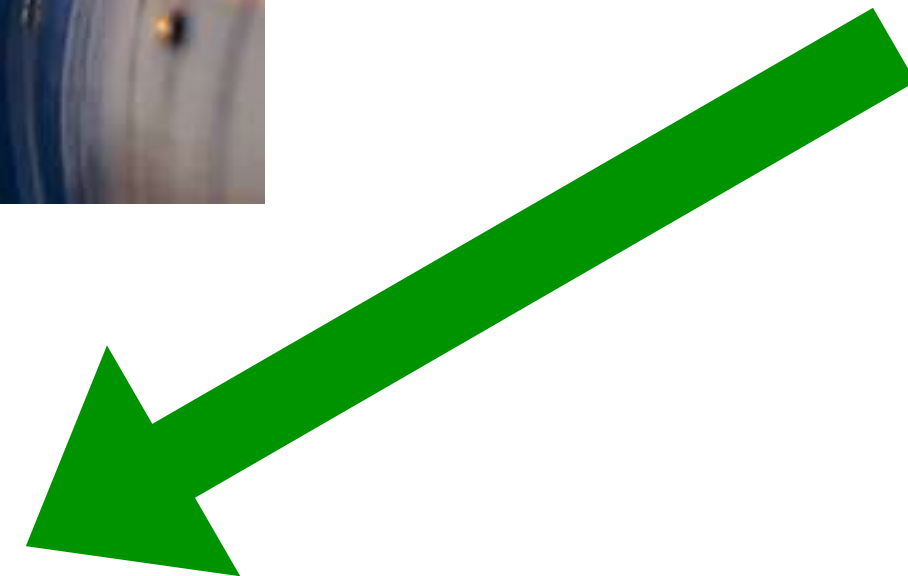
Motivation - Photon collider



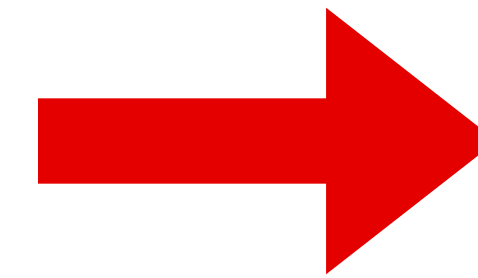
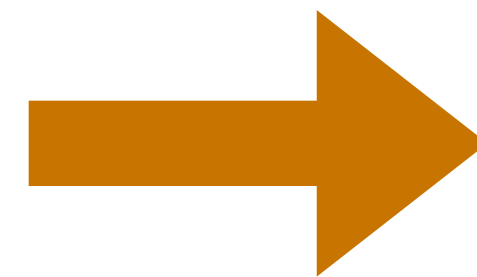
Motivation - Photon collider



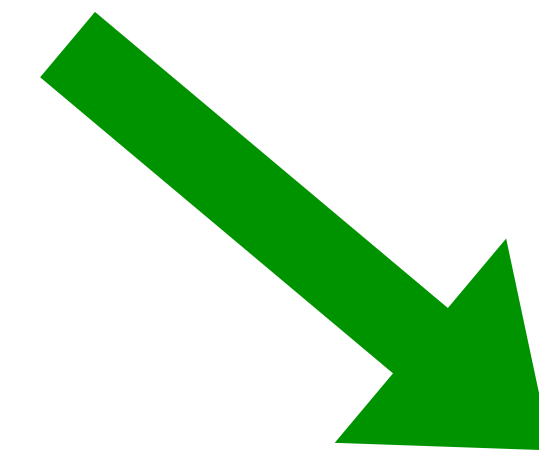
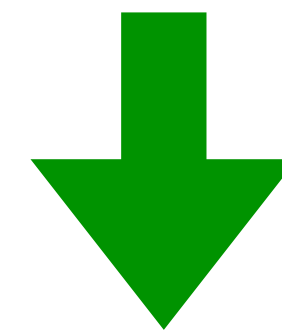
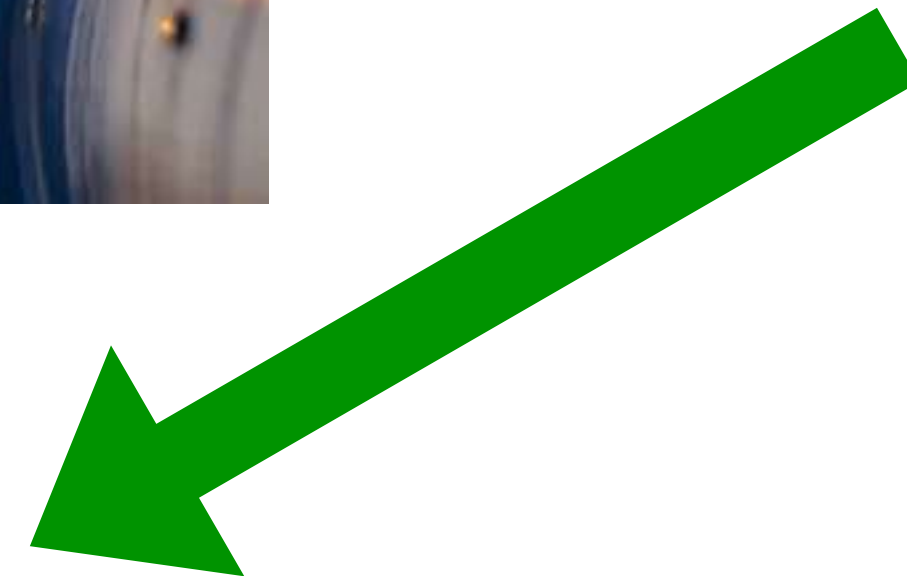
Photon collider



Motivation - Photon collider



Photon collider



- Measuring the trilinear Higgs coupling is one of the goals of all future colliders

$$gg \rightarrow hh, \quad e^+e^- \rightarrow Zhh, \quad e^+e^- \rightarrow \nu_e\bar{\nu}_e hh, \quad \gamma\gamma \rightarrow hh$$

Why photon colliders?

- Photon physics is complementary to other collider searches
- Different channels for searches
- High cross-sections at lower centre-of-mass energy

- Especially for Higgs physics

- $\gamma\gamma \rightarrow h$

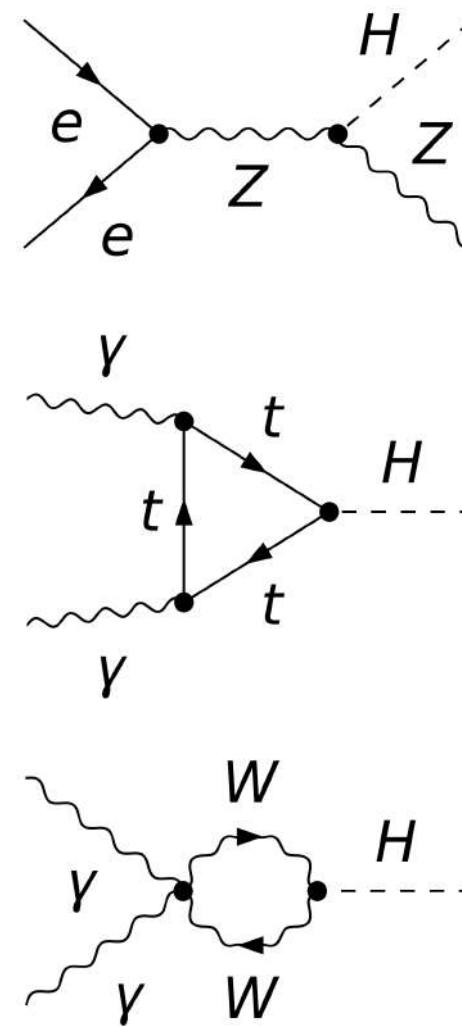
[U.S. Qureshi, T. Barklow, A. Schwartzman '26]

- And SUSY, production of charged particles

- $e\gamma \rightarrow \tilde{e}\tilde{\chi}_1^0$

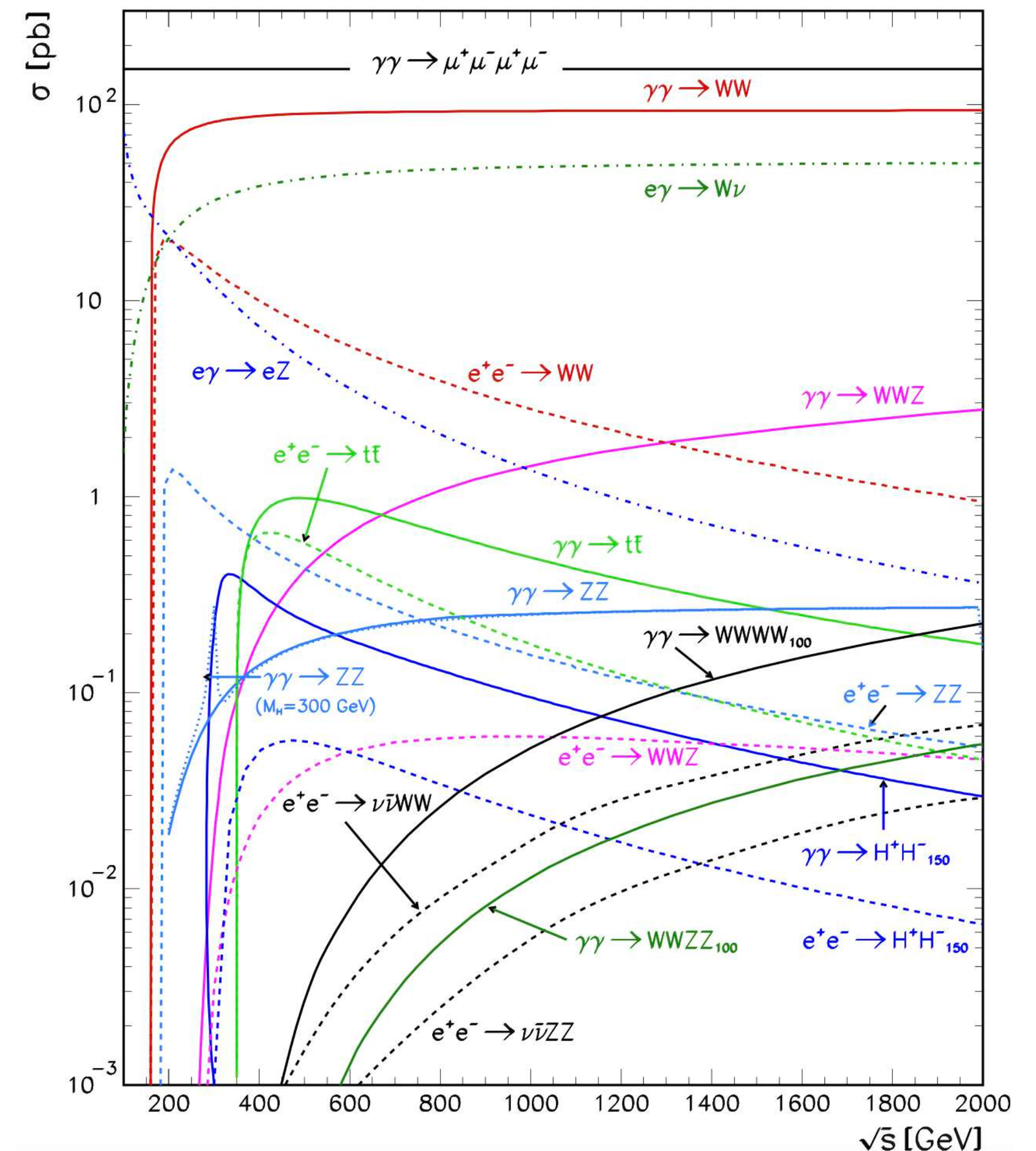
[A. Goto and T. Kon '92][F. Cuypersa, G. J. van Oldenborgha, R. Rückl '92]

[S. Kanemura '01][M. Mühlleitner et al. '06]



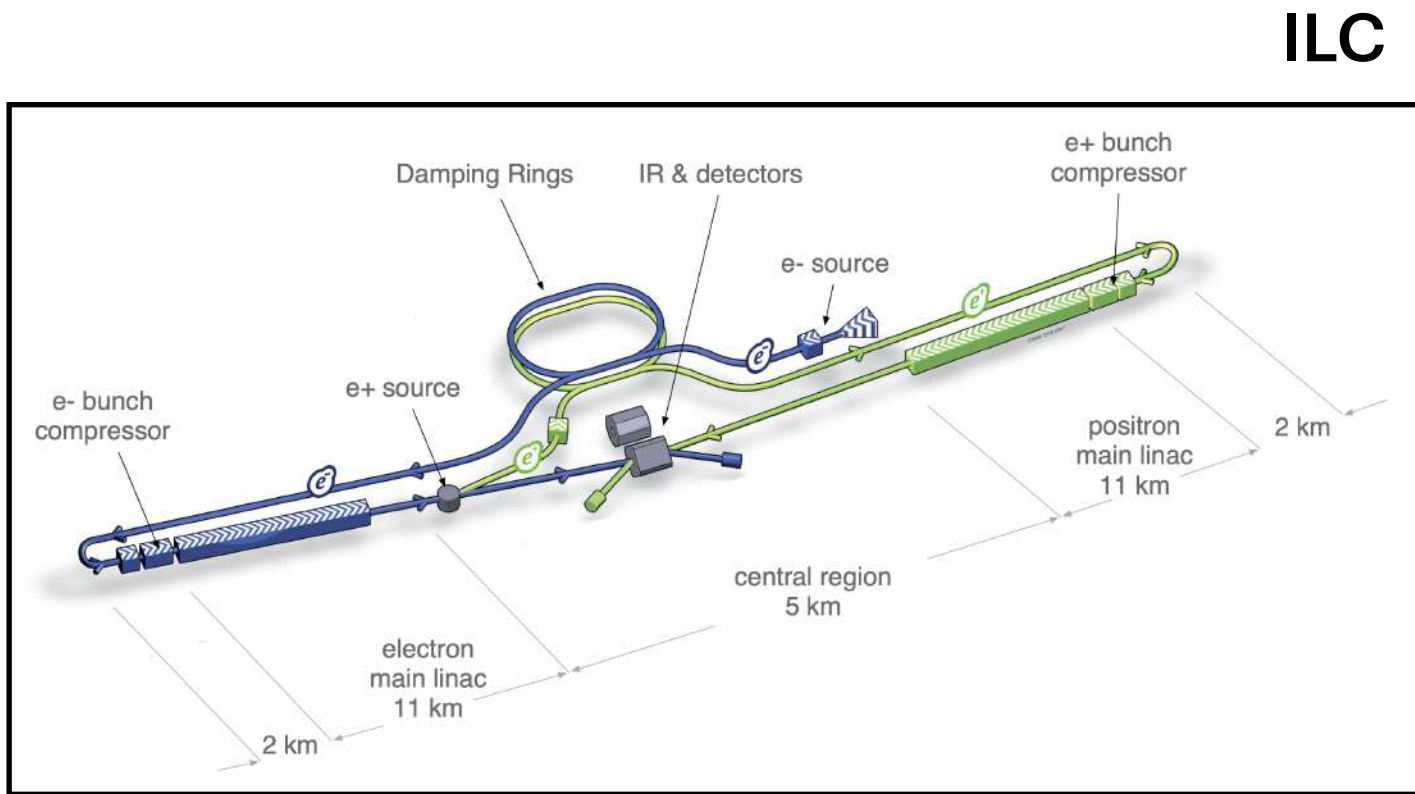
[H. Abramowicz, MB et al. '25]

[B. Badelek et al. '04]

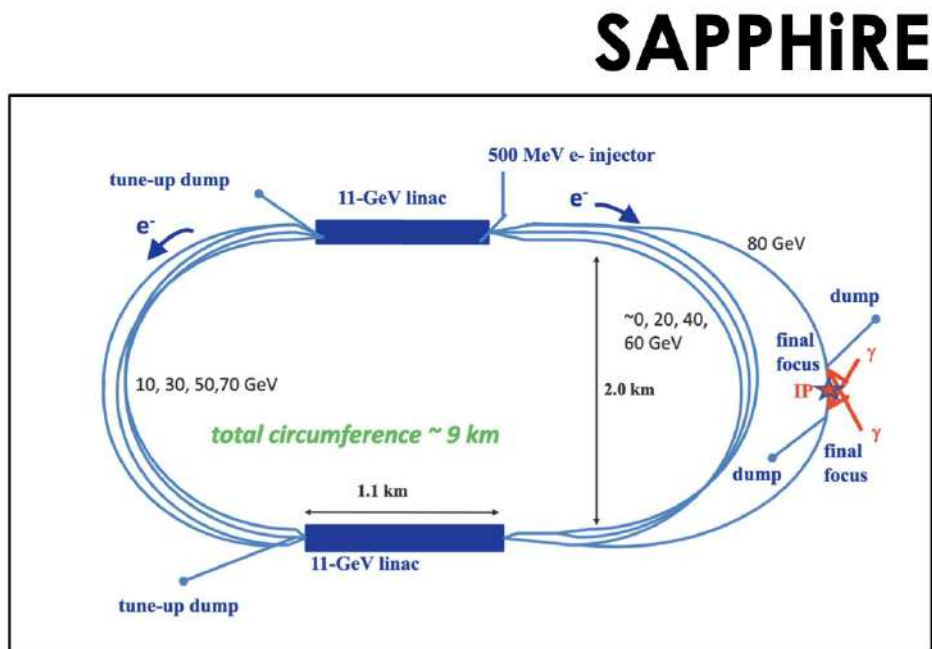


Photon collider designs

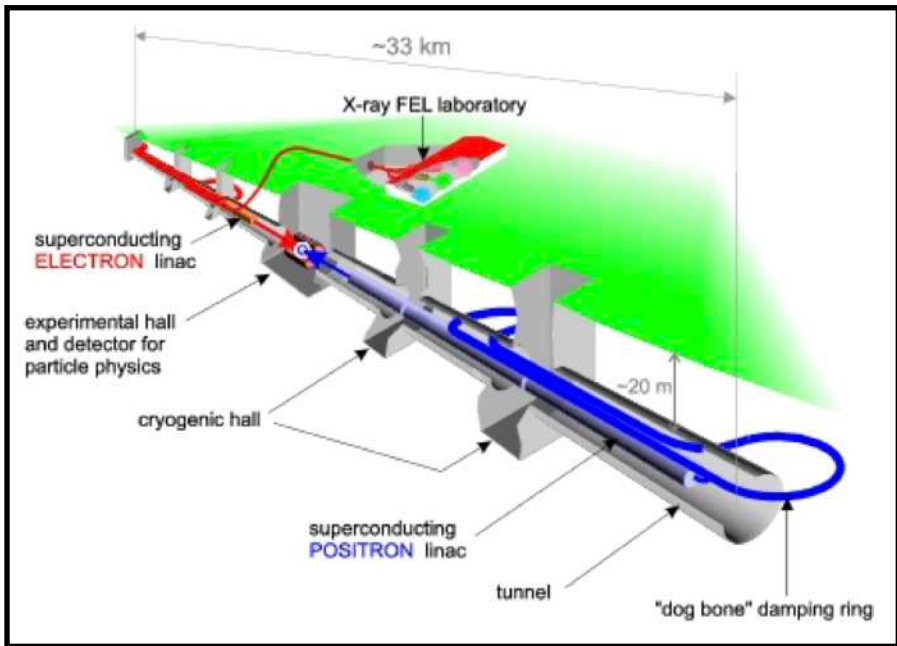
- Many future linear collider designs have been considered for photon collider upgrades
- Expanding the physics reach at low center-of-mass energies



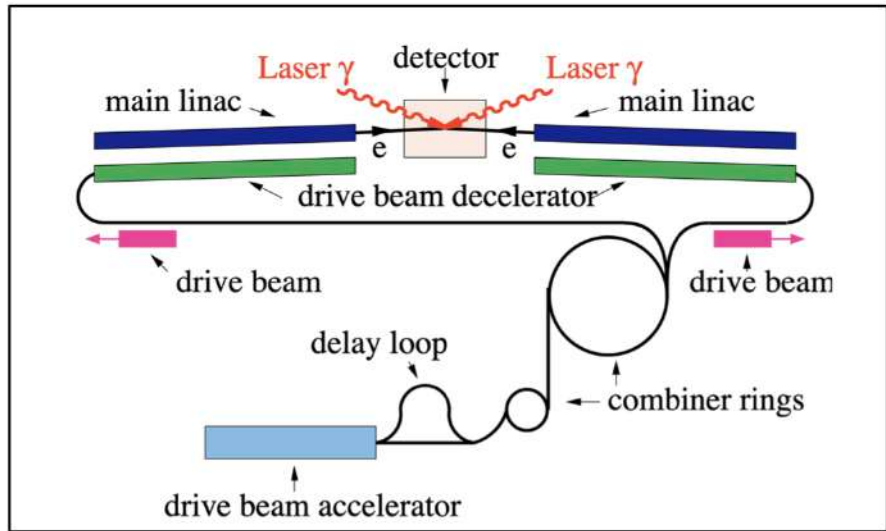
TESLA



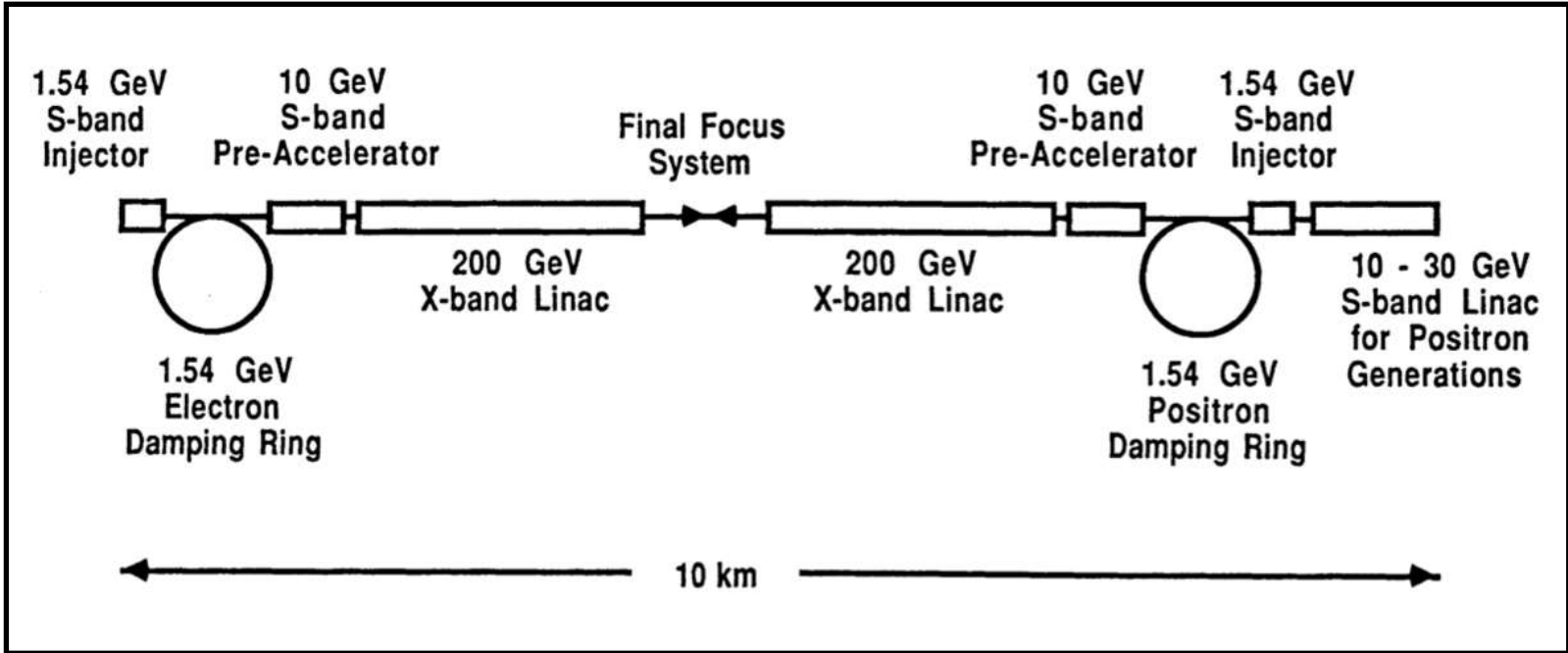
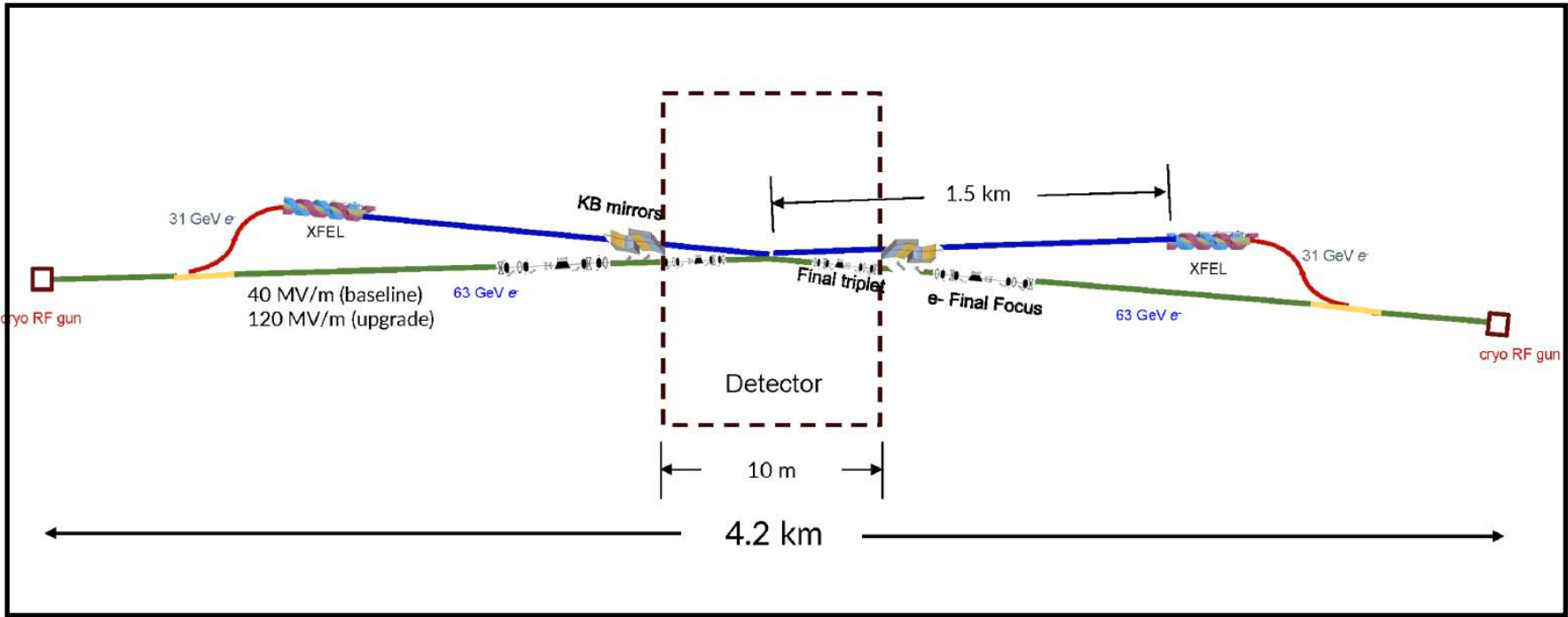
SAPHIRE



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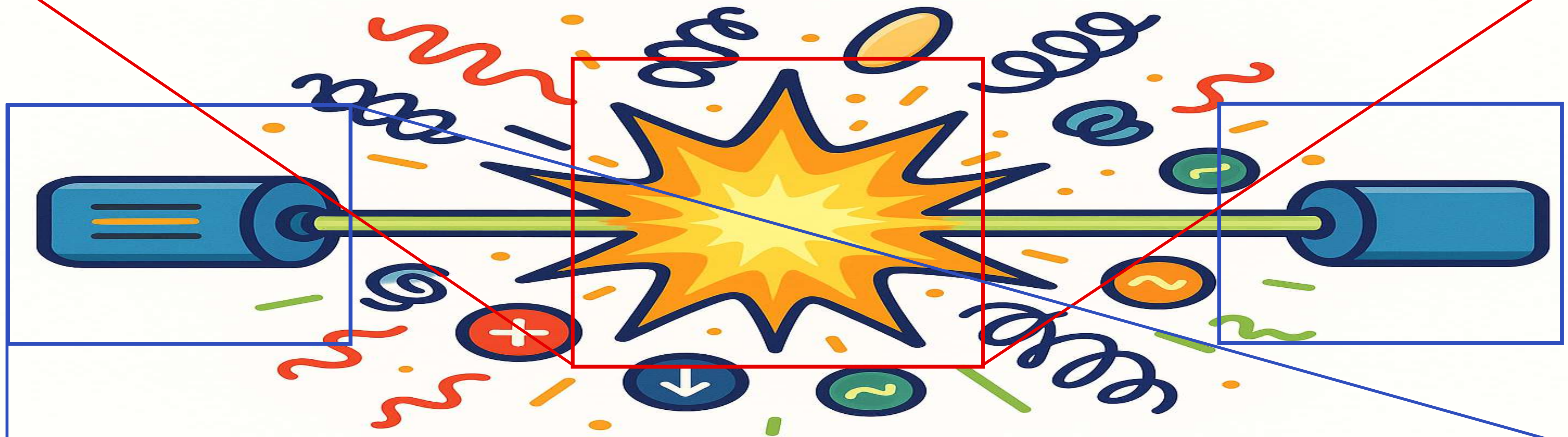


XCC



JLC

Study of $\gamma\gamma$ -processes



How do we produce high-energy photons?

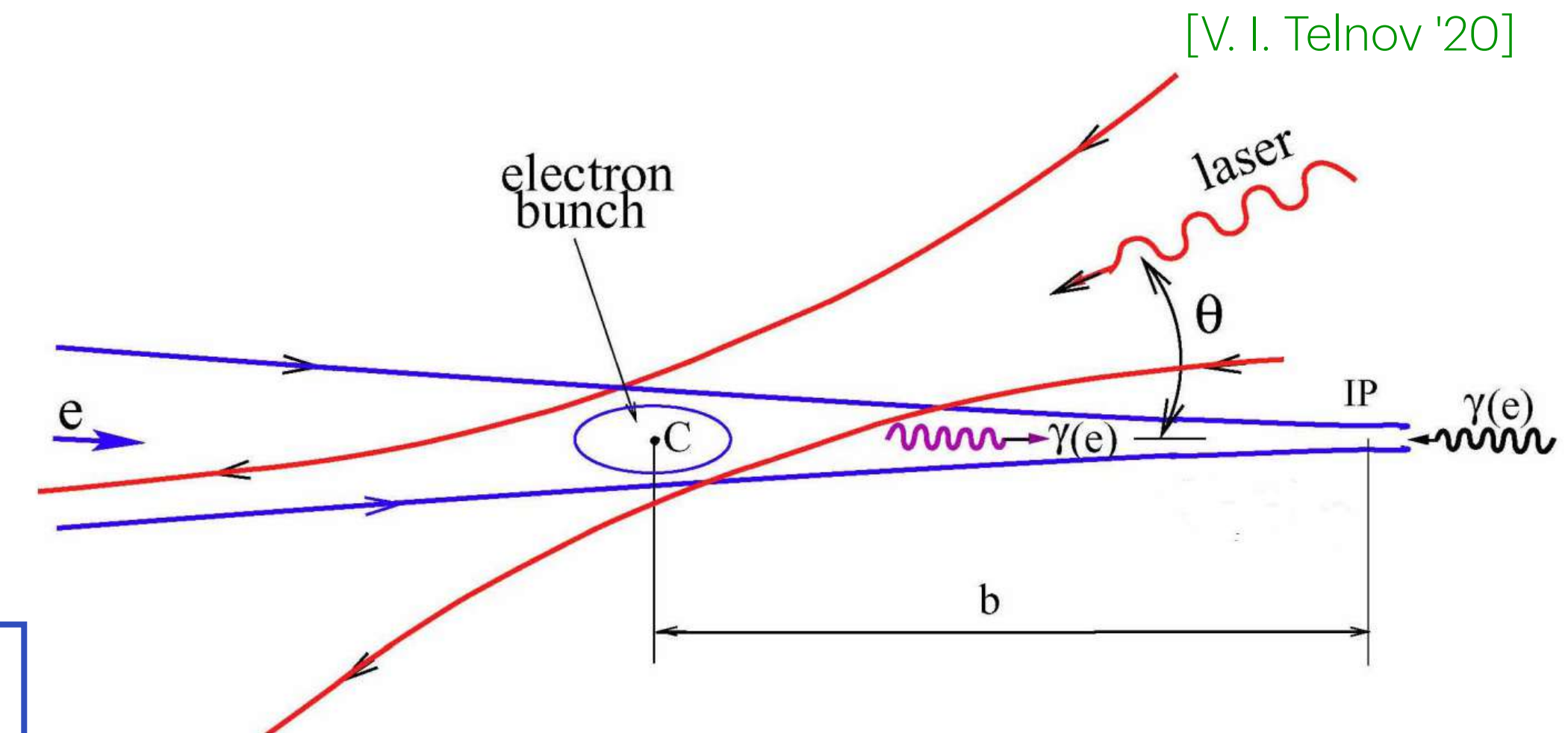
Study of $\gamma\gamma$ -processes



How do we produce high-energy photons?

Principles of a photon collider

- Simple laser collider is too weak (eV to $\sim$ MeV)
- Compton scattering at ee -collider
- Getting access to $\gamma\gamma$ and γe processes



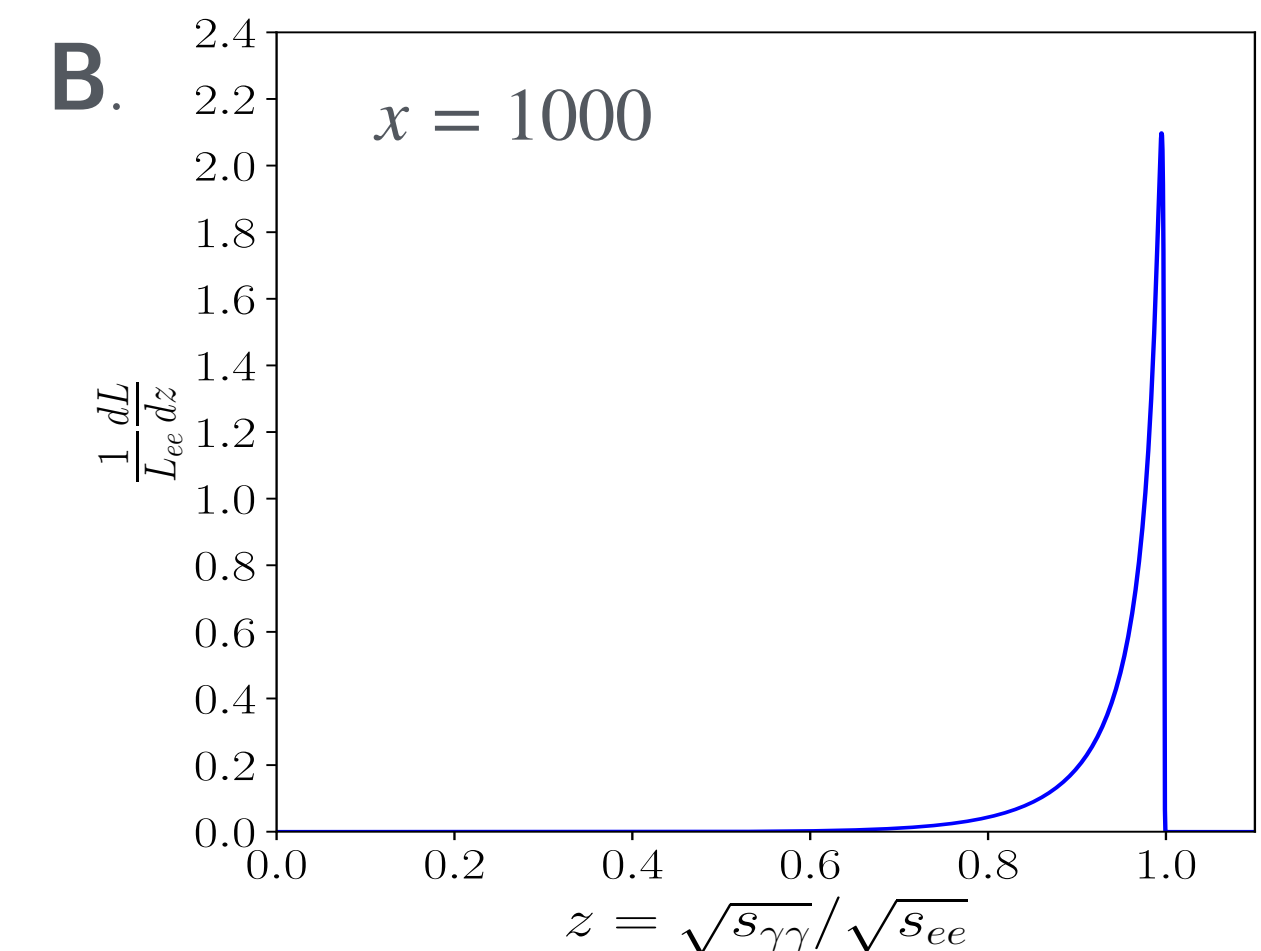
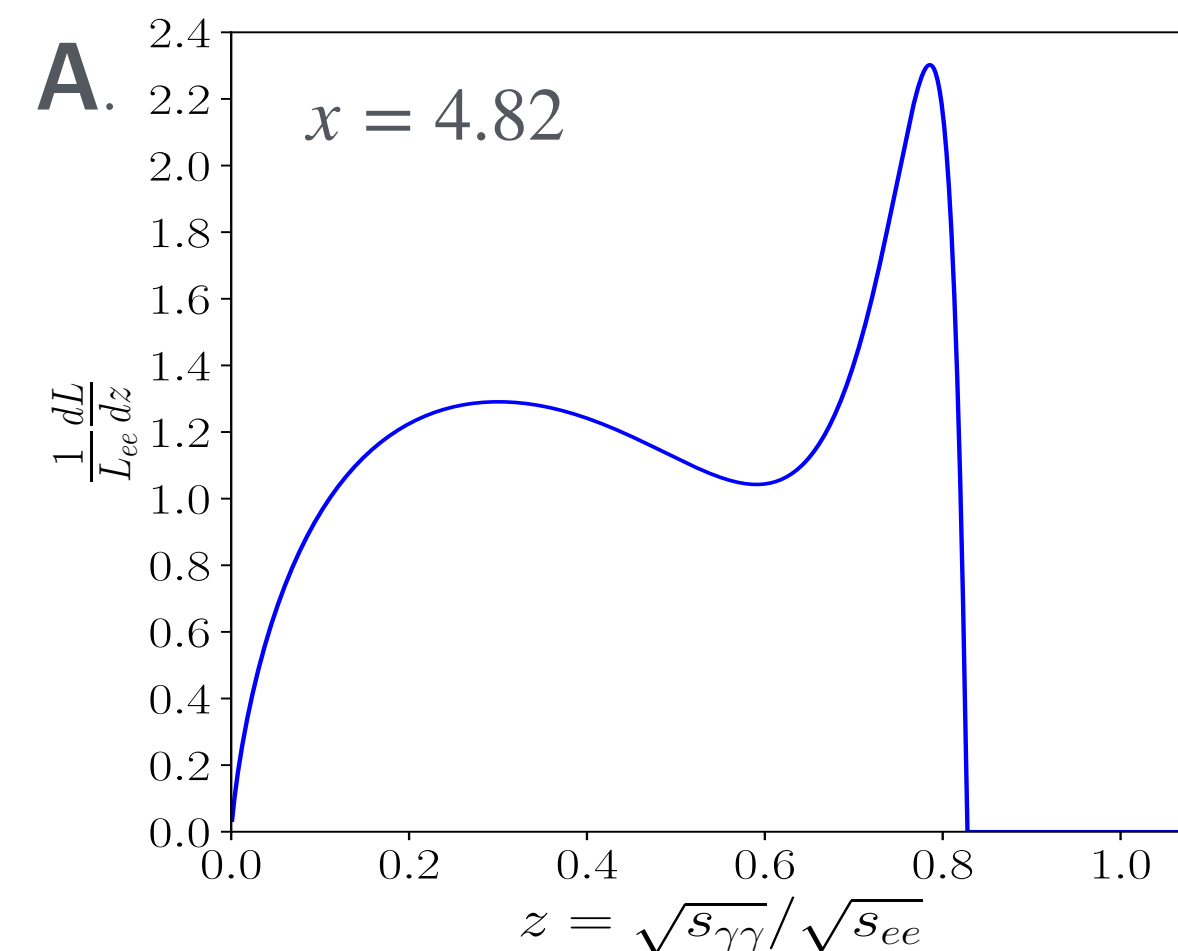
$$E_{\gamma}^{\max} \approx \frac{x}{1+x} E_e, \quad x = \frac{4E_e \omega_0}{m_e^2} \simeq 15.3 \left[\frac{E_e}{\text{TeV}} \right] \left[\frac{\omega_0}{\text{eV}} \right] = 19 \left[\frac{E_e}{\text{TeV}} \right] \left[\frac{\mu\text{m}}{\lambda} \right]$$

Type of laser is decisive

A. Optical $x \lesssim 4.82$

[T. Barklow et al. '21]

B. XFEL-like (XCC) $x \geq 1000$



Polarisation effects

- The total luminosity depends on the product of the electron polarisation λ_e and laser photon polarisation P_c : $\frac{1}{k^2 L_{ee}} \frac{dL_{\gamma\gamma}}{dz} = F(x, y, \lambda_e P_c, \rho)$
- But the scattered photon polarisation also depends on the incoming beam polarisation individually

$$\langle \lambda_\gamma \rangle = \frac{\lambda_e x r [1 + (1 - y)(2r - 1)^2] - P_c (2r - 1) [(1 - y)^{-1} + 1 - y]}{(1 - y)^{-1} + 1 - y - 4r(1 - r) - \lambda_e P_c x r (2 - y)(2r - 1)}$$

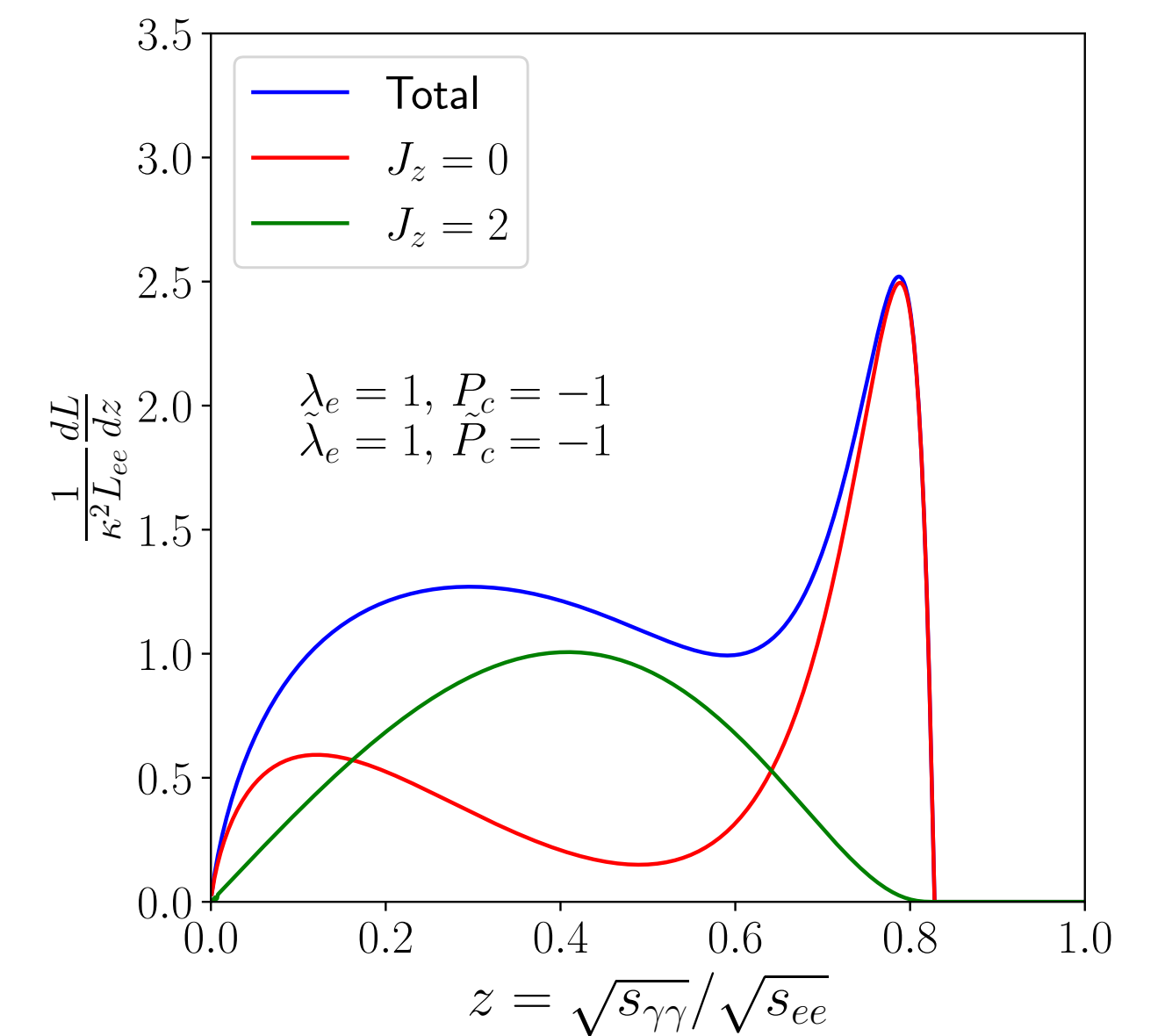
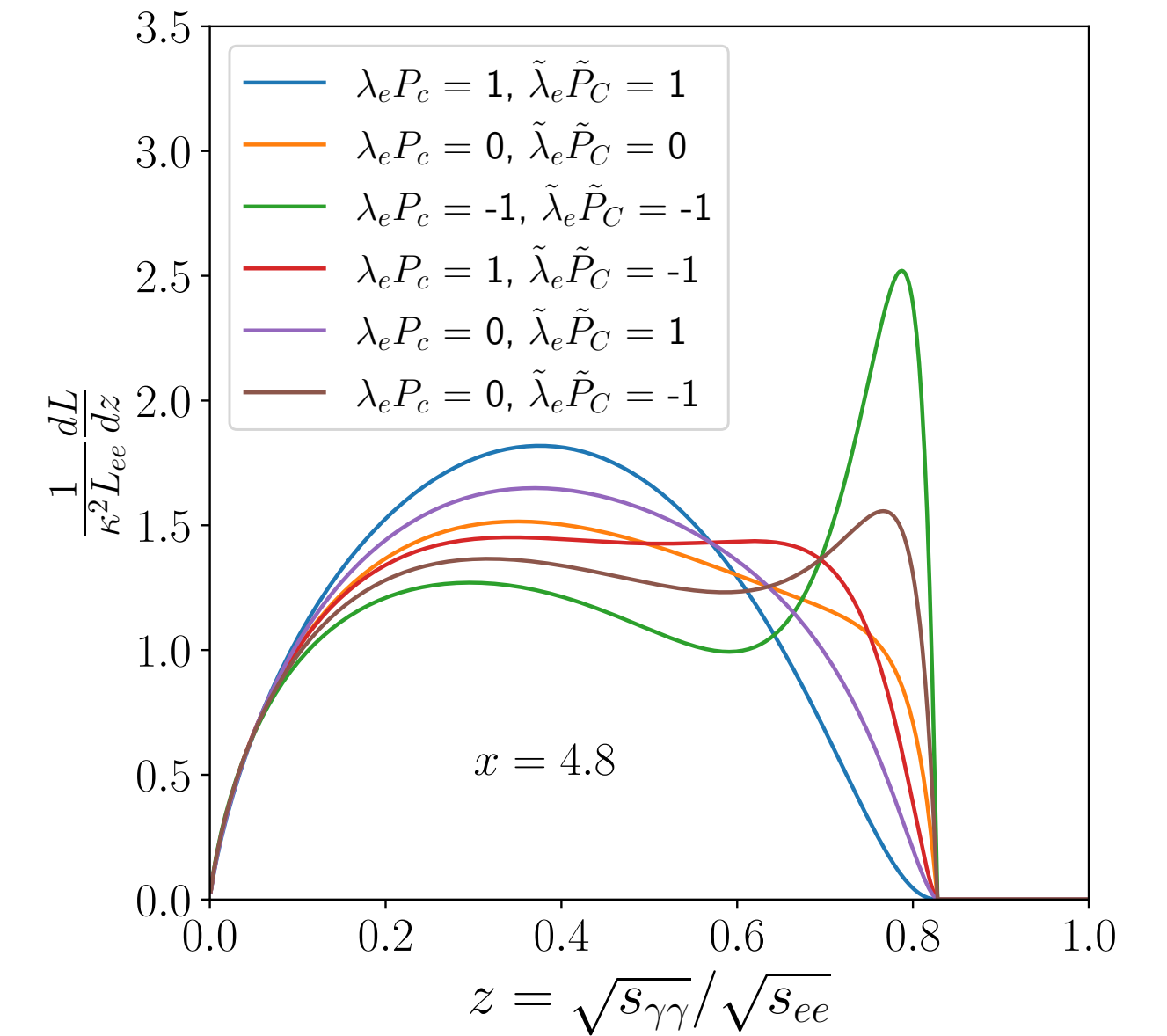
- Two distinct total angular momentum states can be defined

$$J_z = 0$$

Photons with: equal polarisation

$$J_z = 2$$

opposite polarisation



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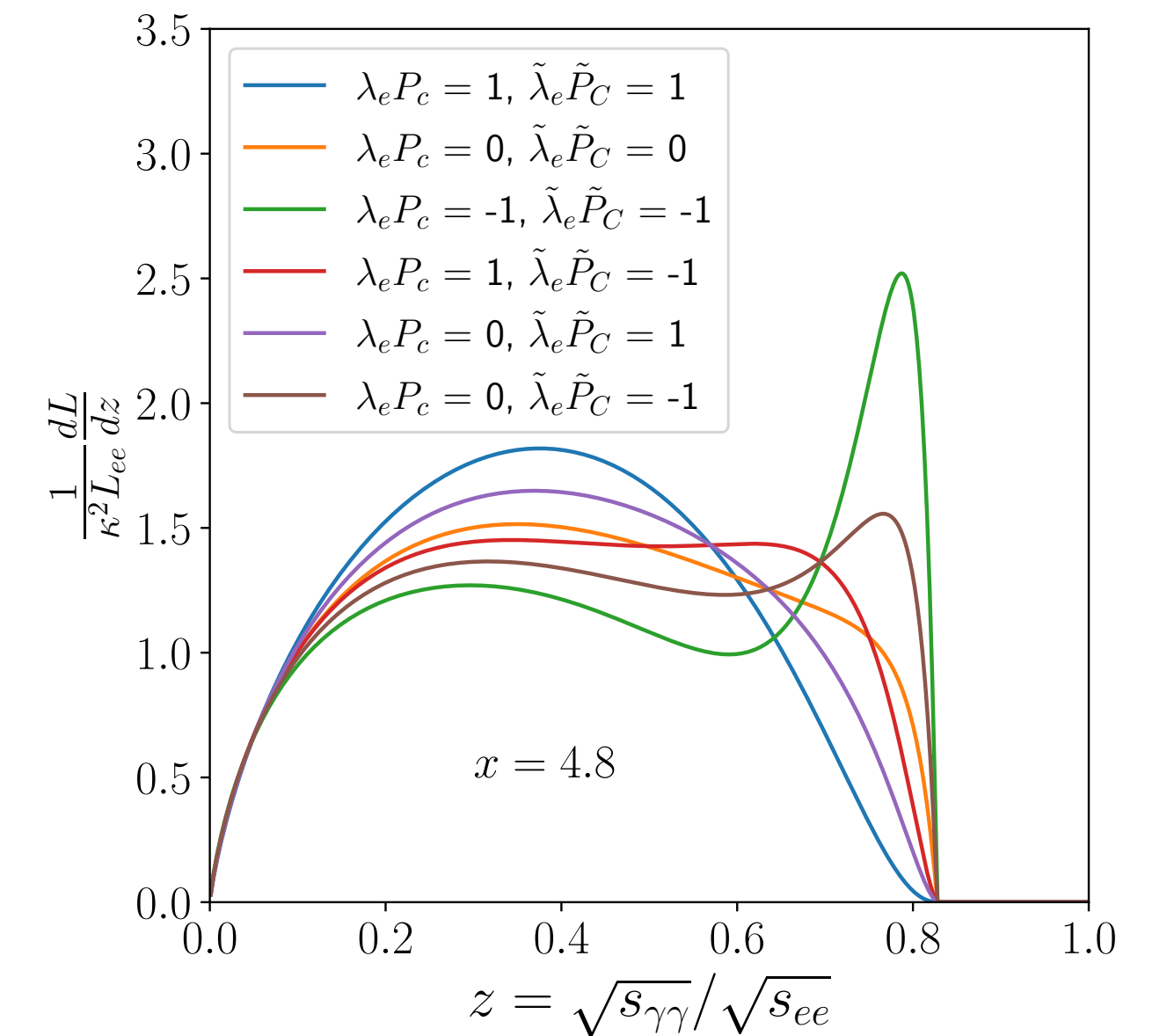
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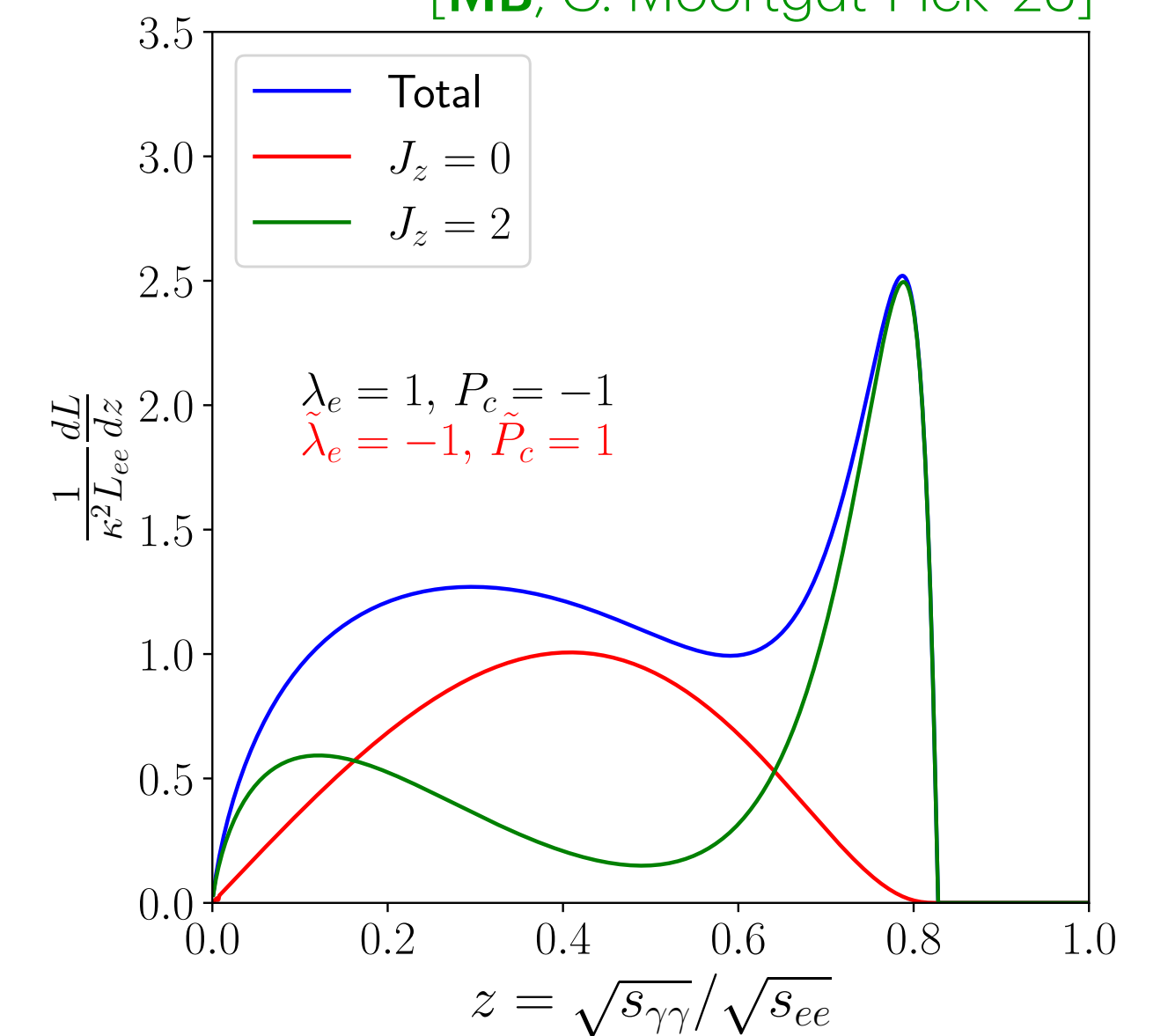
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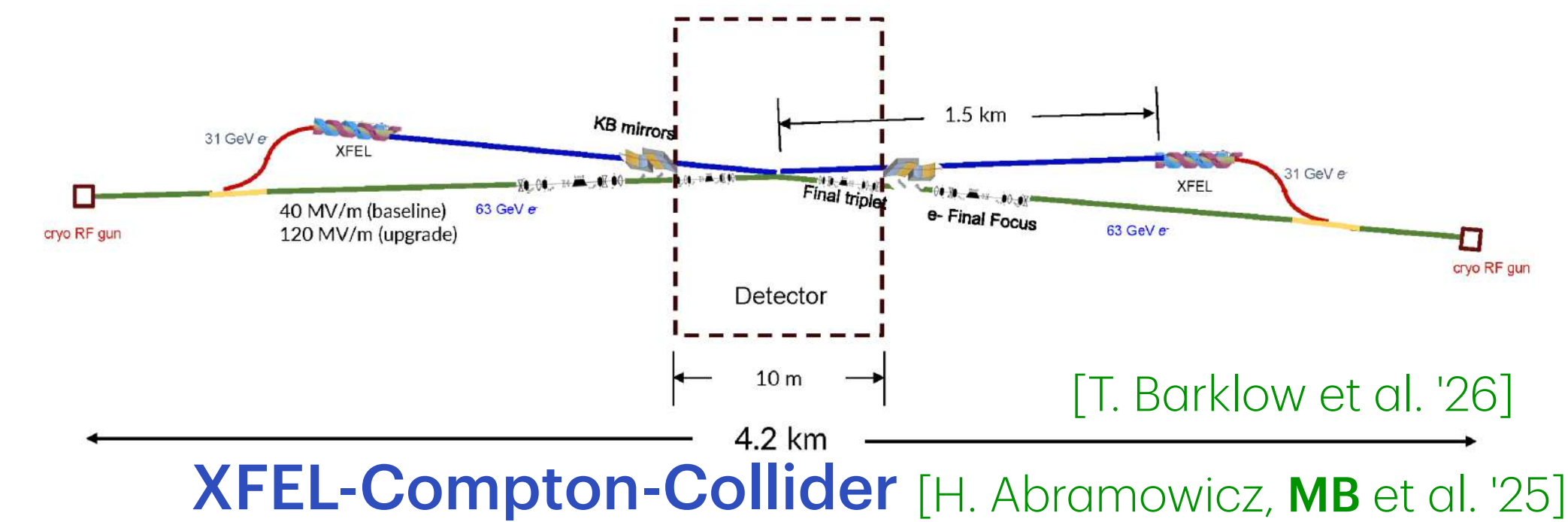
[MB, G. Moortgat-Pick '26]



Photon collider designs

Optical collider design

- Mercury laser system or Nd:YAG laser
- $\lambda = 1047 \text{ nm}$ or $\lambda = 1064 \text{ nm}$
- $x \lesssim 4.82$
- $\lambda_e P_c = -0.9$
- $\sqrt{s_{\gamma\gamma}^{\max}} \approx 0.8 \sqrt{s_{ee}}$
- Two photon collider energies: $\sqrt{s_{\gamma\gamma}} \approx 270 \text{ GeV}$ and $\sqrt{s_{\gamma\gamma}} \approx 380 \text{ GeV}$
- Electron collider energies:
 - $\sqrt{s_{ee}} = 380 \text{ GeV}$
 - $\sqrt{s_{ee}} = 550 \text{ GeV}$



- XFEL photon source
- $\lambda \approx 2 \text{ nm}$
- $x \geq 1000$
- $\lambda_e P_c = 0.9$
- $\sqrt{s_{\gamma\gamma}^{\max}} = \sqrt{s_{ee}}$
- Two photon collider energies: $\sqrt{s_{\gamma\gamma}} = 270 \text{ GeV}$ and $\sqrt{s_{\gamma\gamma}} = 380 \text{ GeV}$
- Electron collider energies:
 - $\sqrt{s_{ee}} = 270 \text{ GeV}$
 - $\sqrt{s_{ee}} = 380 \text{ GeV}$

Photon collider designs

- We used the Monte-Carlo code CAIN [P. Chen et al. '95] to simulate realistic photon collider spectra
- The optical collider designs are based on ILC photon collider design [F. Bechtle et al. '06] and LCF at CERN [H. Abramowicz, MB et al. '25]
- The XCC designs are based on C³ [K.L. Bane et al. '18][T. Barklow et al. '22] but have also been checked using a SCRF accelerator [H. Abramowicz, MB et al. '25]

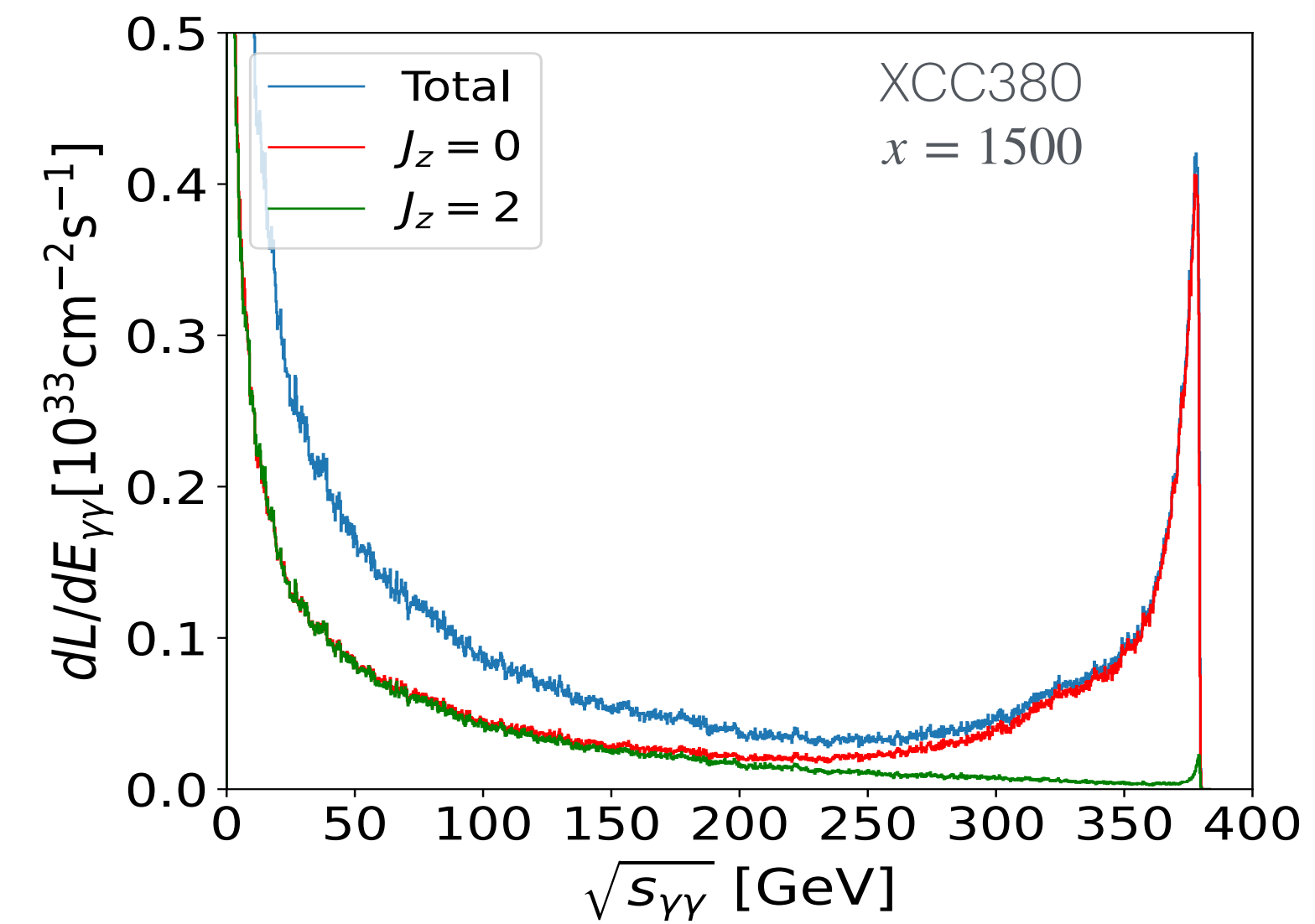
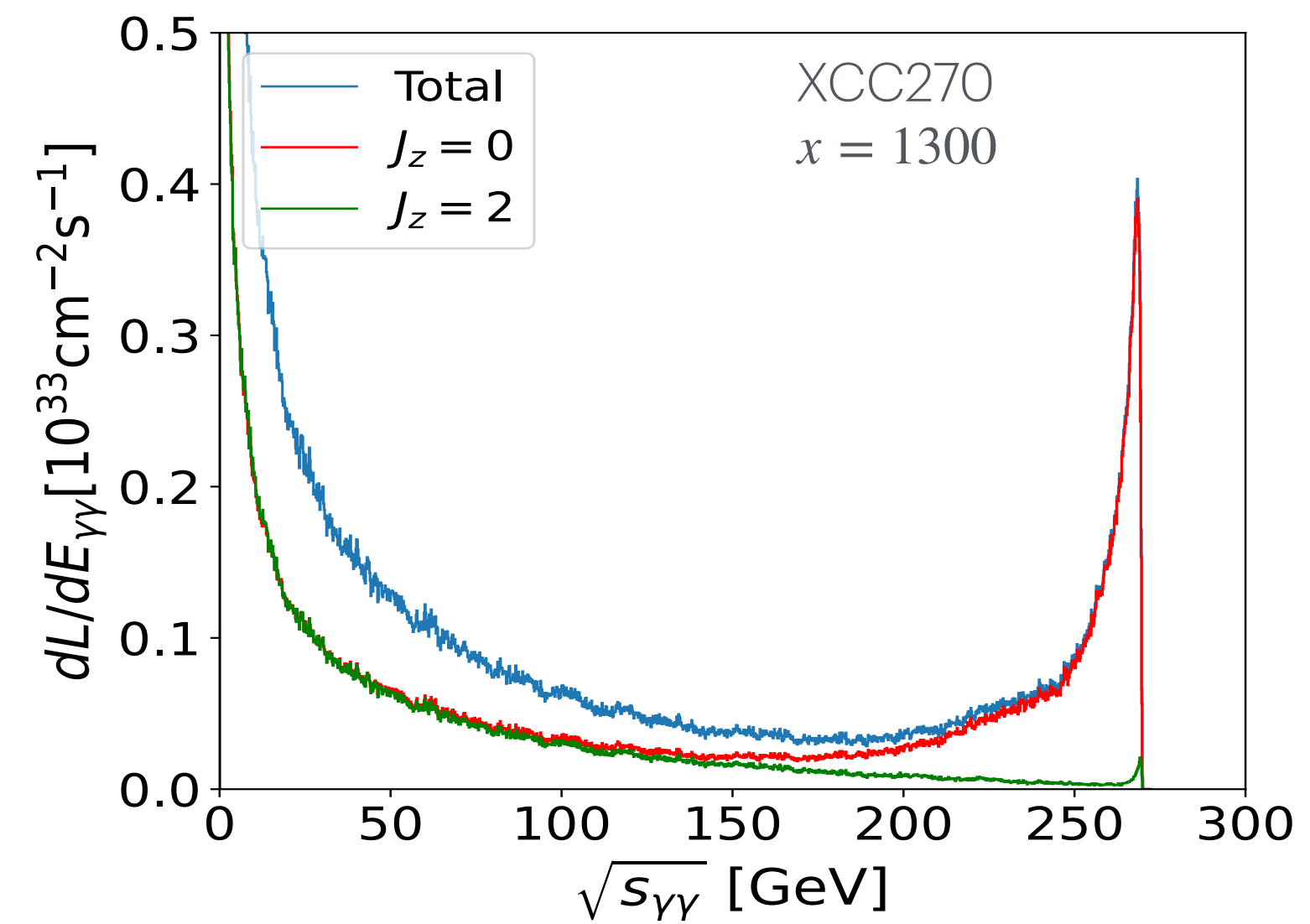
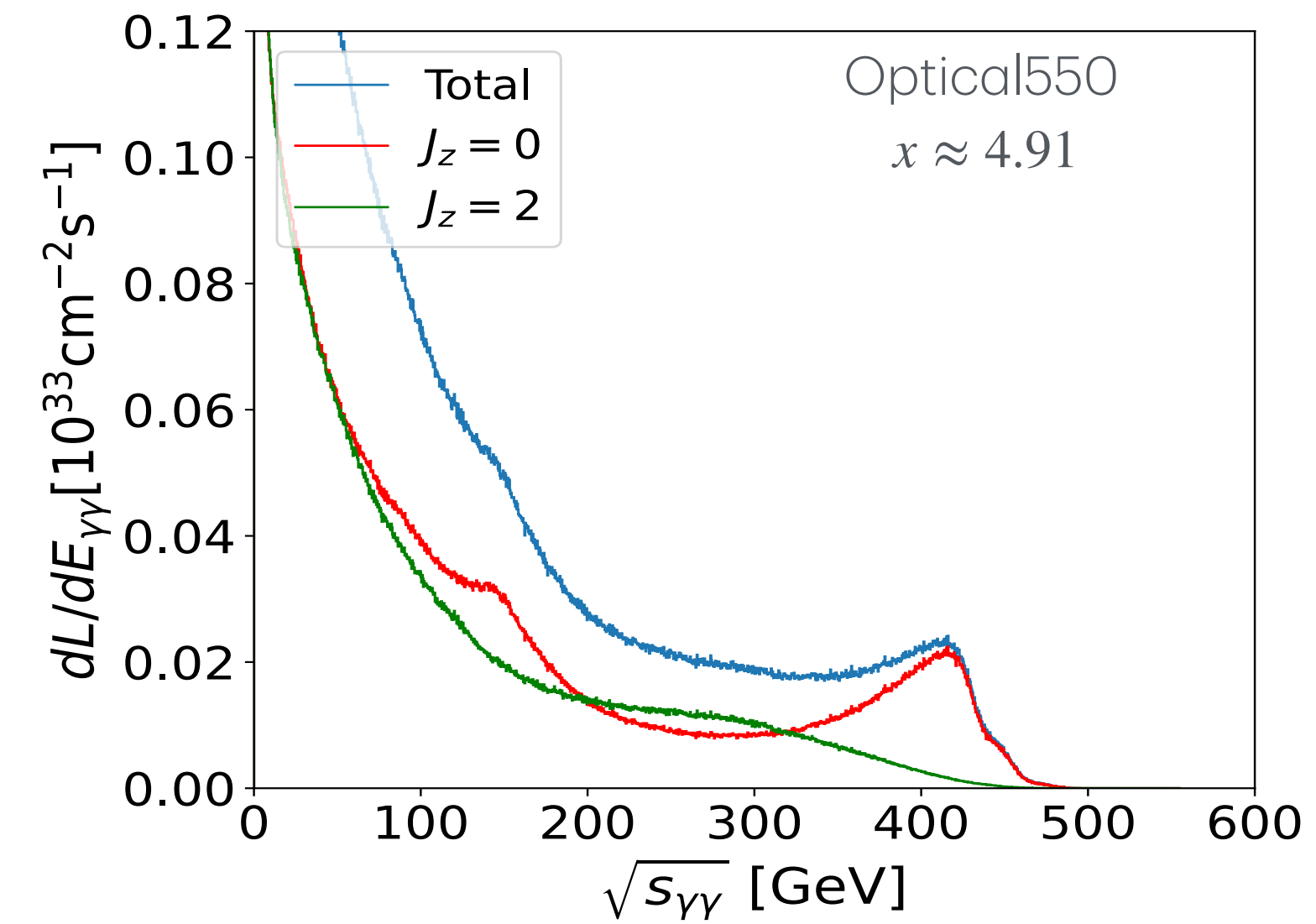
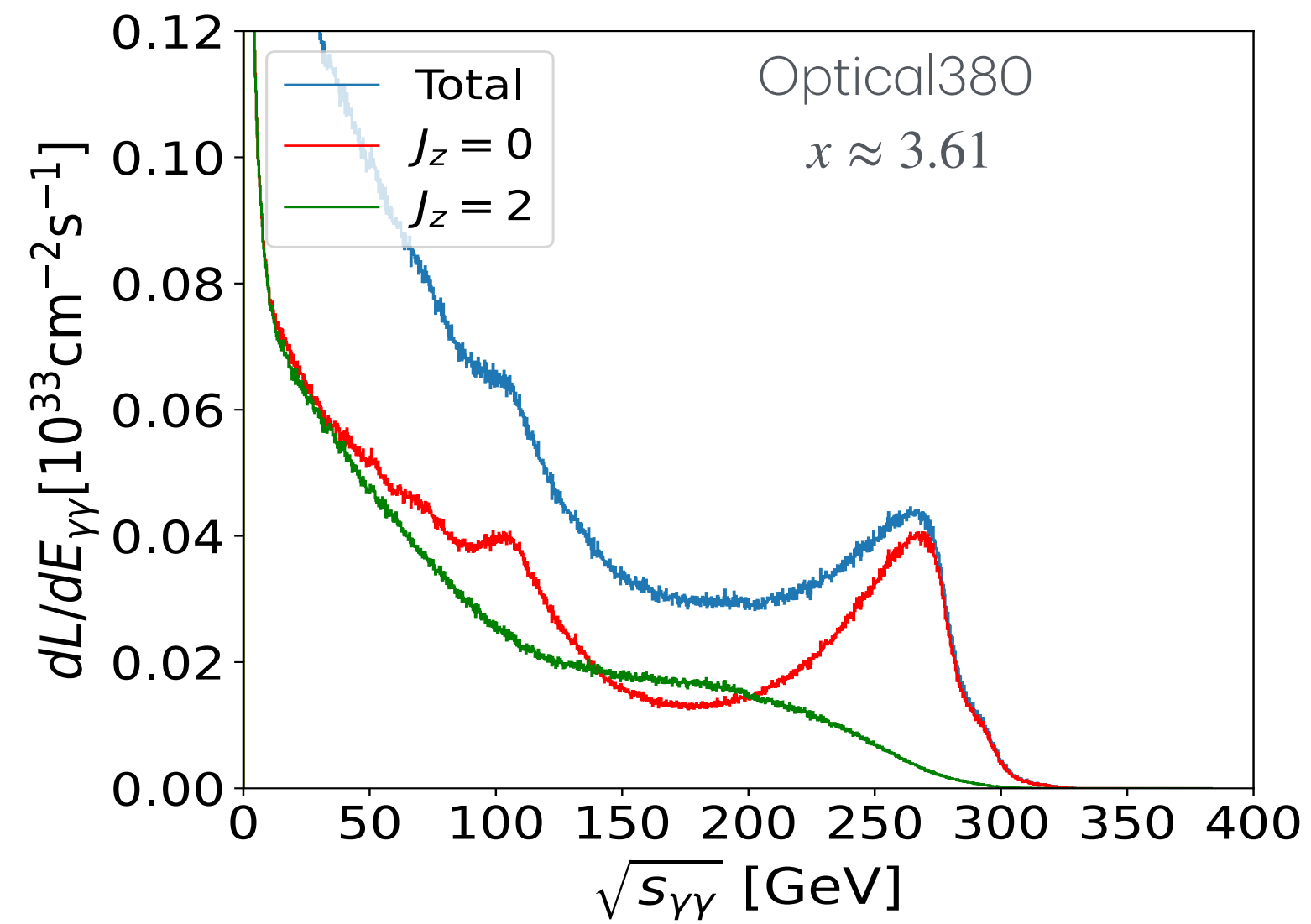
Optical collider

Parameters	$\sqrt{s_{\gamma\gamma}} = 270 \text{ GeV}$	$\sqrt{s_{\gamma\gamma}} = 420 \text{ GeV}$
Electron energy E_0 [GeV]	190	275
β_x/β_y [μm]	1500/300	11000/482
$\gamma\varepsilon_x/\gamma\varepsilon_y$ [nm·mrad]	2500/30	10000/35
σ_x/σ_y [nm]	88/4.3	452/5.6
σ_z [μm]	300	300
Bunch charge [$10^{10} e^-$]	2	2
Collision rate [kHz]	14.2	26.25
Crossing angle [mrad]	20	20
b (<i>C-IP</i> distance) [mm]	2.6	2.6
Laser λ (energy) [μm (eV)]	1 (1.17)	1 (1.17)
Laser flash energy [J]	9	9
Laser pulse duration [ps]	1.5	1.5
f# $\equiv F/D$ of laser system	17	17
Parameter x and ξ^2	3.61, 0.34	4.91, 0.39
Power density laser [$10^{21}\text{W}/\text{m}^2$]	9.3	9.3
$\alpha_{e\lambda}$ [mrad]	47	47

XCC

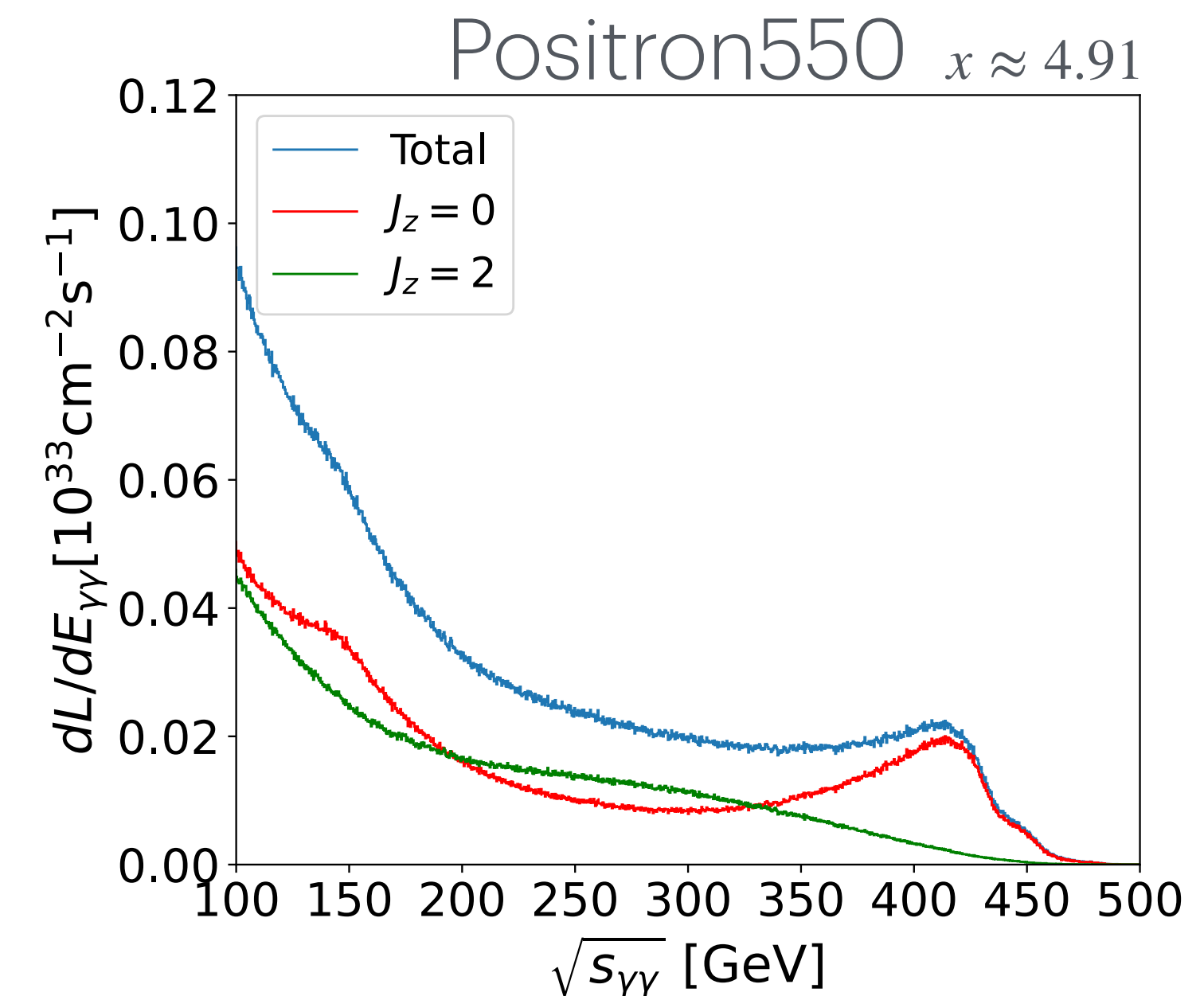
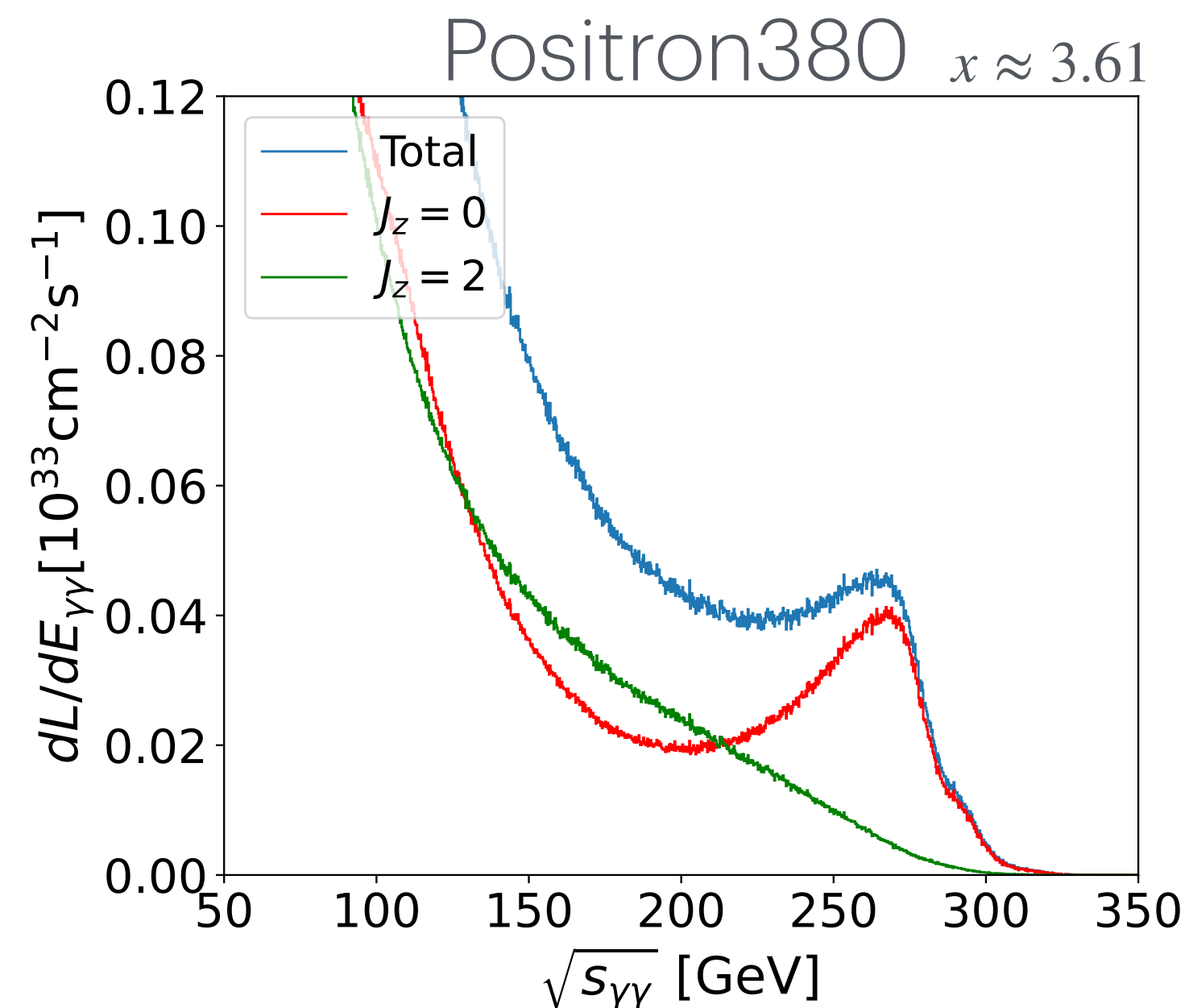
Parameters	$\sqrt{s_{\gamma\gamma}} = 270 \text{ GeV}$	$\sqrt{s_{\gamma\gamma}} = 380 \text{ GeV}$
Electron energy E_0 [GeV]	135	190
β_x/β_y [μm]	10/10	10/10
$\gamma\varepsilon_x/\gamma\varepsilon_y$ [nm·mrad]	60/60	60/60
σ_x/σ_y [nm]	1.5/1.5	1.5/1.5
σ_z [μm]	10	10
Bunch charge [$10^{10} e^-$]	0.62	0.62
Collision rate [kHz]	14.52	11.16
Crossing angle [mrad]	2	2
b (<i>C-IP</i> distance) [mm]	0.04	0.04
Laser λ (energy) [nm (keV)]	2.0 (0.61)	2.4 (0.52)
Laser flash energy [J]	1	1
Laser pulse length [fs]	33	67
Parameter x and ξ^2	1300, 0.76	1500, 1.13
Power density laser [$10^{27}\text{W}/\text{m}^2$]	5.3	5.3
$\alpha_{e\lambda}$ [mrad]	1	1

Photon collider designs



Photon collider designs - e^+e^- base

- Positron beams have not been considered in photon collider designs
 - Polarisation is very important, now can reach 60% polarised e^+ beams [S. Riemann, F. Staufenberg, G. Moortgat-Pick, A. Ushakov '14] [J.A. Clarke et al. '08] [K. Alharbi et al. '21]
- Same Compton scattering, but additional photon collisions at lower energies
- Designs work for optical collider setups, not tested for XCC



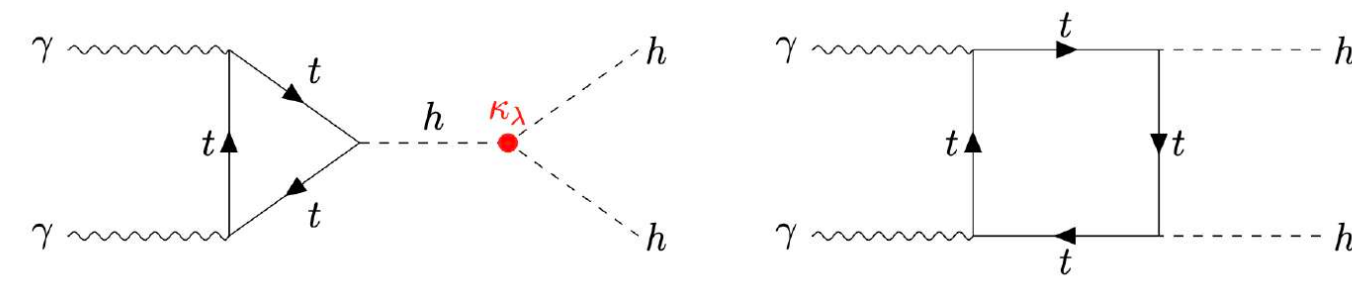
Study of $\gamma\gamma$ -processes



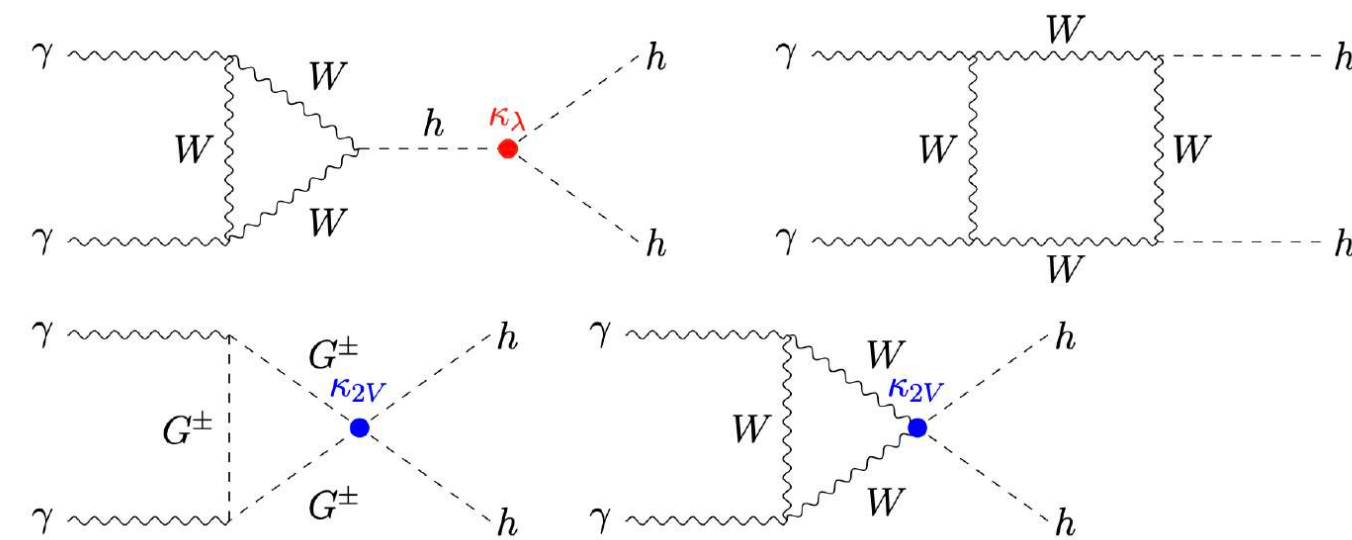
How do we produce high-energy photons?

Higgs pair production

- Similar structure to $gg \rightarrow hh$



- With additional gauge sector contributions



- Compute the helicity amplitudes for the process $\gamma\gamma \rightarrow hh$

$$\frac{d\hat{\sigma}(\lambda_1, \lambda_2)}{dt} = \frac{\alpha^2 \alpha_W^2}{32\pi \hat{s}^2} |M(\lambda_1, \lambda_2)|^2$$

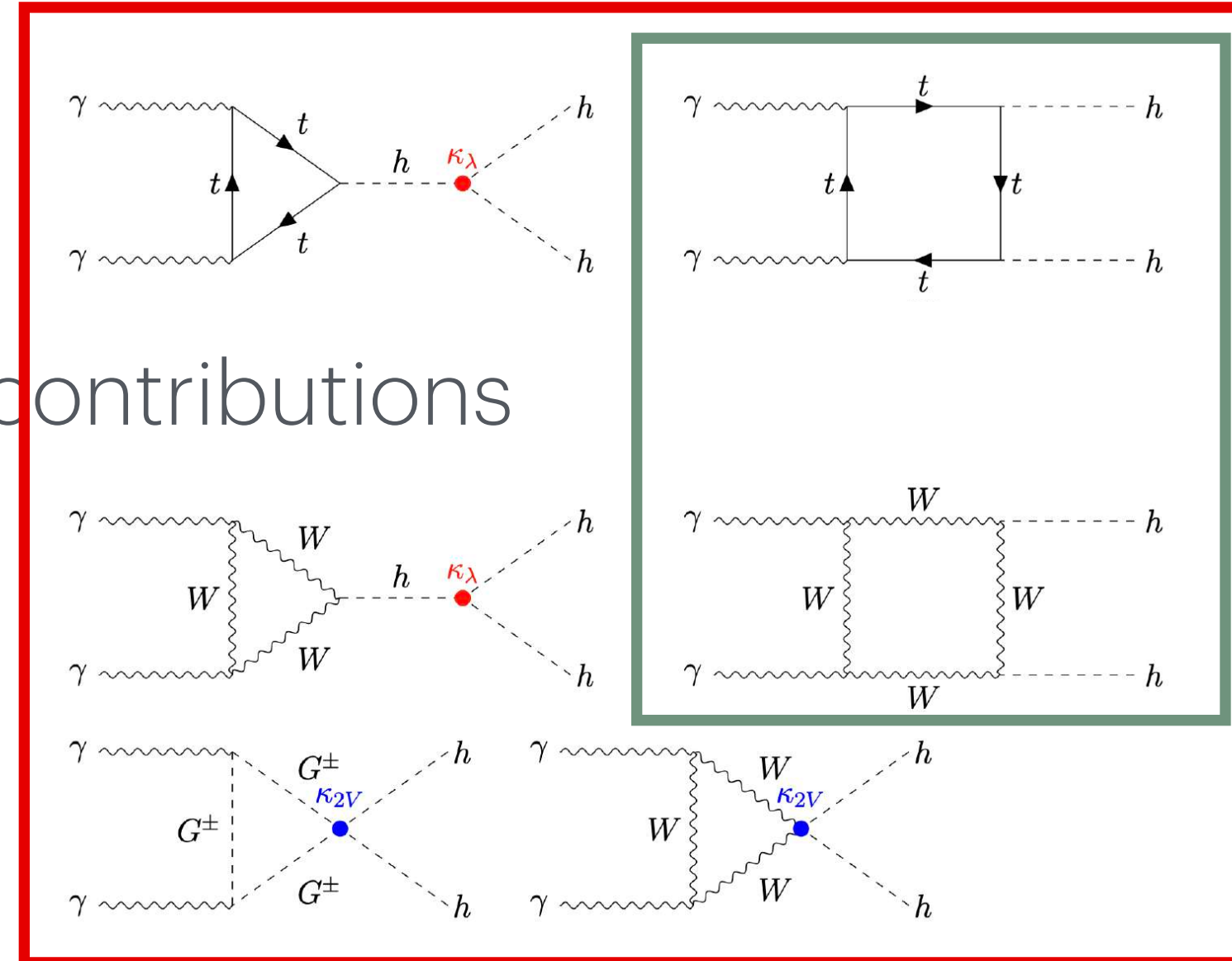
$$\kappa_i \equiv \frac{g_i}{g_i^{\text{SM},(0)}}$$

Higgs pair production

- Similar structure to $gg \rightarrow hh$

$$\kappa_i \equiv \frac{g_i}{g_i^{\text{SM},(0)}}$$

- With additional gauge sector contributions



- Compute the helicity amplitudes for the process $\gamma\gamma \rightarrow hh$
- Some diagrams only contribute to $J_z = 2$ and κ_λ only contained in the $J_z = 0$ contribution

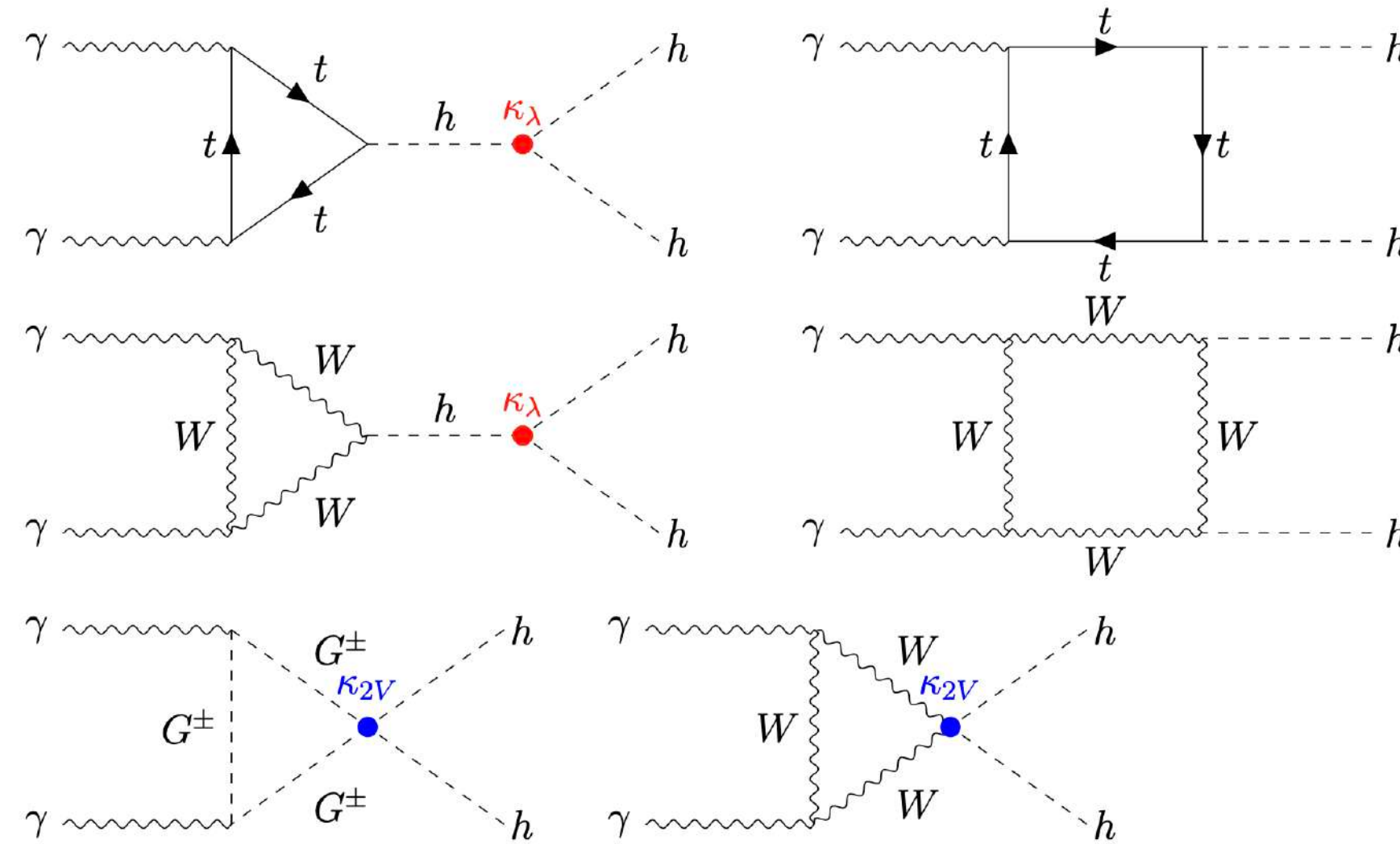
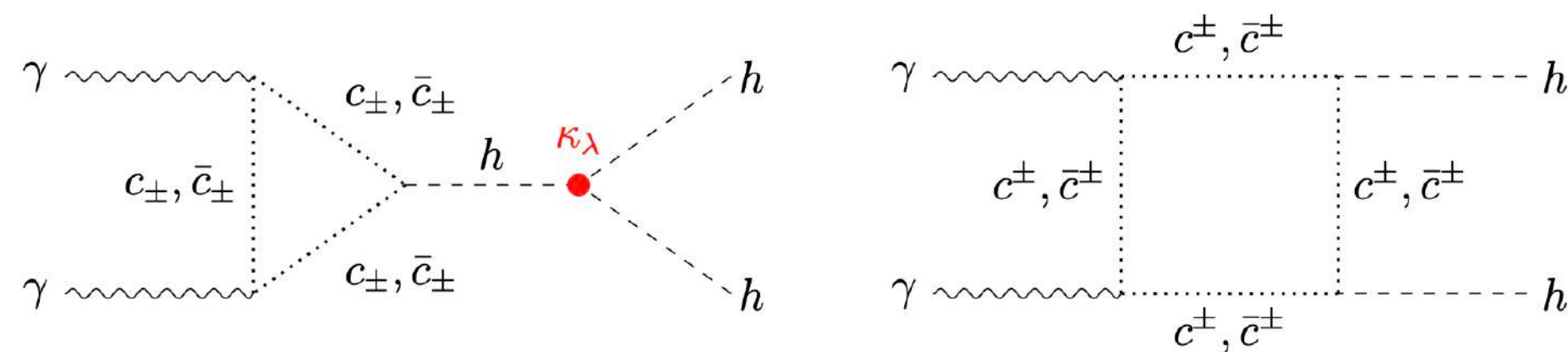
Higgs pair production - Frameworks

κ -framework

- Rescaling of couplings by hand
- Gauge invariance not guaranteed
- No consistent handling of higher-order / loop corrections

Couplings in this analysis:

- κ_λ
- κ_{2V}

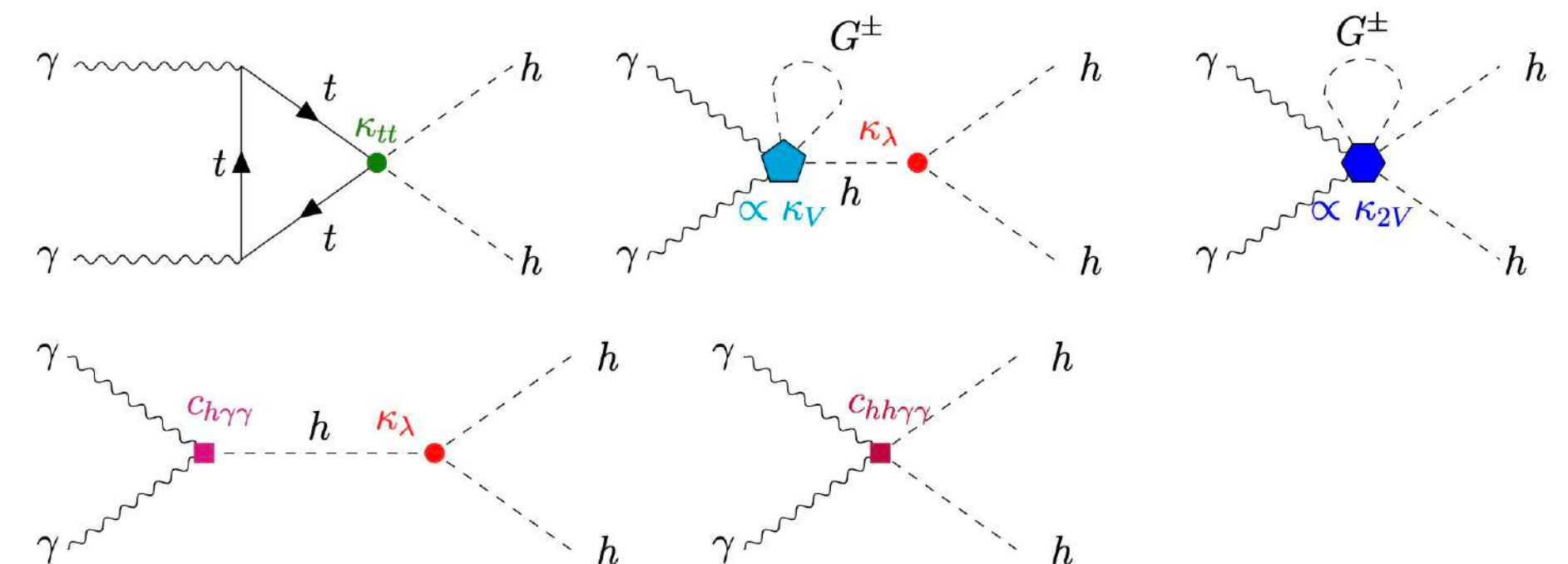


Higgs Effective Field Theory

- Independent study of single and double Higgs couplings
- Consistent EFT framework

Couplings in this analysis:

- κ_λ
- κ_{2V}
- κ_{tt}



Higgs Effective Field Theory

- The physical Higgs Boson is treated as a singlet, independent of the electroweak Goldstones

$$h, \quad U(x) = \exp \left(i \frac{\sigma^a \pi^a(x)}{v} \right)$$

- Electroweak symmetry breaking is realised nonlinear
- With the leading order Lagrangian

$$\begin{aligned} \mathcal{L}_{\text{HEFT}}^{\text{LO}} = & \mathcal{L}_{\text{gauge}} + \mathcal{L}_{\text{fermions}} + \mathcal{L}_{\text{Yukawa}} \\ & + \frac{v^2}{4} \mathcal{F}_C(h/v) \text{Tr} \left[\mathbf{D}_\mu \mathbf{U}^\dagger \mathbf{D}^\mu \mathbf{U} \right] + \frac{1}{2} \partial_\mu h \partial^\mu h - V(h) \end{aligned}$$

- Higgs interactions with Goldstone bosons and gauge bosons are described by the flare function $\mathcal{F}_C(h/v)$

$$\mathcal{F}_C(h) = 1 + 2a_C \frac{h}{v} + b_C \frac{h^2}{v^2} + c_C \frac{h^3}{v^3} + \dots$$

Higgs Effective Field Theory

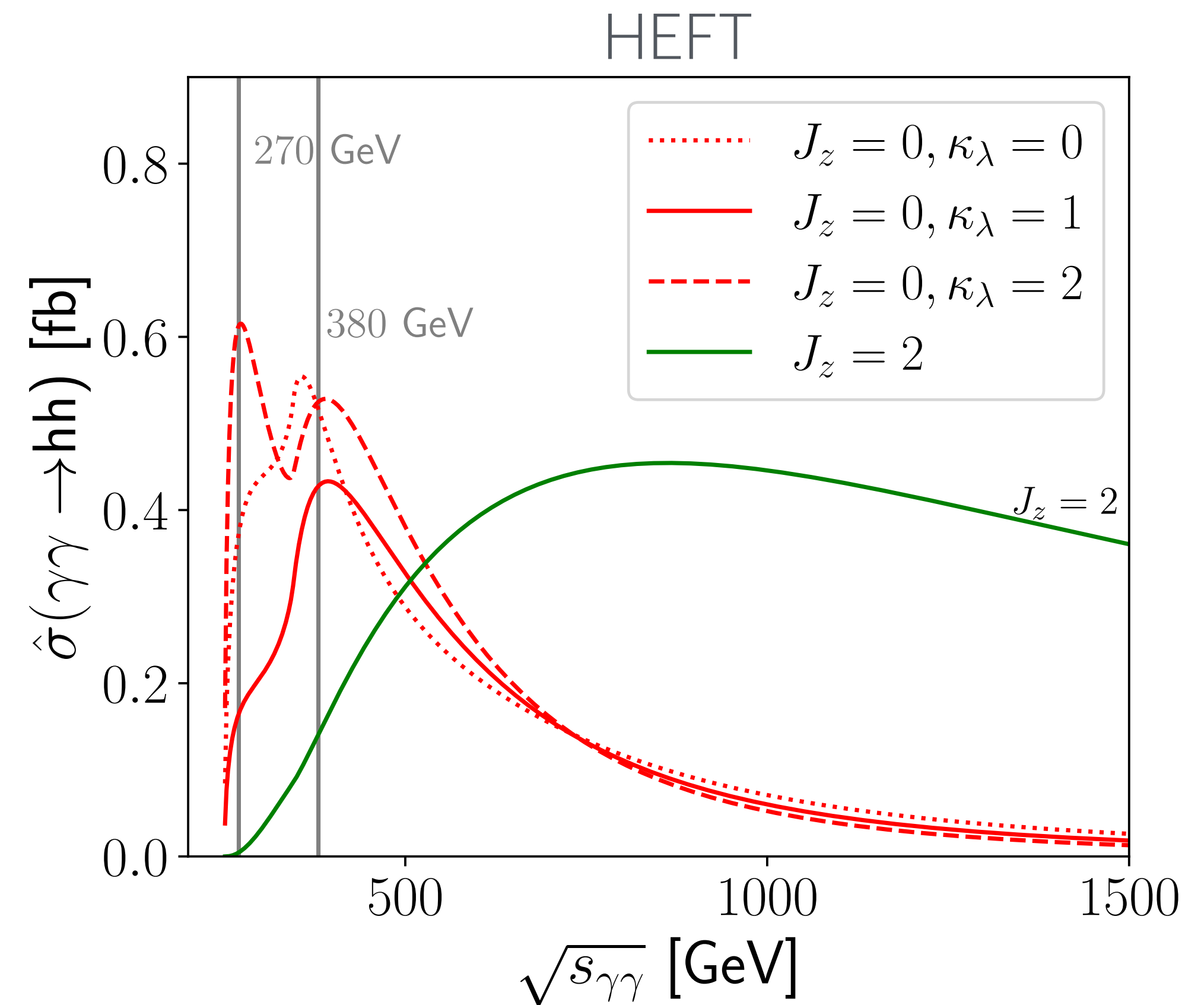
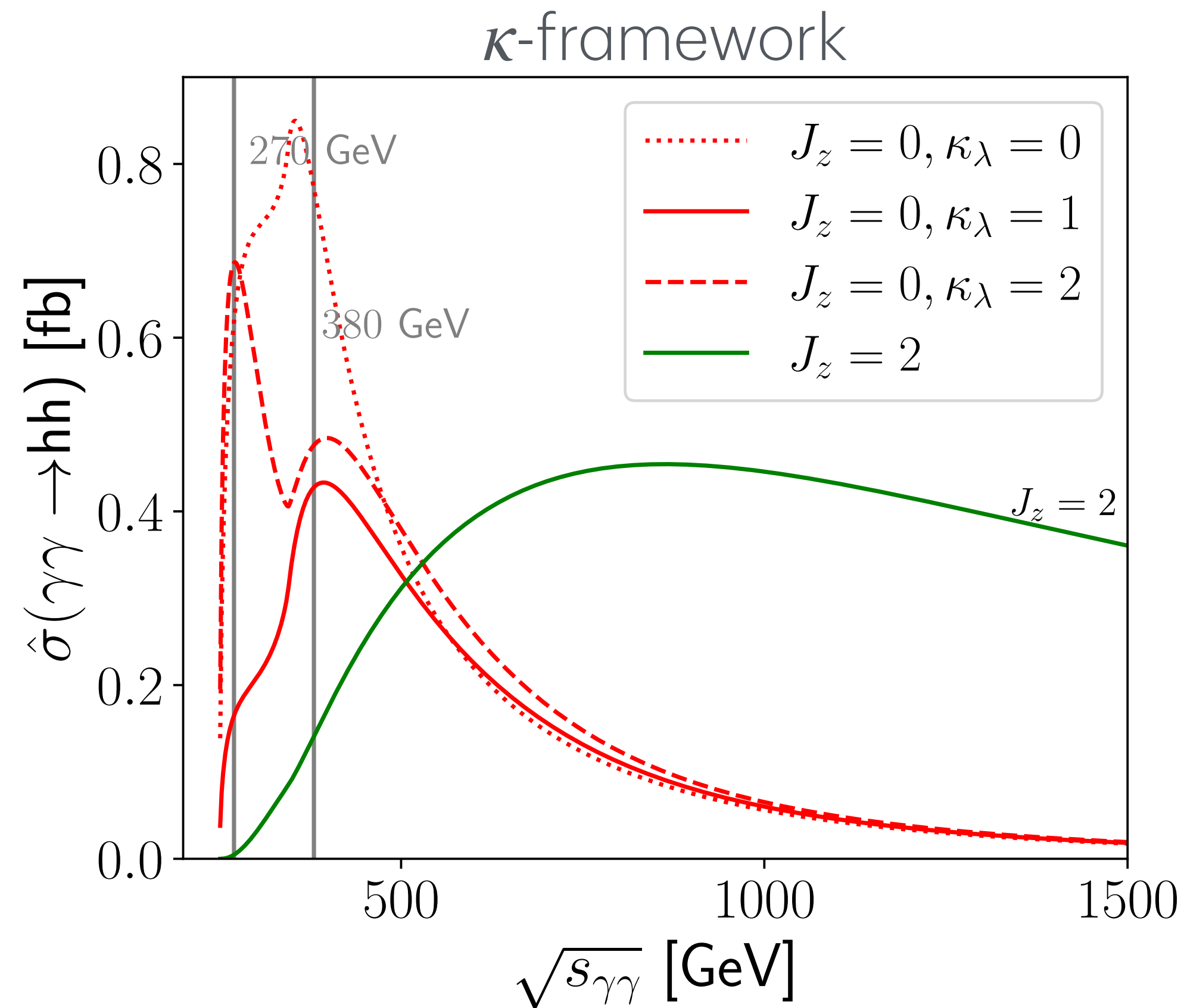
- The effective Lagrangian

$$\begin{aligned} \mathcal{L}_{\text{eff}} \supset & -\frac{1}{2} \kappa_\lambda \frac{m_h^2}{v} h^3 - \frac{y_t v}{\sqrt{2}} \bar{t} \left(\kappa_t \frac{h}{v} + \kappa_{tt} \frac{h^2}{v^2} \right) t \\ & + \frac{v^2}{4} \left(2\kappa_V \frac{h}{v} + \kappa_{2V} \frac{h^2}{v^2} \right) \text{Tr} \left[\mathbf{D}_\mu \mathbf{U}^\dagger \mathbf{D}^\mu \mathbf{U} \right] \\ & - \frac{\alpha}{16\pi} F_{\mu\nu} F^{\mu\nu} \left(c_{h\gamma\gamma} \frac{h}{v} + c_{hh\gamma\gamma} \frac{h^2}{v^2} \right) \end{aligned}$$

- In the SM: $\kappa_\lambda = 1, \kappa_V = 1, \kappa_{2V} = 1, \kappa_t = 1, \kappa_{tt} = 0, c_{h\gamma\gamma} = 0, c_{hh\gamma\gamma} = 0$
- We keep the other couplings at their SM value

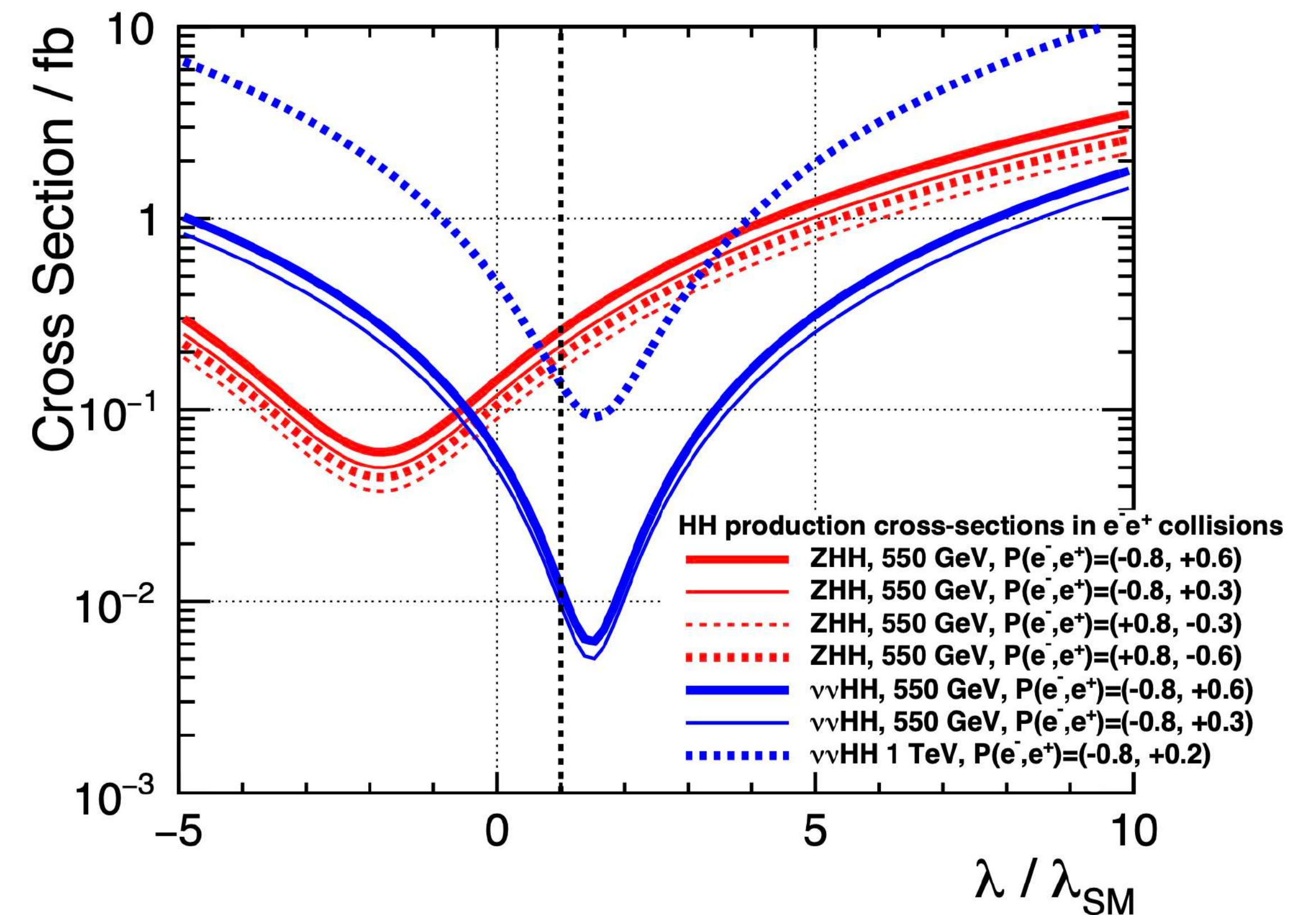
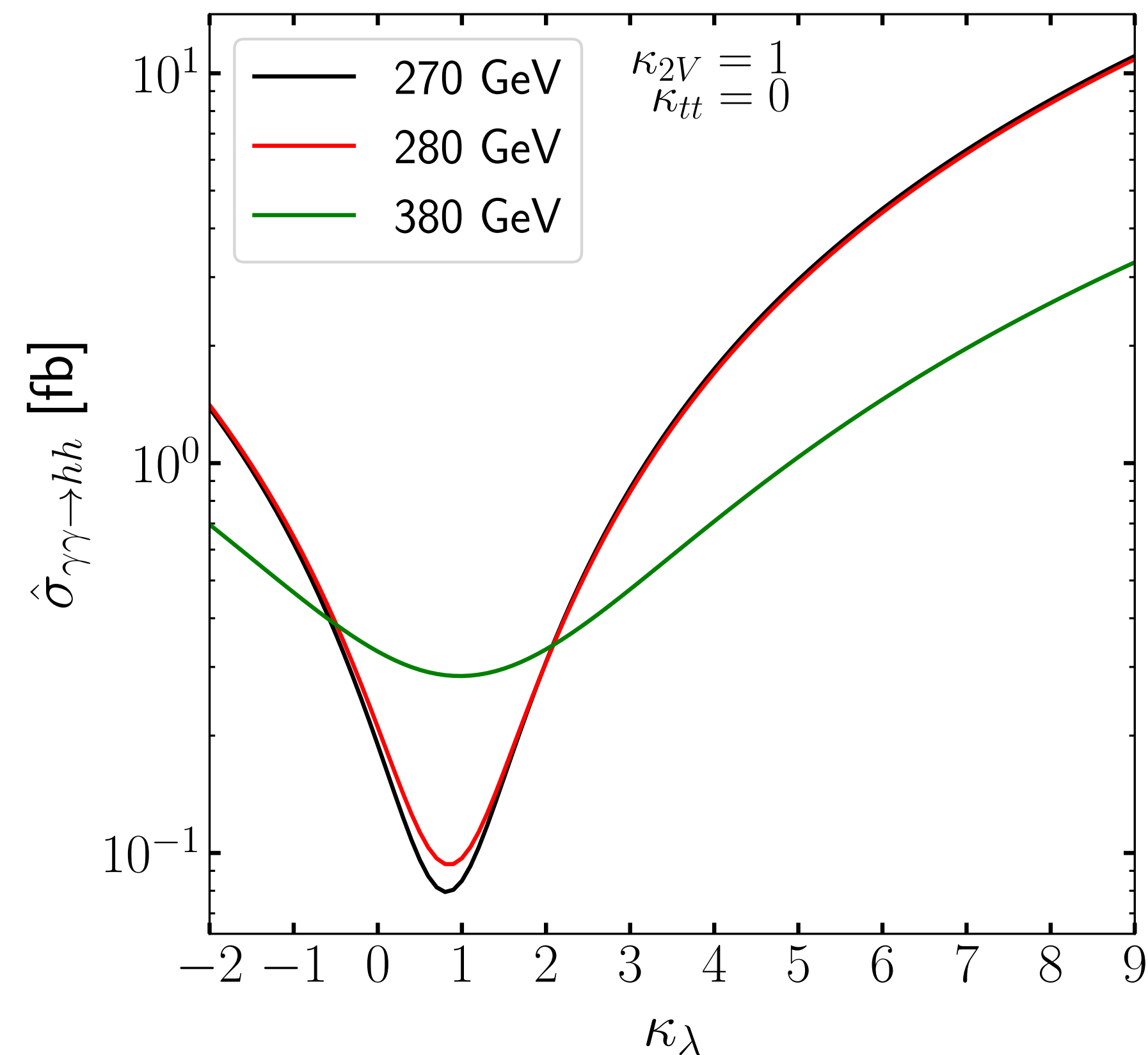
Higgs pair production - Partonic level

- Both frameworks agree for the SM cross-section
- Disagree for $\kappa_\lambda \neq 1$, need a consistent framework $\rightarrow$ HEFT



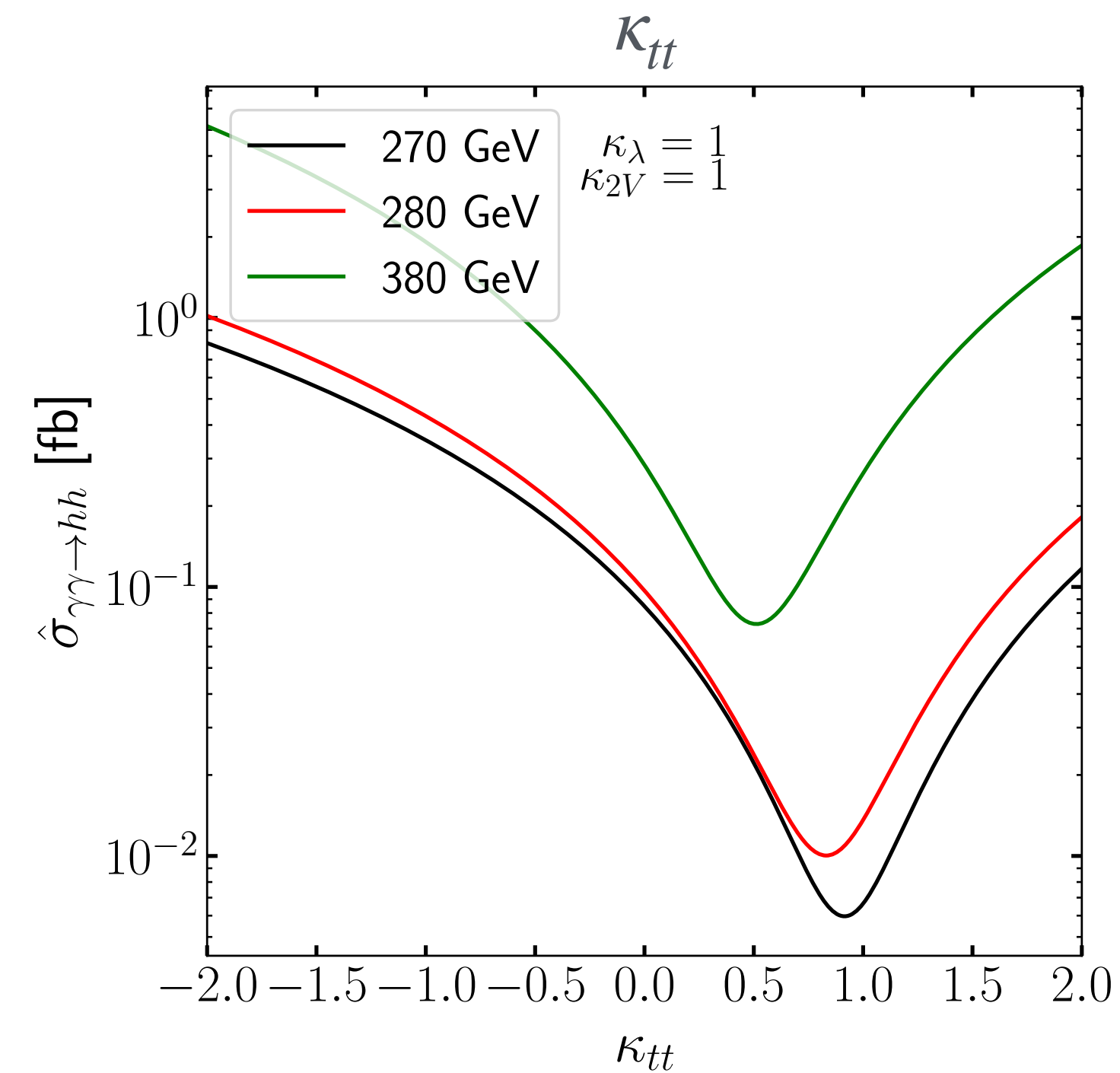
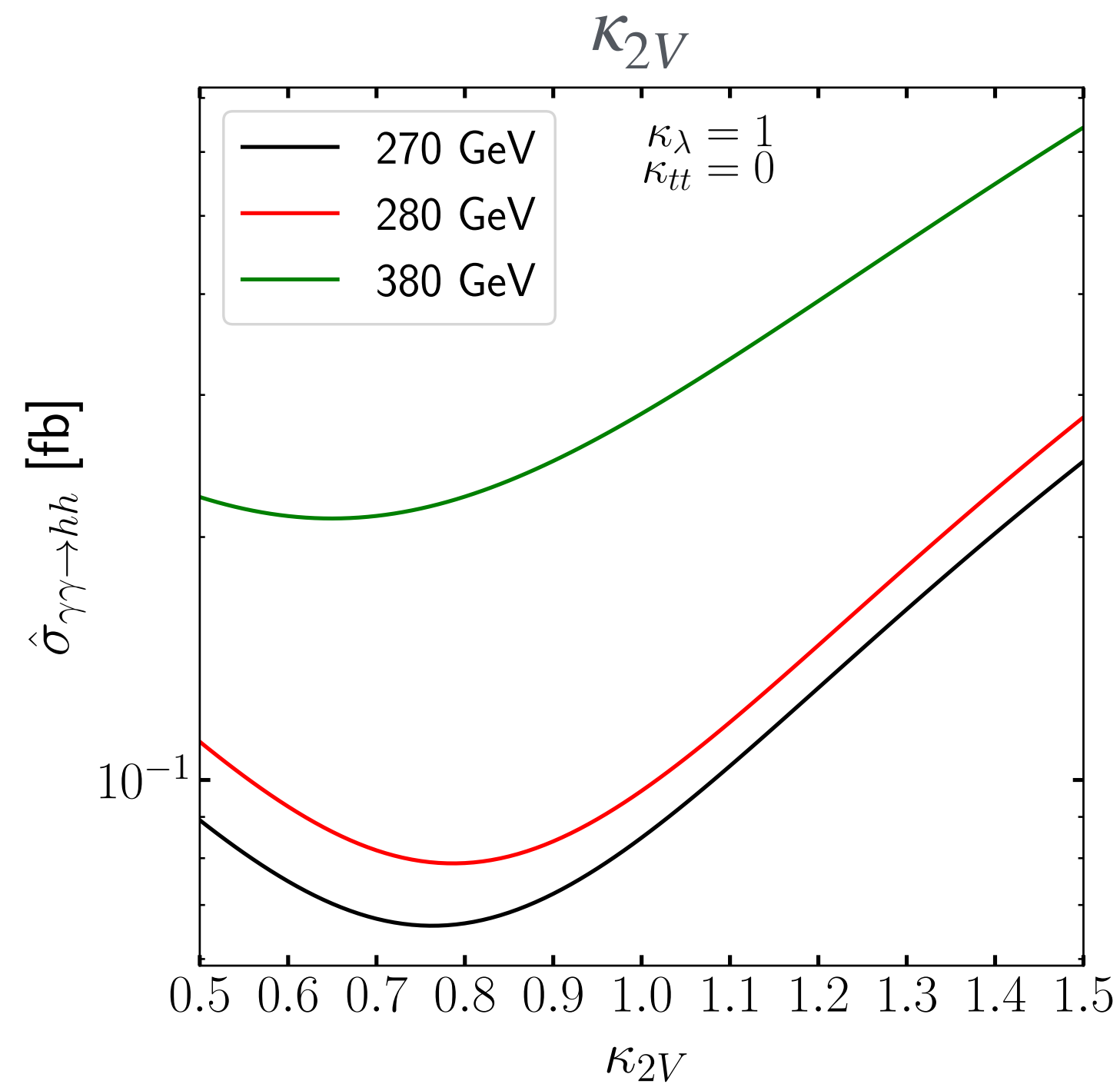
Higgs pair production - Partonic level

- The minimum is at different values of κ_λ , offers complementarity between different collider searches
 - Lower energies are better for sensitivity studies due to the strong dependency on κ_λ [H. Abramowicz, **MB** et al. '25]
- e^+e^- collider



Higgs pair production - Partonic level

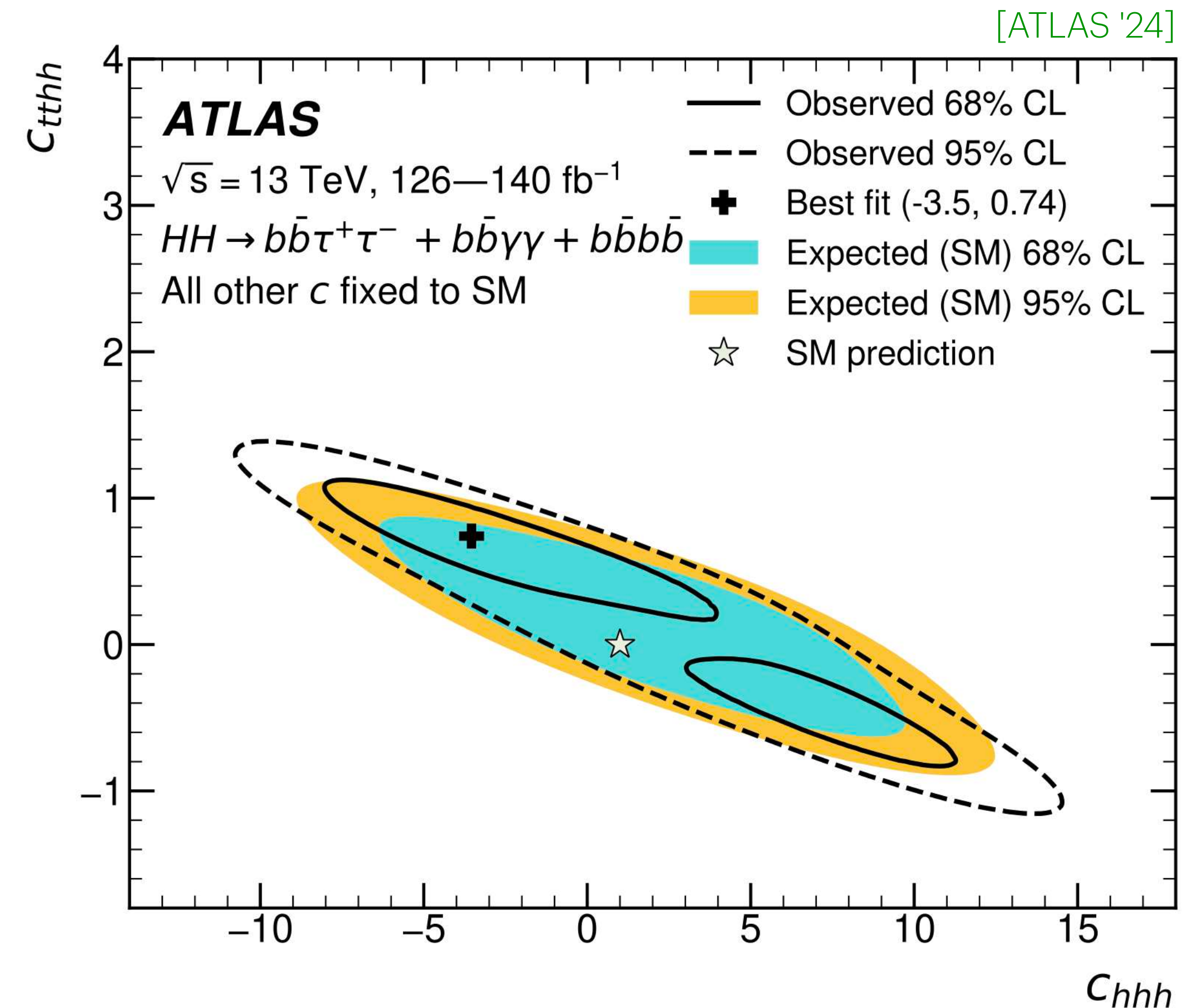
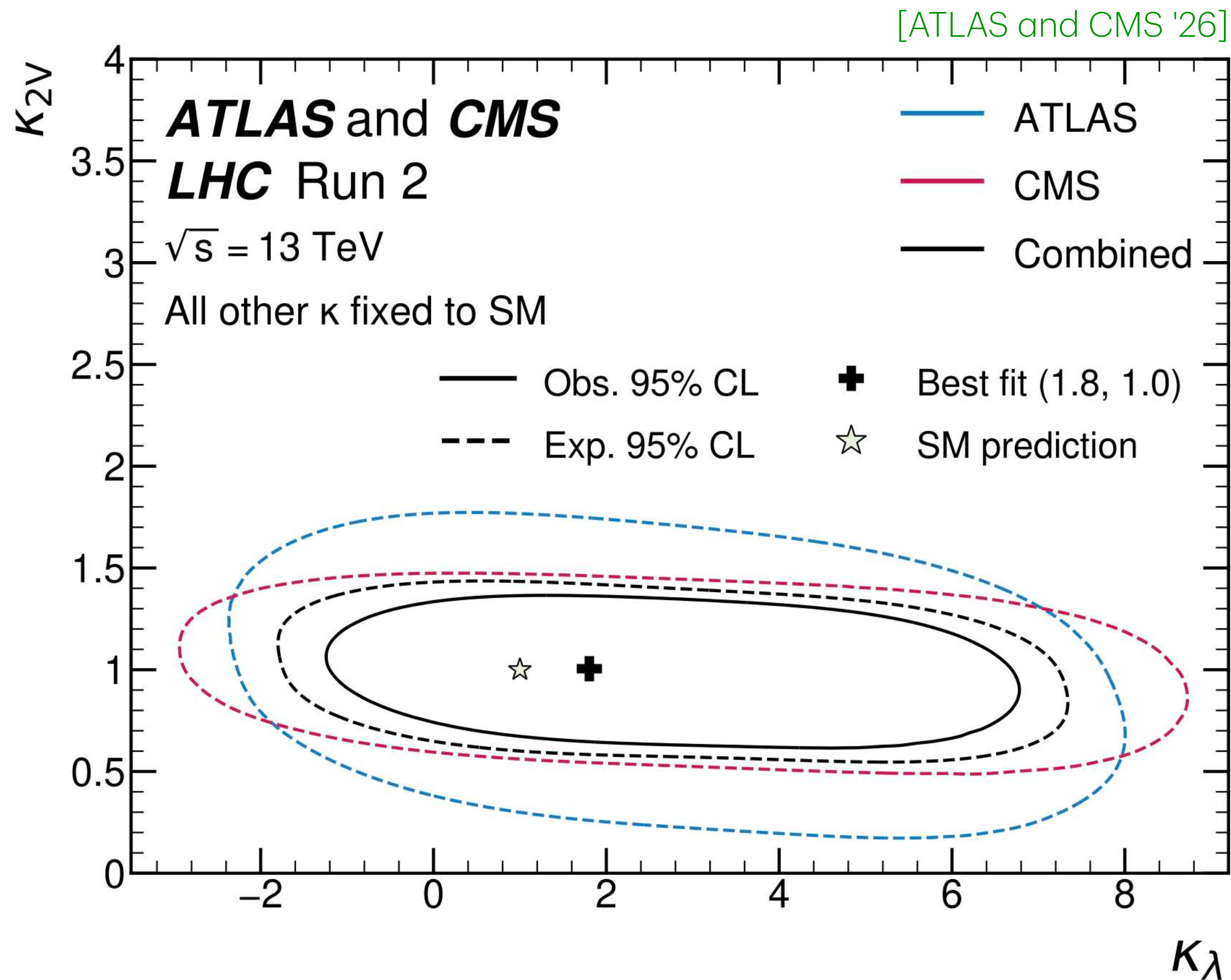
- The dependency on κ_{2V} and κ_{tt} does not show different behaviour at the different energies
- No s-channel suppression compared to the κ_λ dependency



Higgs pair production - Limits

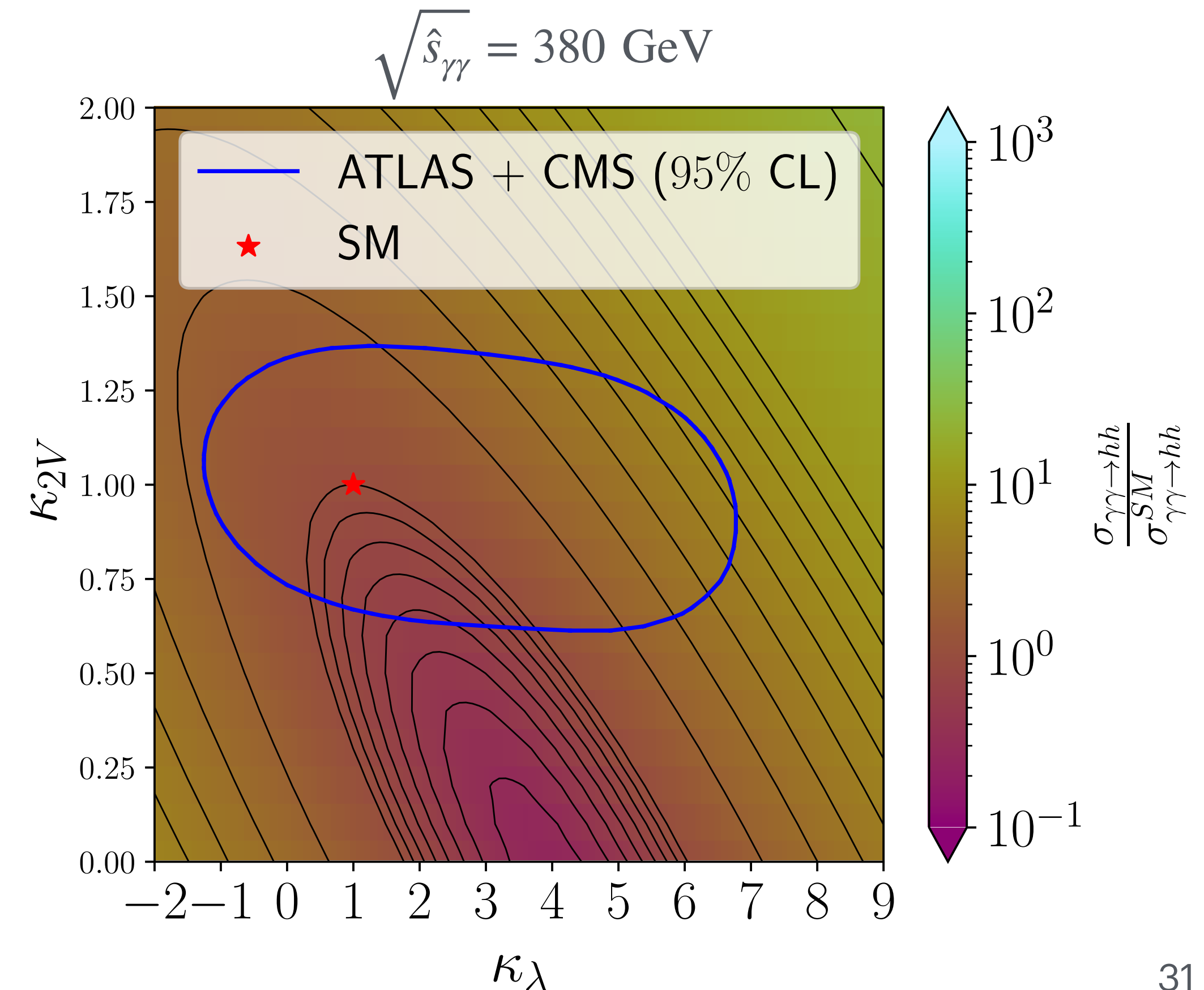
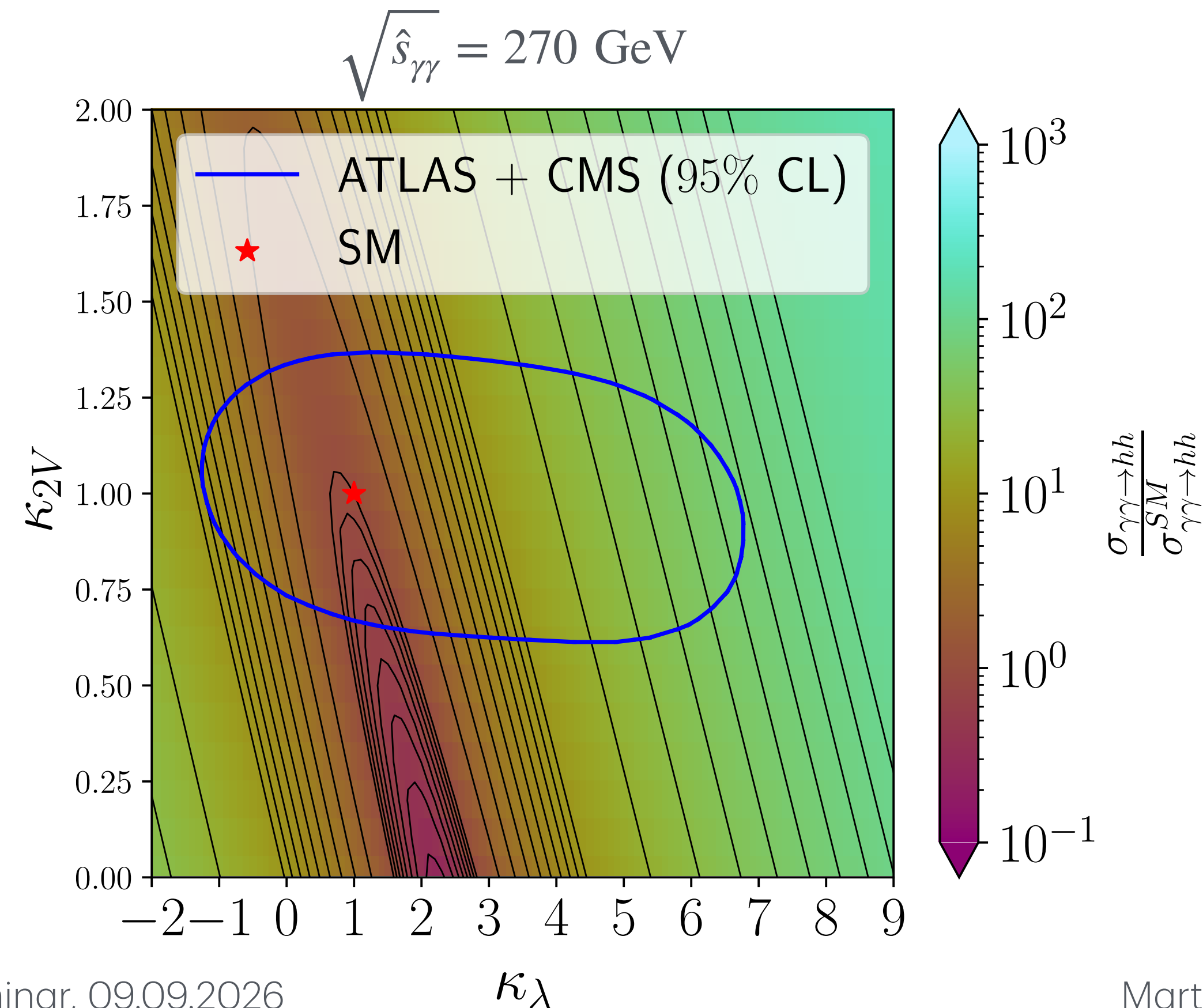
$[-0.71 < \kappa_\lambda < 6.1]$ [ATLAS and CMS '26]
while all other couplings
fixed to SM

- By including additional couplings the sensitivity on κ_λ decreases



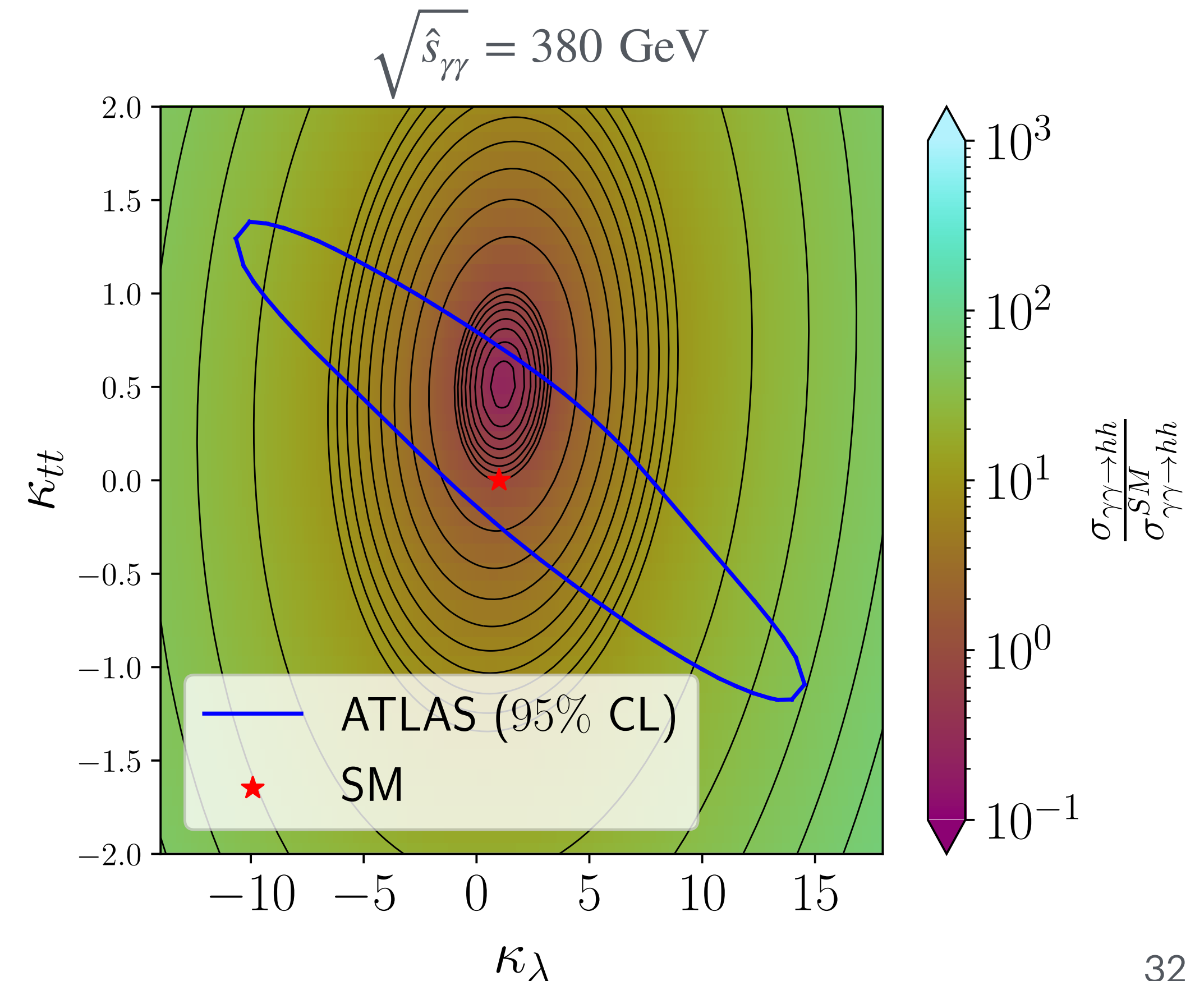
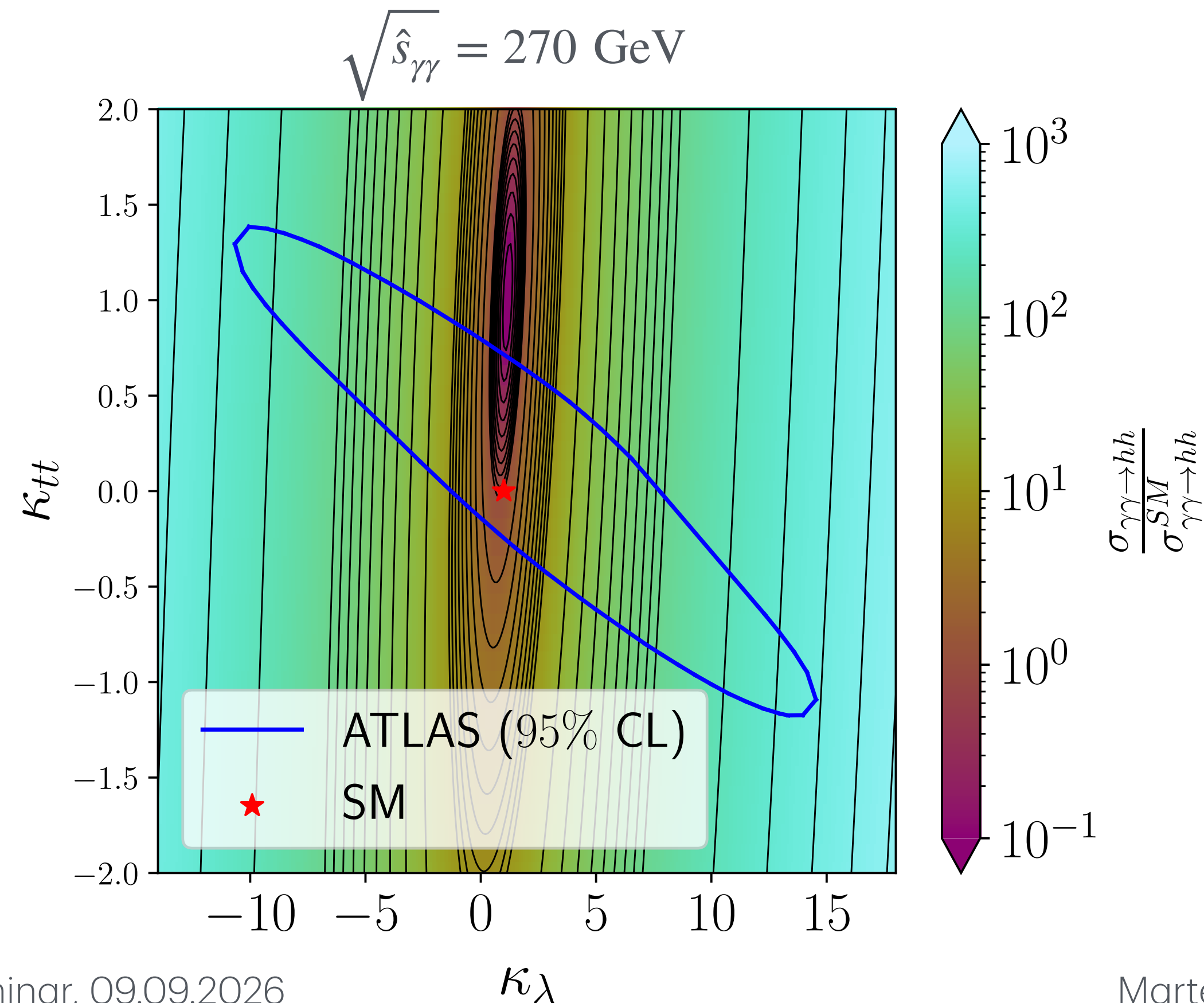
Higgs pair production - Partonic level

- The contour in the $\kappa_\lambda/\kappa_{2V}$ plane at 270 GeV is close to vertical behaviour, while at 380 GeV at larger angle
- Cross-section dependency shown within the combined 95% CL from ATLAS and CMS [ATLAS and CMS '26]

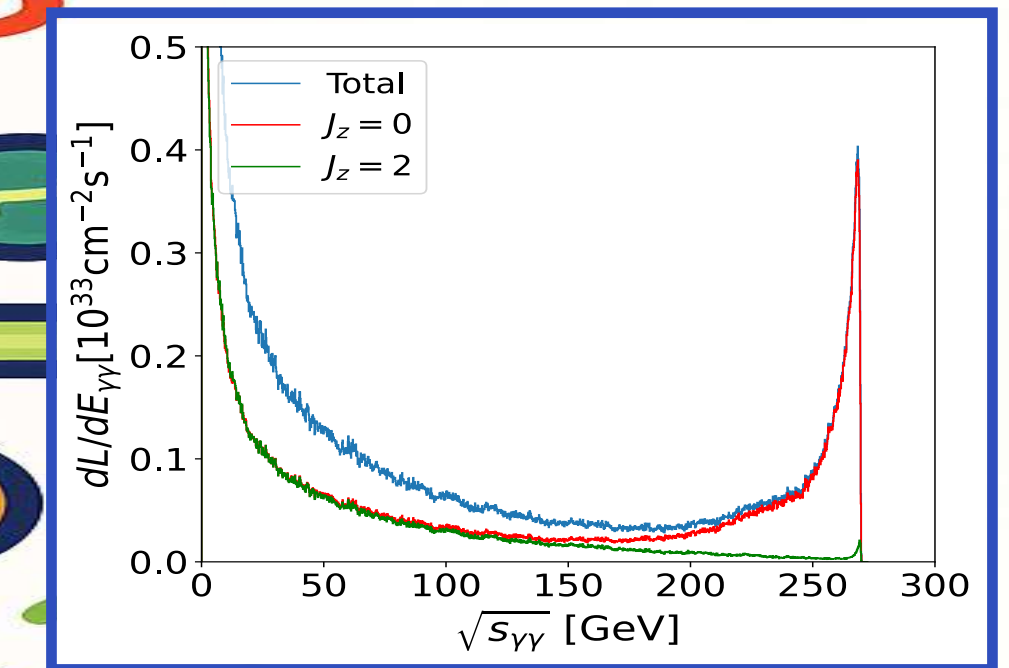
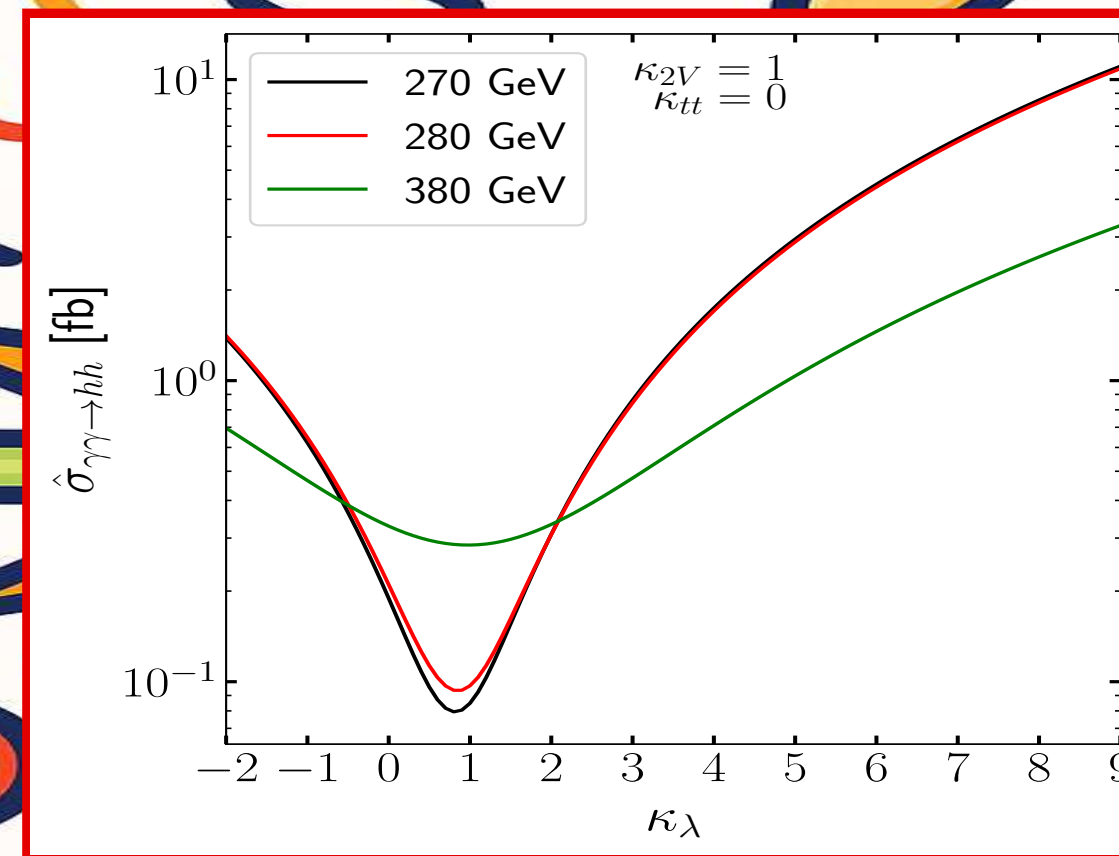
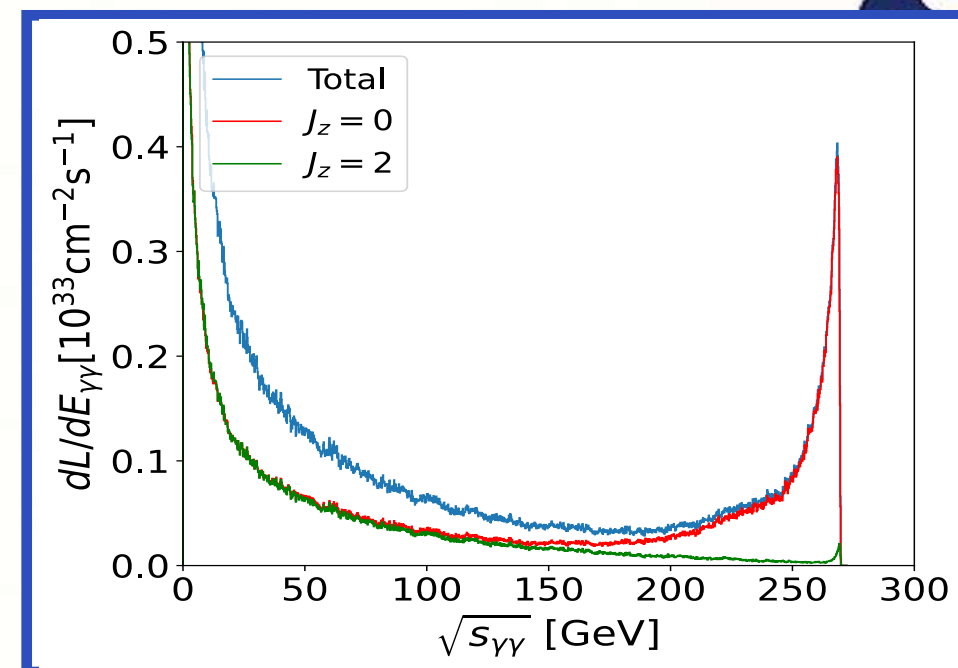


Higgs pair production - Partonic level

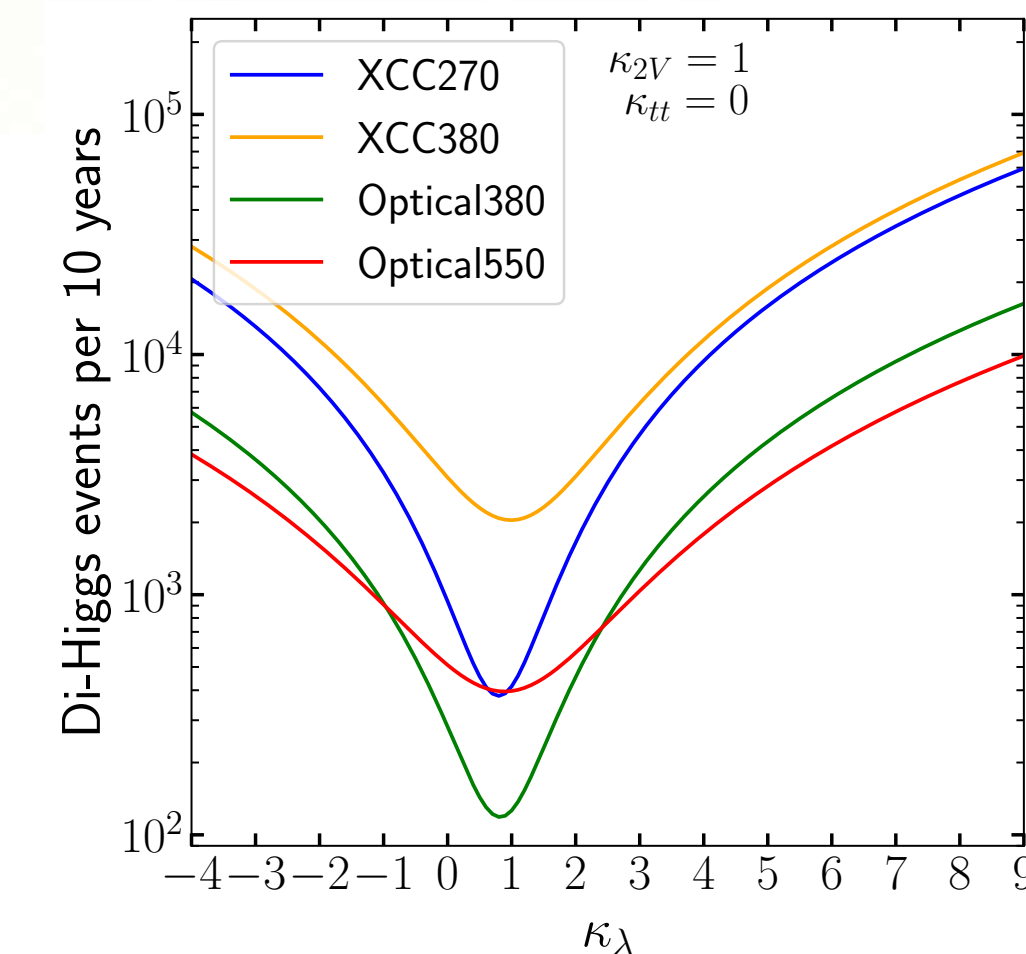
- The contours in the $\kappa_\lambda/\kappa_{tt}$ plane at both energies arrange in nearly vertical ellipses, with very tight spacing at 270 GeV
- Cross-section dependency shown within the 95% CL from ATLAS [ATLAS '24]



Photon collider - Higgs pair production



- Take both results and smash them together

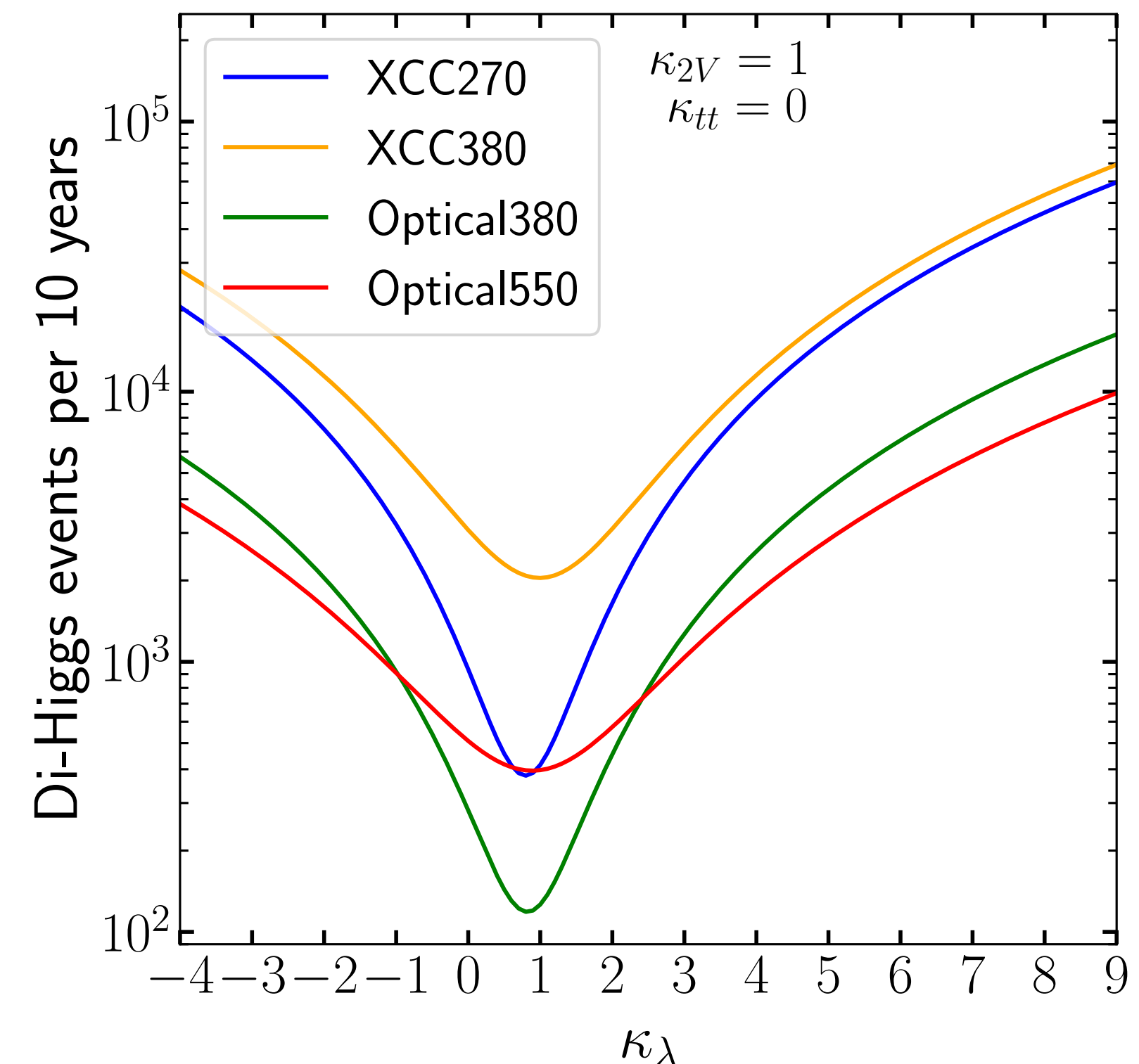


Photon collider - Higgs pair production

- Combine the **partonic cross section** and the **luminosity spectra**

$$\sigma = \int_{4m_h^2/\hat{s}}^{y_m^2} d\tau \frac{1}{2} \left[\frac{1}{L_{\gamma\gamma}} \frac{dL_{\gamma\gamma}^{++}}{d\tau} \hat{\sigma}_{++}(\hat{s}) + \frac{1}{L_{\gamma\gamma}} \frac{dL_{\gamma\gamma}^{+-}}{d\tau} \hat{\sigma}_{+-}(\hat{s}) \right]$$

- 10 years of runtime ($1.6 * 10^7$ seconds per year)
- SM Higgs pair production
 - Optical380: 126 events
 - Optical550: 397 events
 - XCC270: 415 events
 - XCC380: 2043 events



Photon collider - Higgs pair production

- Combine the **partonic cross section** and the **luminosity spectra**

$$\sigma = \int_{4m_h^2/\hat{s}}^{y_m^2} d\tau \frac{1}{2} \left[\frac{1}{L_{\gamma\gamma}} \frac{dL_{\gamma\gamma}^{++}}{d\tau} \hat{\sigma}_{++}(\hat{s}) + \frac{1}{L_{\gamma\gamma}} \frac{dL_{\gamma\gamma}^{+-}}{d\tau} \hat{\sigma}_{+-}(\hat{s}) \right]$$

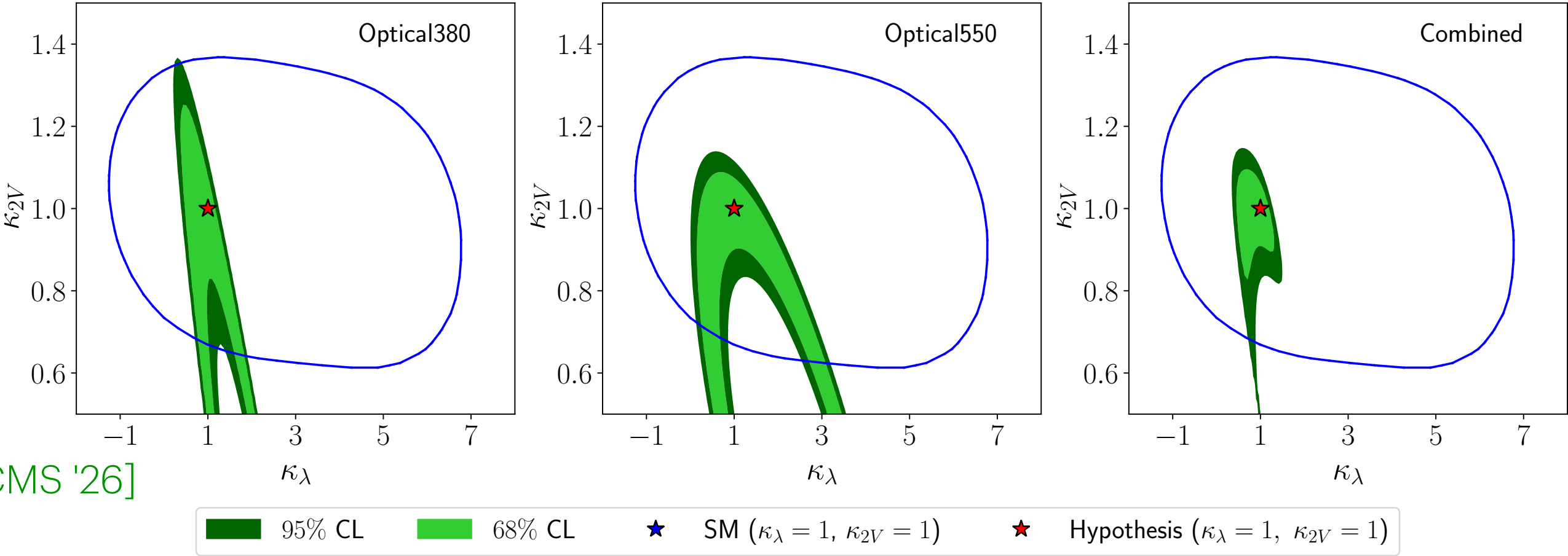
- 10 years of runtime ($1.6 * 10^7$ seconds per year)
- Run 5 years at each of the energies and combine the sensitivity

$$\chi^2 = \sum \frac{(O_i - E_i)^2}{E_i + (\epsilon E_i)^2}$$

- For each choice of true κ_λ , κ_{2V} and κ_{tt} confidence level regions have to be calculated

Optical collider - $\kappa_\lambda/\kappa_{2V}$ analysis

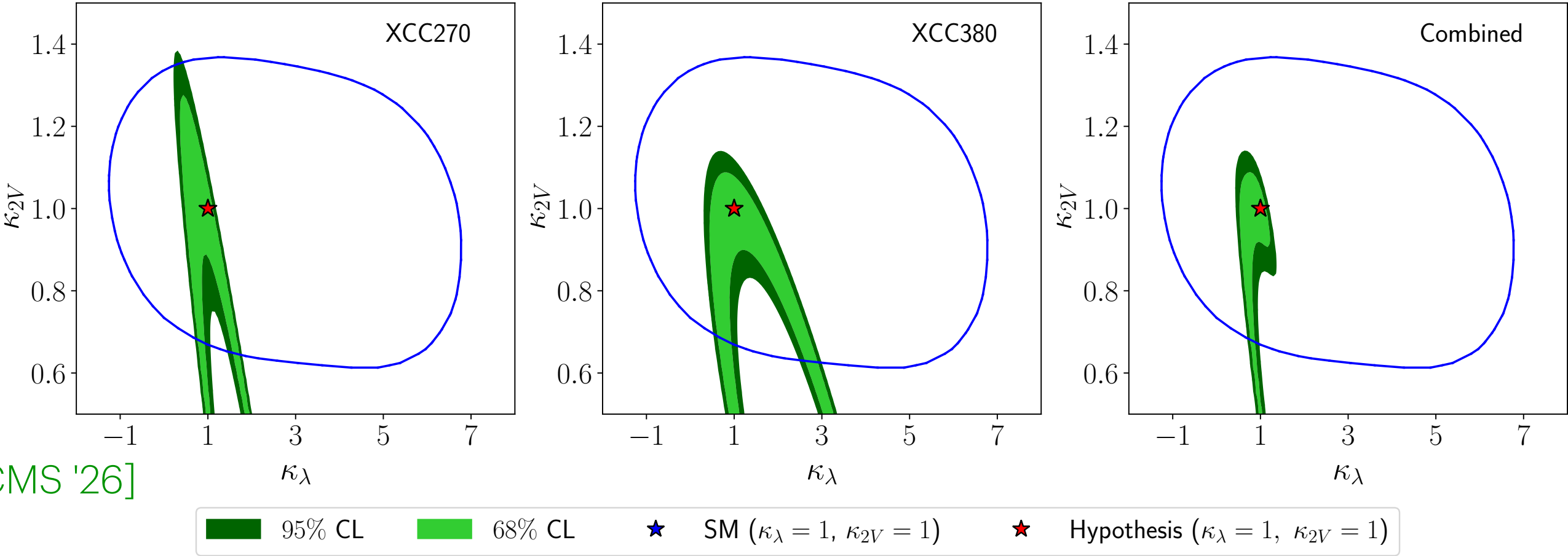
- Possible to achieve more precise limits
- In the full plane between κ_λ and κ_{2V}
- Current limits by ATLAS and CMS:
 - $\kappa_\lambda \in [-0.71, 6.1]$ and $\kappa_{2V} \in [0.73, 1.3]$ [ATLAS, CMS '26]



$\kappa_\lambda^{\text{true}} = 1, \kappa_{2V}^{\text{true}} = 1$ 95 % CL	Both parameters free	Both parameters free	$\kappa_{2V} = 1$	$\kappa_\lambda = 1$
Optical380	$\kappa_\lambda \in [0.25, 1.81]$	$\kappa_{2V} \in [0.66, 1.35]$	$\kappa_\lambda \in [0.44, 1.23]$	$\kappa_{2V} \in [0.68, 1.13]$
Optical550	$\kappa_\lambda \in [0.05, 3.06]$	$\kappa_{2V} \in [0.66, 1.14]$	$\kappa_\lambda \in [0.12, 1.62]$	$\kappa_{2V} \in [0.84, 1.10]$
Optical combined	$\kappa_\lambda \in [0.38, 1.49]$	$\kappa_{2V} \in [0.68, 1.15]$	$\kappa_\lambda \in [0.44, 1.23]$	$\kappa_{2V} \in [0.87, 1.08]$

XCC - $\kappa_\lambda/\kappa_{2V}$ analysis

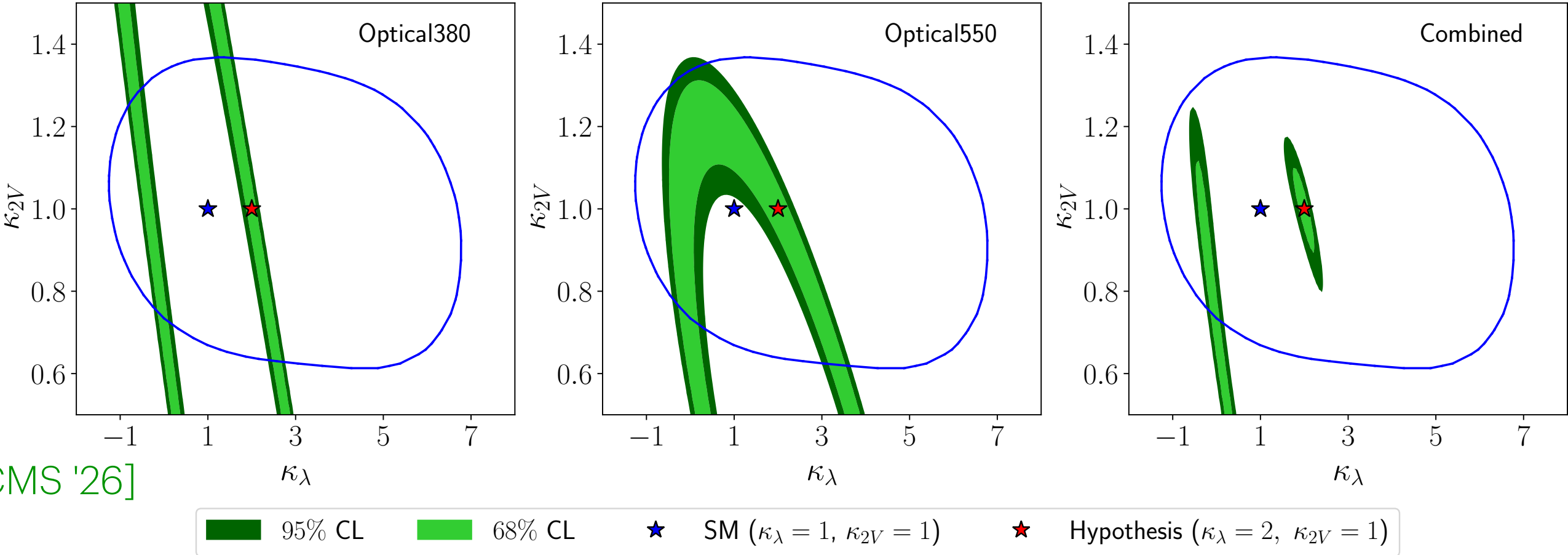
- Possible to achieve more precise limits
- In the full plane between κ_λ and κ_{2V}
- Current limits by ATLAS and CMS:
 - $\kappa_\lambda \in [-0.71, 6.1]$ and $\kappa_{2V} \in [0.73, 1.3]$ [ATLAS, CMS '26]



$\kappa_\lambda^{\text{true}} = 1, \kappa_{2V}^{\text{true}} = 1$ 95 % CL	Both parameters free	Both parameters free	$\kappa_{2V} = 1$	$\kappa_\lambda = 1$
XCC270	$\kappa_\lambda \in [0.25, 1.75]$	$\kappa_{2V} \in [0.66, 1.35]$	$\kappa_\lambda \in [0.44, 1.16]$	$\kappa_{2V} \in [0.68, 1.10]$
XCC380	$\kappa_\lambda \in [0.31, 2.86]$	$\kappa_{2V} \in [0.66, 1.14]$	$\kappa_\lambda \in [0.38, 1.62]$	$\kappa_{2V} \in [0.68, 1.10]$
XCC combined	$\kappa_\lambda \in [0.44, 1.36]$	$\kappa_{2V} \in [0.68, 1.14]$	$\kappa_\lambda \in [0.51, 1.16]$	$\kappa_{2V} \in [0.89, 1.07]$

Optical collider - $\kappa_\lambda/\kappa_{2V}$ analysis

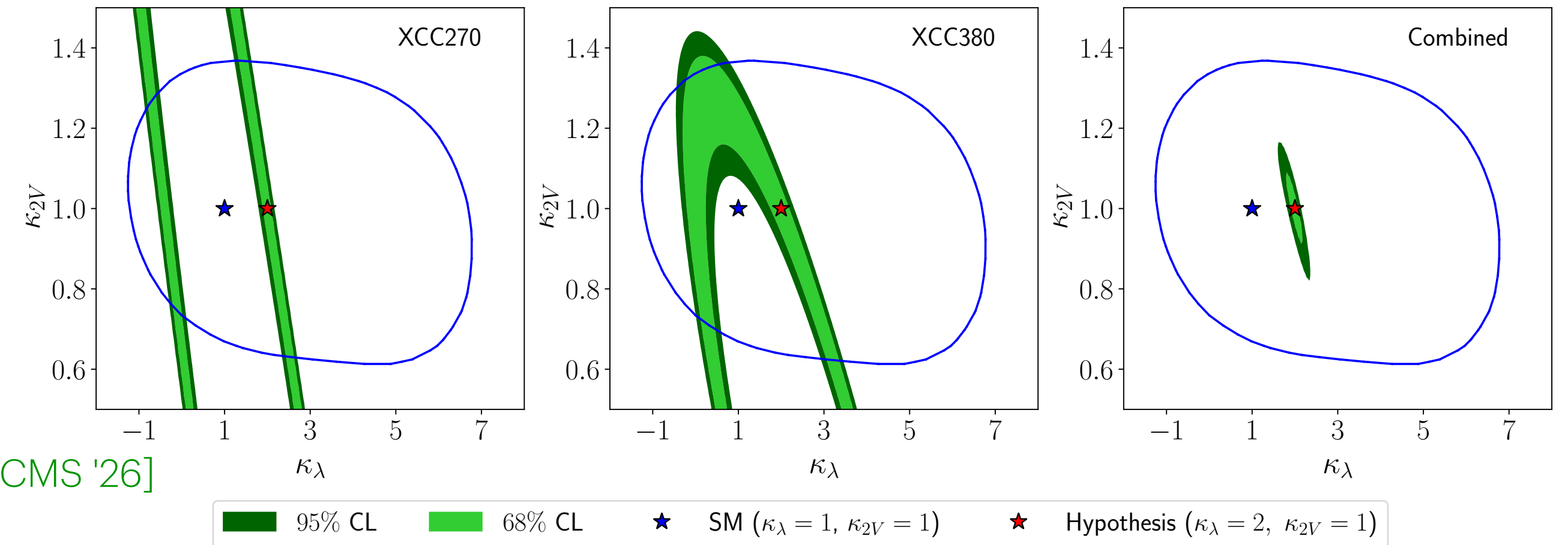
- Possible to achieve more precise limits
- In the full plane between κ_λ and κ_{2V}
- Current limits by ATLAS and CMS:
 - $\kappa_\lambda \in [-0.71, 6.1]$ and $\kappa_{2V} \in [0.73, 1.3]$ [ATLAS, CMS '26]



$\kappa_\lambda^{\text{true}} = 2, \kappa_{2V}^{\text{true}} = 1$ 95 % CL	Both parameters free	Both parameters free	$\kappa_{2V} = 1$	$\kappa_\lambda = 2$
Optical380	$\kappa_\lambda \in [-0.80, 0.18] \cup [1.16, 2.66]$	$\kappa_{2V} \in [0.66, 1.36]$	$\kappa_\lambda \in [-0.47, -0.21] \cup [1.81, 2.14]$	$\kappa_{2V} \in [0.87, 1.08]$
Optical550	$\kappa_\lambda \in [-0.60, 3.51]$	$\kappa_{2V} \in [0.66, 1.35]$	$\kappa_\lambda \in [-0.54, 0.25] \cup [1.42, 2.34]$	$\kappa_{2V} \in [0.90, 1.07]$
Optical combined	$\kappa_\lambda \in [-0.60, 0.12] \cup [1.55, 2.40]$	$\kappa_{2V} \in [0.73, 1.25]$	$\kappa_\lambda \in [-0.47, -0.21] \cup [1.81, 2.14]$	$\kappa_{2V} \in [0.93, 1.05]$

XCC - $\kappa_\lambda/\kappa_{2V}$ analysis

- Possible to achieve more precise limits
- In the full plane between κ_λ and κ_{2V}
- Current limits by ATLAS and CMS:
 - $\kappa_\lambda \in [-0.71, 6.1]$ and $\kappa_{2V} \in [0.73, 1.3]$ [ATLAS, CMS '26]

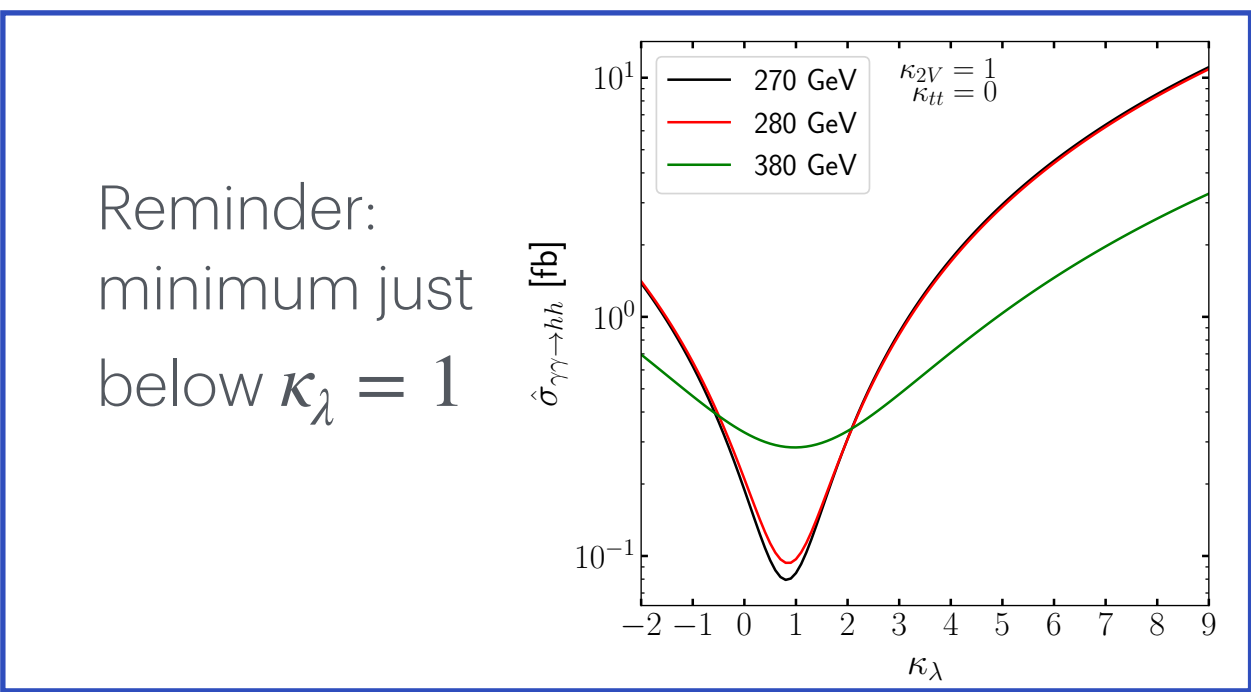


$\kappa_\lambda^{\text{true}} = 2, \kappa_{2V}^{\text{true}} = 1$ 95 % CL	Both parameters free	Both parameters free	$\kappa_{2V} = 1$	$\kappa_\lambda = 2$
XCC270	$\kappa_\lambda \in [-0.80, 0.05] \cup [1.29, 2.60]$	$\kappa_{2V} \in [0.66, 1.36]$	$\kappa_\lambda \in [-0.54, -0.28] \cup [1.88, 2.14]$	$\kappa_{2V} \in [0.89, 1.08]$
XCC380	$\kappa_\lambda \in [-0.41, 3.32]$	$\kappa_{2V} \in [0.66, 1.36]$	$\kappa_\lambda \in [-0.28, 0.31] \cup [1.68, 2.21]$	$\kappa_{2V} \in [0.92, 1.07]$
XCC combined	$\kappa_\lambda \in [1.62, 2.34]$	$\kappa_{2V} \in [0.82, 1.16]$	$\kappa_\lambda \in [1.88, 2.08]$	$\kappa_{2V} \in [0.95, 1.04]$

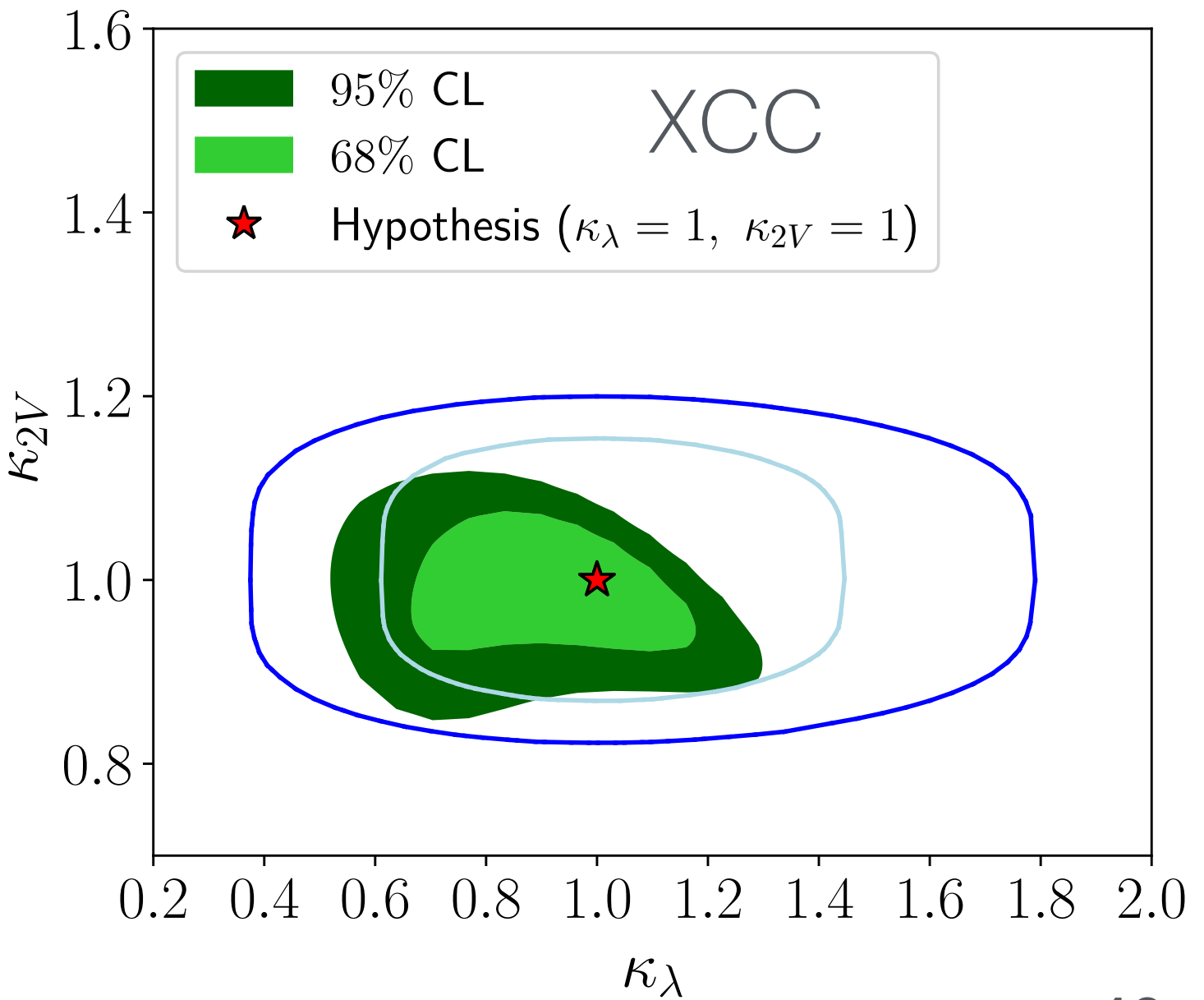
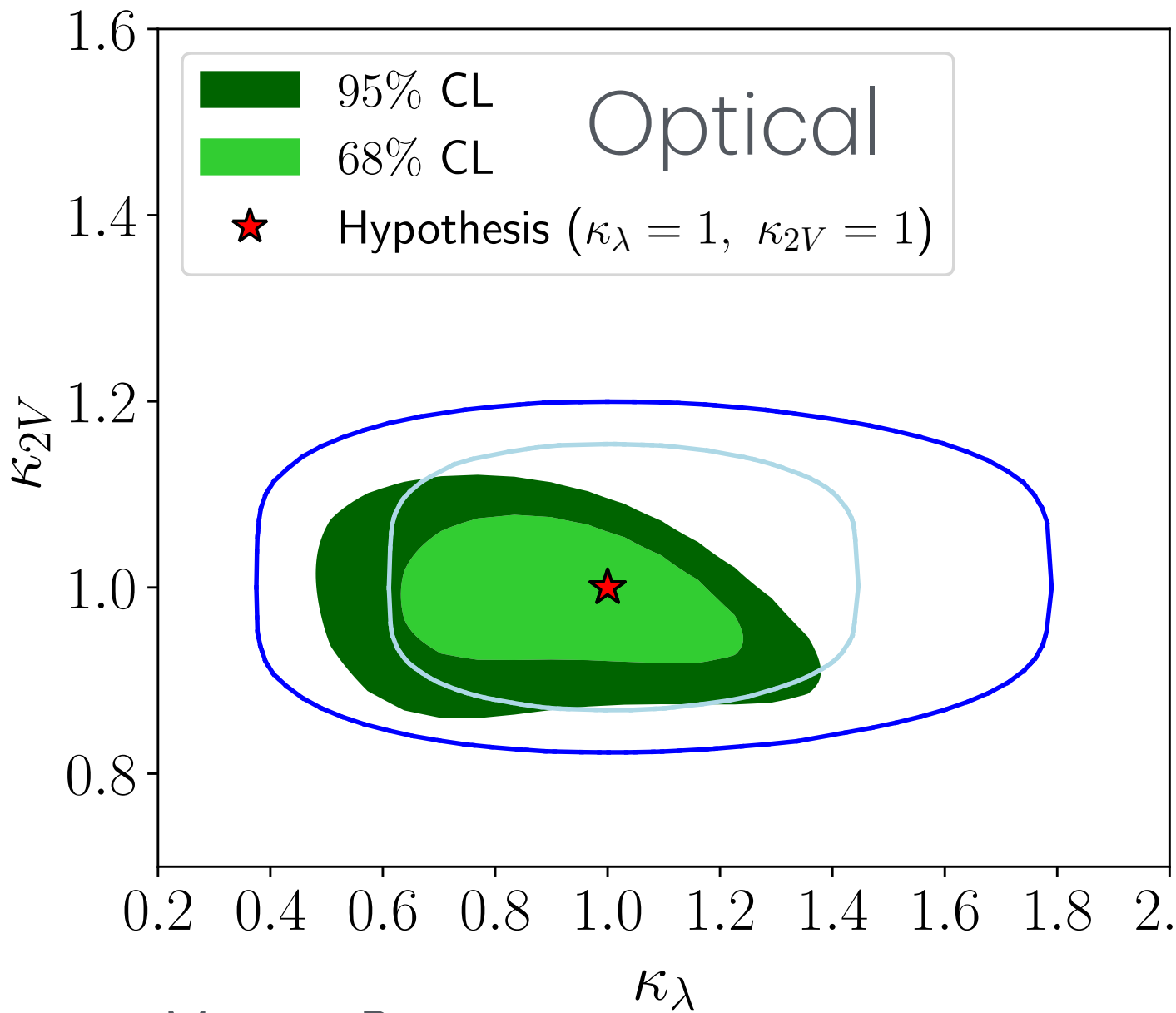
HL-LHC + Photon collider

- Combining the projected limits for HL-LHC with the photon collider

	κ_λ 68% CL	κ_λ 95% CL	κ_{2V} 68% CL	κ_{2V} 95% CL	$\kappa_{2V} = 1$ 68% CL	$\kappa_{2V} = 1$ 95% CL
Optical combined	$\kappa_\lambda \in [0.64, 1.23]$	$\kappa_\lambda \in [0.51, 1.36]$	$\kappa_{2V} \in [0.92, 1.07]$	$\kappa_{2V} \in [0.86, 1.12]$	$\kappa_\lambda \in [0.77, 1.10]$	$\kappa_\lambda \in [0.57, 1.23]$
XCC combined	$\kappa_\lambda \in [0.70, 1.16]$	$\kappa_\lambda \in [0.57, 1.29]$	$\kappa_{2V} \in [0.93, 1.07]$	$\kappa_{2V} \in [0.85, 1.12]$	$\kappa_\lambda \in [0.83, 1.03]$	$\kappa_\lambda \in [0.64, 1.16]$



- XCC at 68% CL manages -17%, +3% for $\kappa_{2V} = 1$
- Asymmetry due to the location of the cross-section minimum

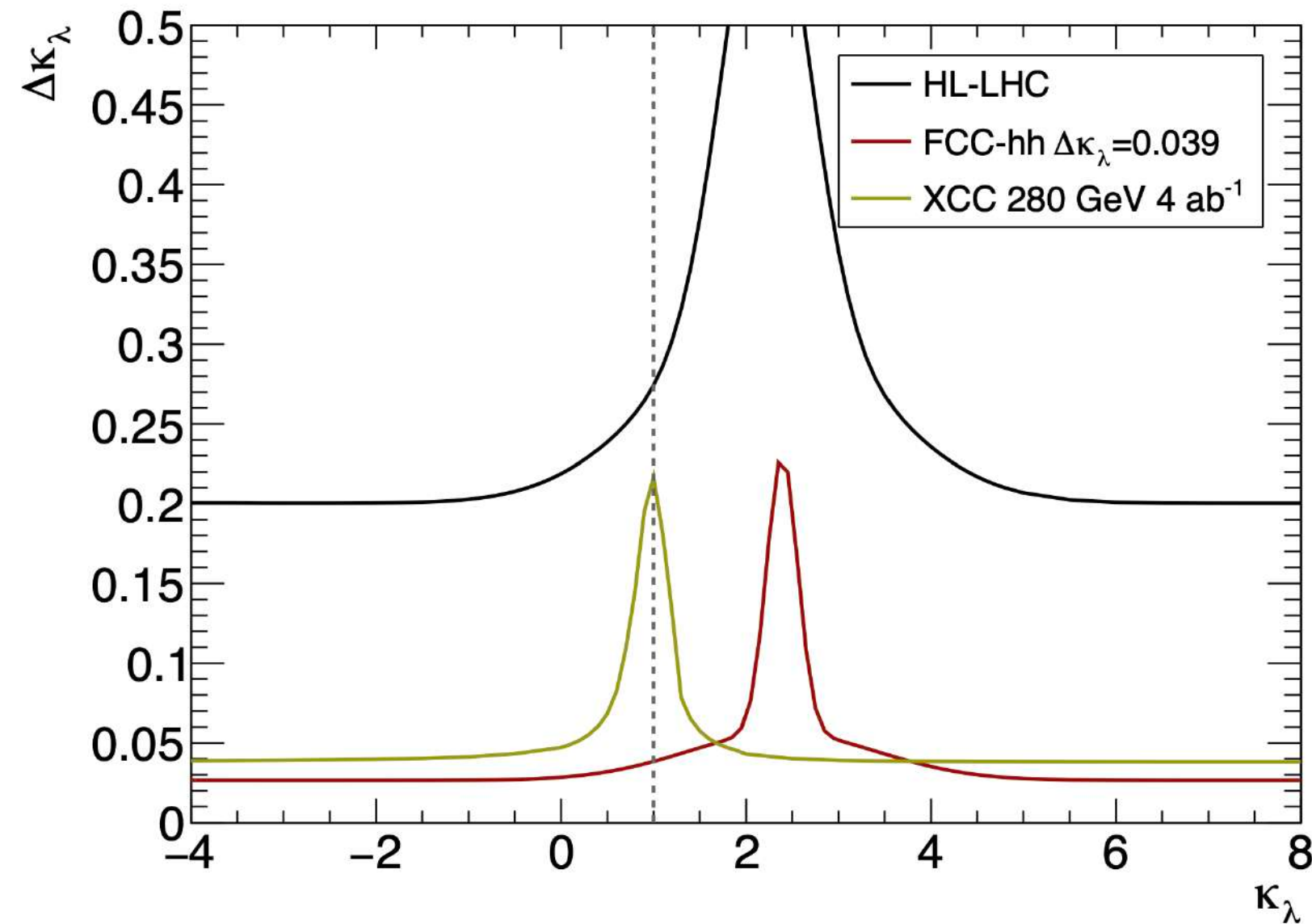
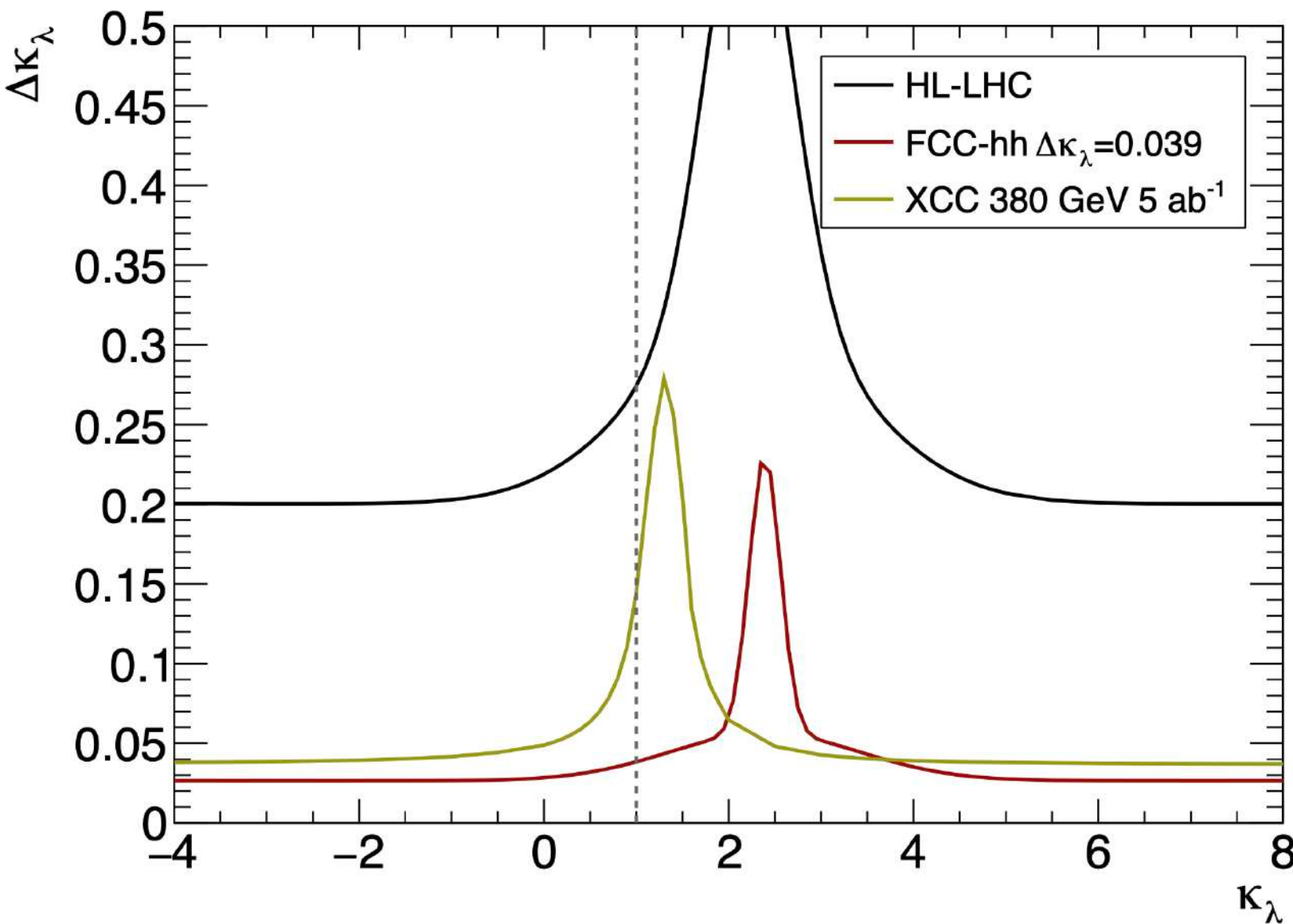


Photon collider - Detector analysis

[S. A. Castelazo, U. S. Qureshi, T. Barklow and A. Schwartzman '26]

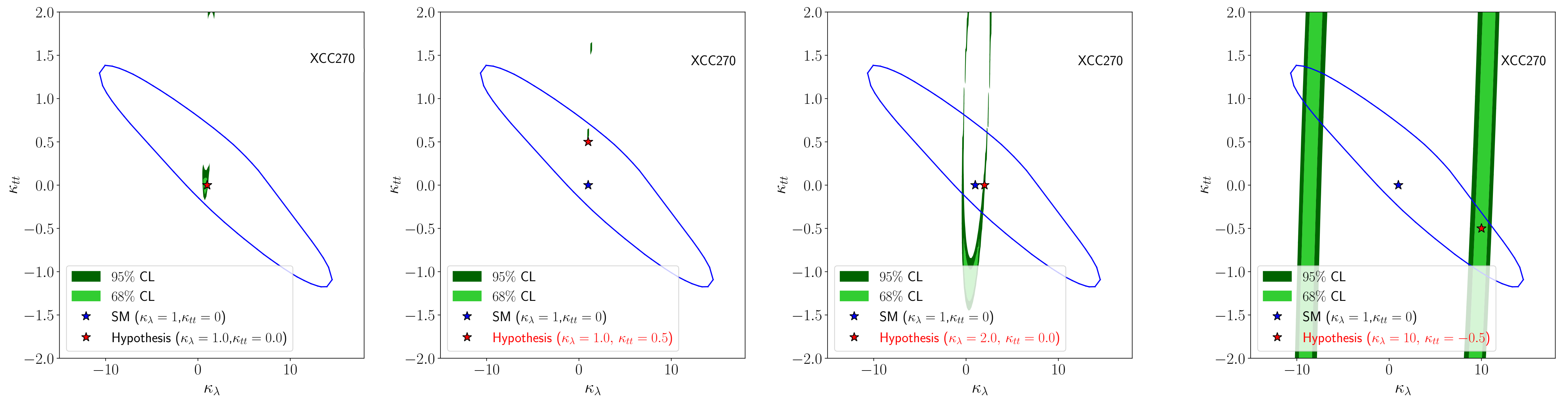
Process	Initial Events	After Preselection	After GA
$\gamma\gamma \rightarrow HH \rightarrow b\bar{b}b\bar{b}$	1,812	703	151
$\gamma\gamma \rightarrow W^+W^- \rightarrow q\bar{q}q\bar{q}$	3,813,000	4,717	2
$\gamma\gamma \rightarrow t\bar{t}$	2,866,000	34,695	29
$\gamma\gamma \rightarrow ZZ \rightarrow q\bar{q}q\bar{q}$	1,378,000	98,527	8
$\gamma\gamma \rightarrow q\bar{q}$	307,700	11,112	6
$\gamma\gamma \rightarrow ZH \rightarrow q\bar{q}H$	8,202	1,004	9
$e\gamma \rightarrow eq\bar{q}$ or $\nu q\bar{q}$	41,195	1,206	1
$e\gamma \rightarrow eq\bar{q}q\bar{q}$ or $\nu q\bar{q}q\bar{q}$	7,681	137	0
$e\gamma \rightarrow eq\bar{q}H$	3,282	217	0
$e^+e^- \rightarrow q\bar{q}$	753,615	54,067	12
$e^+e^- \rightarrow W^+W^-$ or $ZZ \rightarrow q\bar{q}q\bar{q}$	152,212	14,121	2
$e^+e^- \rightarrow ZH \rightarrow q\bar{q}H$	123,698	13,771	62
$e^+e^- \rightarrow t\bar{t}$	57,001	648	1

Table 3. Number of events after GA cuts on the BDT output distributions for the signal and backgrounds.



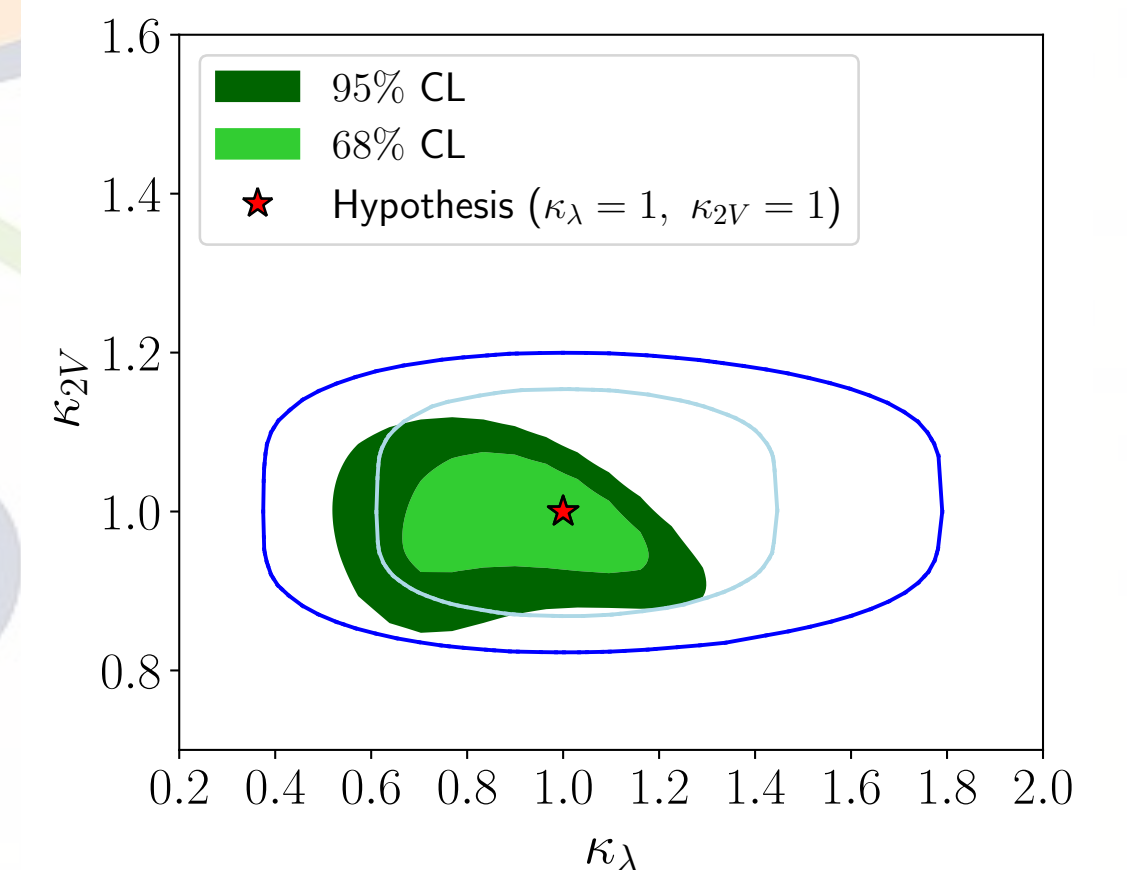
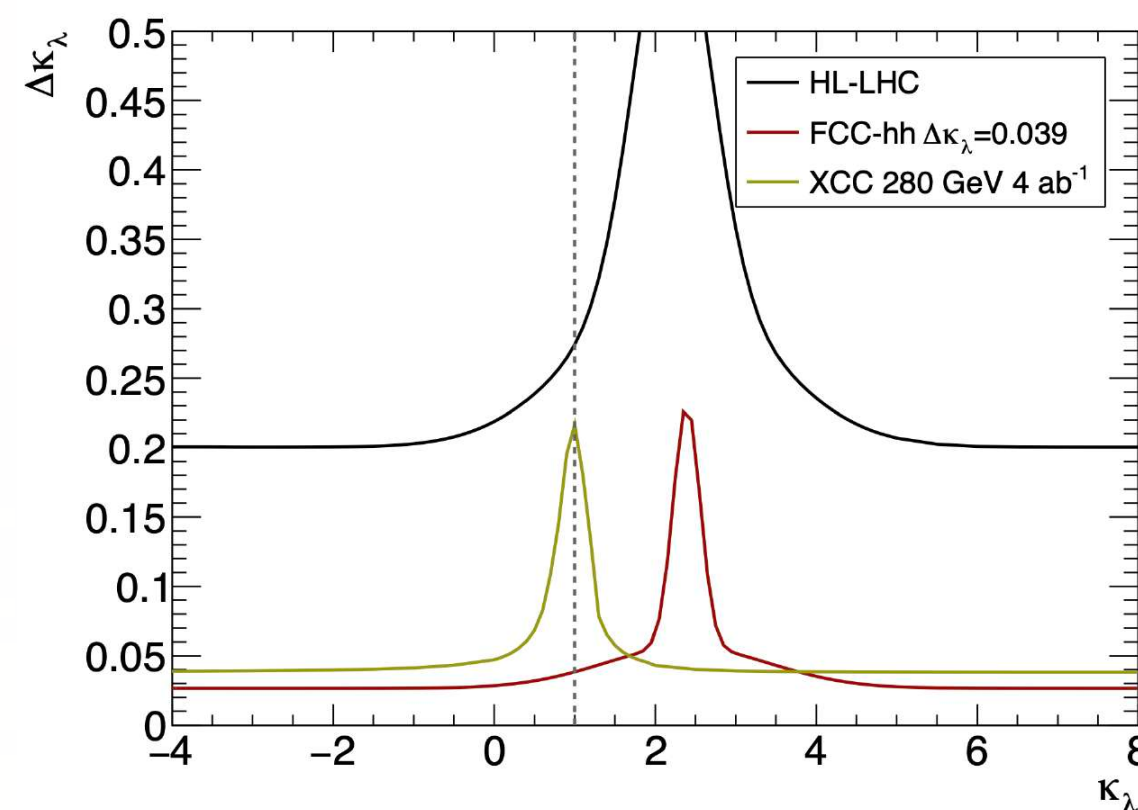
Photon collider - κ_{tt} dependency

- The full XCC270 results in the $\kappa_\lambda/\kappa_{tt}$ plane
- The tight contours lead to very small exclusion ranges
- Slim connected areas between the limits not shown due to plotting limitation



Summary

- Photon colliders offer a vast physics program at the lowest centre-of-mass energies
- Two types of photon colliders: optical collider and the XCC, both offer similar physics reach but vastly different underlying technology
- Combining two energy searches for Higgs pair production offeres the best sensitivity and combining with HL-LHC projected limits gives $\Delta\kappa_\lambda = -27\%, +10\%$ for the optical collider and $\Delta\kappa_\lambda = -17\%, +3\%$ for the XCC at $\kappa_{2V} = 1, \kappa_{tt} = 0$
- Currently still a small community always looking for interested people



Heraeus Photon Collider Workshop

Hosted by the University of Hamburg and DESY, 17-19 March 2027



- Higgs and BSM precision physics at photon colliders
- Physics potential, new BSM searches, event generators, progress in lasers, accelerators and detectors
- **LOC:**
G. Moortgat-Pick (UHH)
J. Braathen (DESY)
- Registration will open in late October/early November
Details and Indico page will follow
- **Please mark your calendars!**



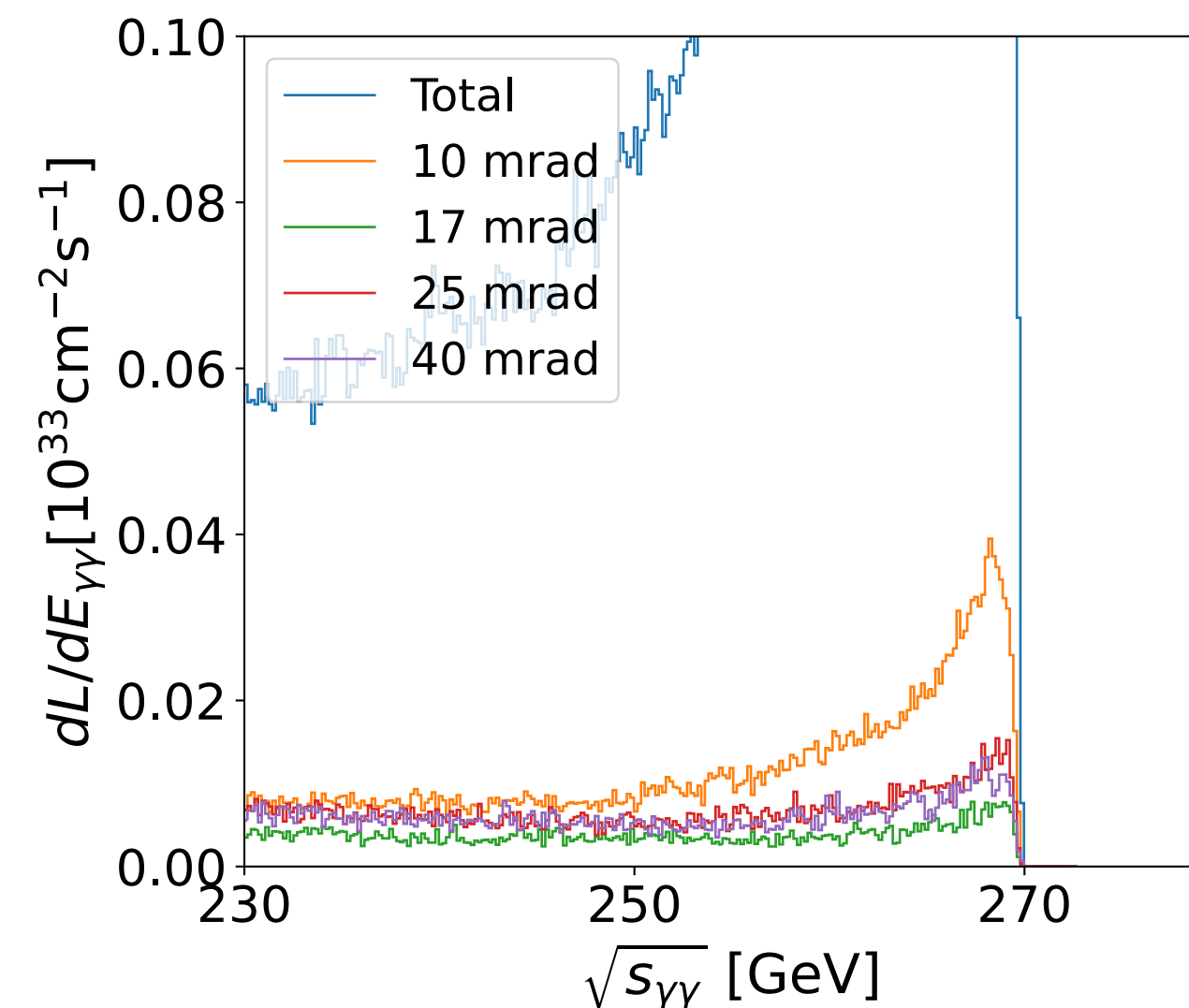
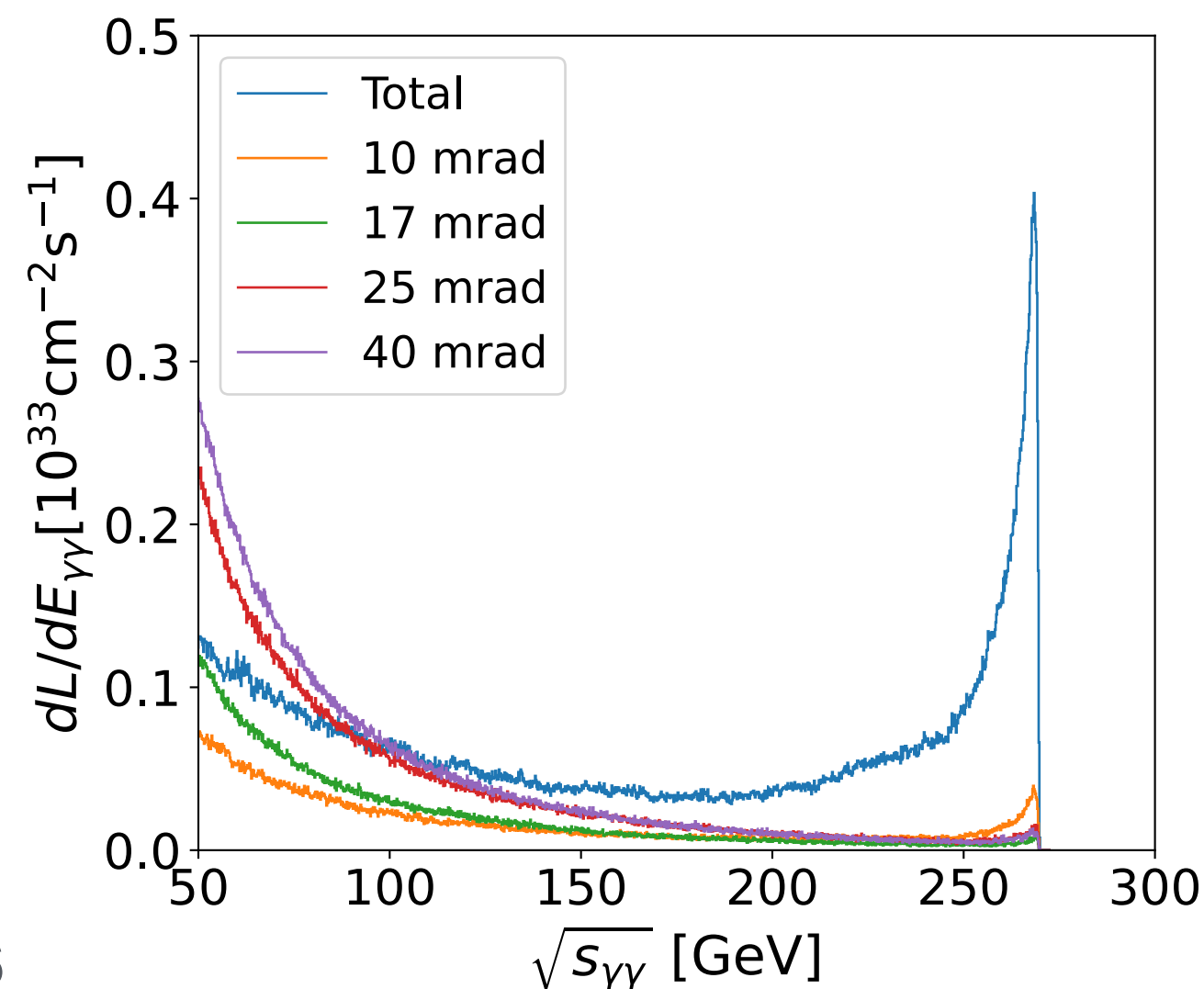
皆様のご参加をお待ちしています!



Thank you for listening

Photon collider - XCC open questions

- The XCC design currently operates with a small crab crossing angle (2 mrad)
- Most future colliders plan for a crab crossing angle of 14-20 mrad
- Can use larger crab crossing for the XCC design, but requires 4 times the laser intensity
- Additionally the angle between the laser and electron beam at conversion point requires a full study within a test facility, due to changes in the angle leading to large drops in luminosity



Why photon colliders?

- Photon collider would be ~60% of the cost of a comparable e⁺e⁻-collider [T. Barklow et al. '23]
- Requires detailed study for modern optical collider
- No need for a positron source allows for large savings

Compared to the capital cost of \$3.7B for one specific C³-250 scenario [M. Bai et al. '21]

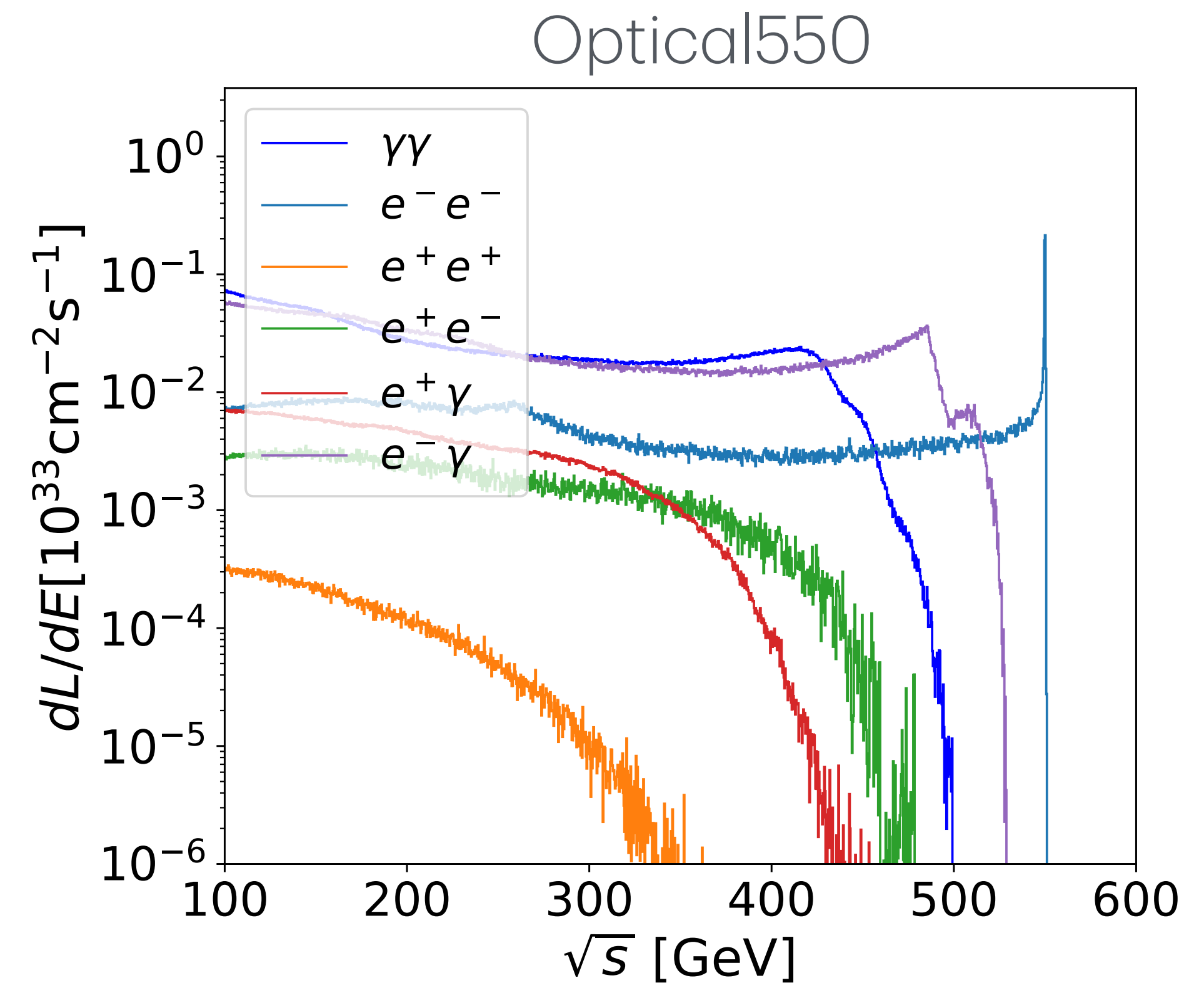
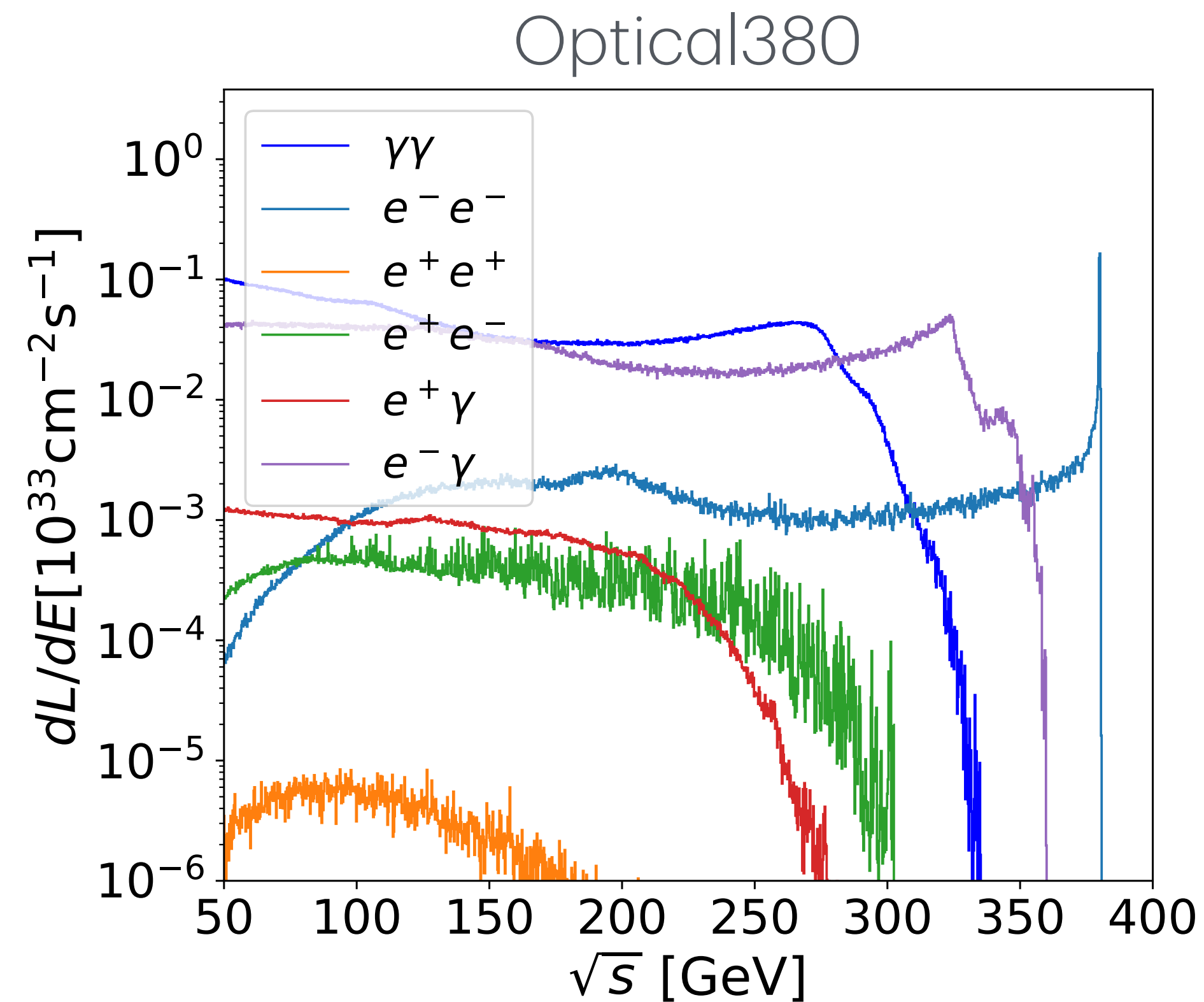
[T. Barklow et al. '23]

	Sub-Domain	%	%
Sources	Injectors	9	26
	FEL	9	
	Beam Transport	9	
Main Linac	Cryomodule	9	30
	C-band Klystron	22	
BDS	Beam Delivery and Final Focus	7	15
	IR	8	
Support Infrastructure	Civil Engineering	5	28
	Common Facilities	18	
	Cryo-plant	6	
Total	2.3B\$	100	100

Table 5. Initial estimate of XCC cost breakout using the C³-250 cost model.

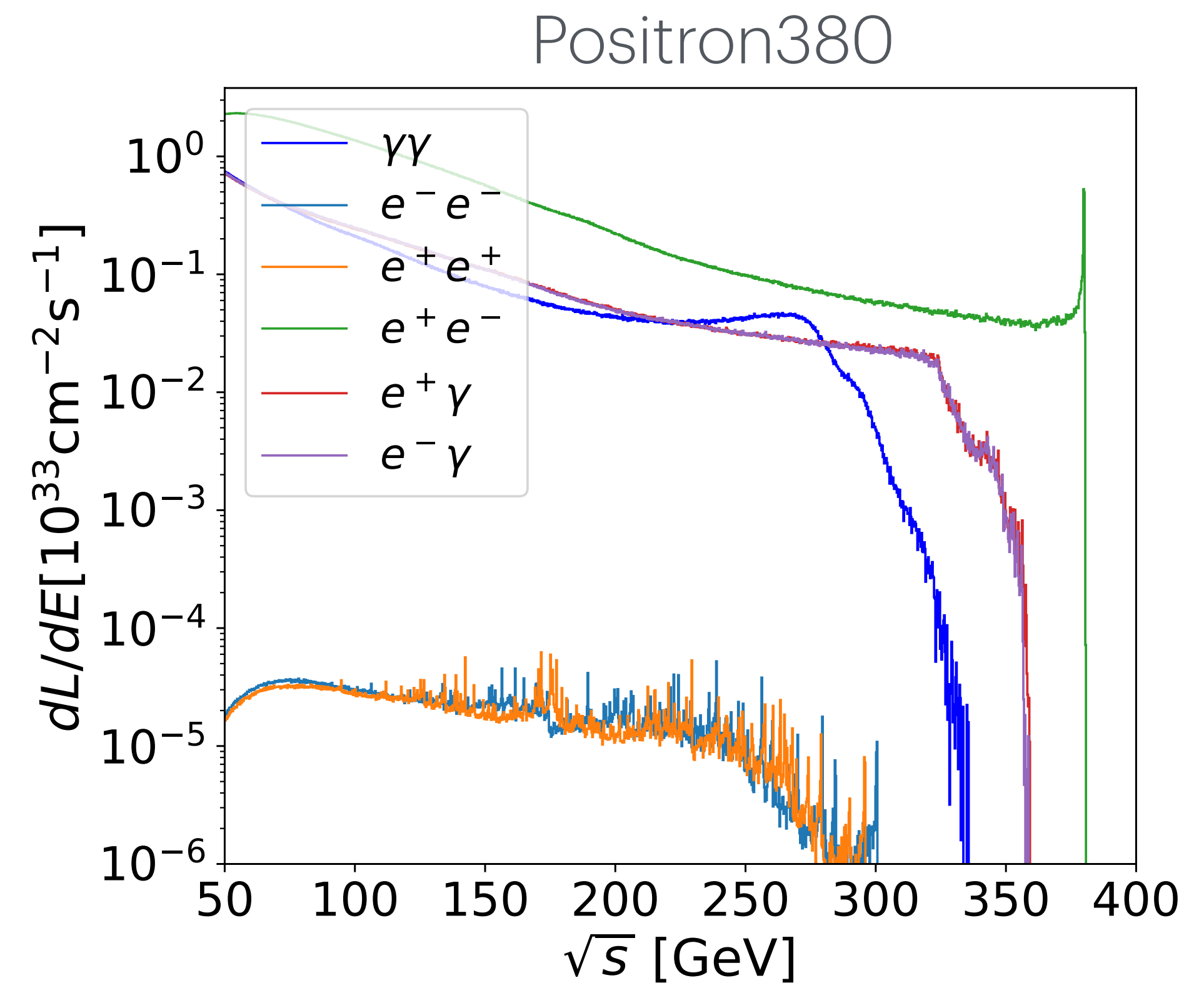
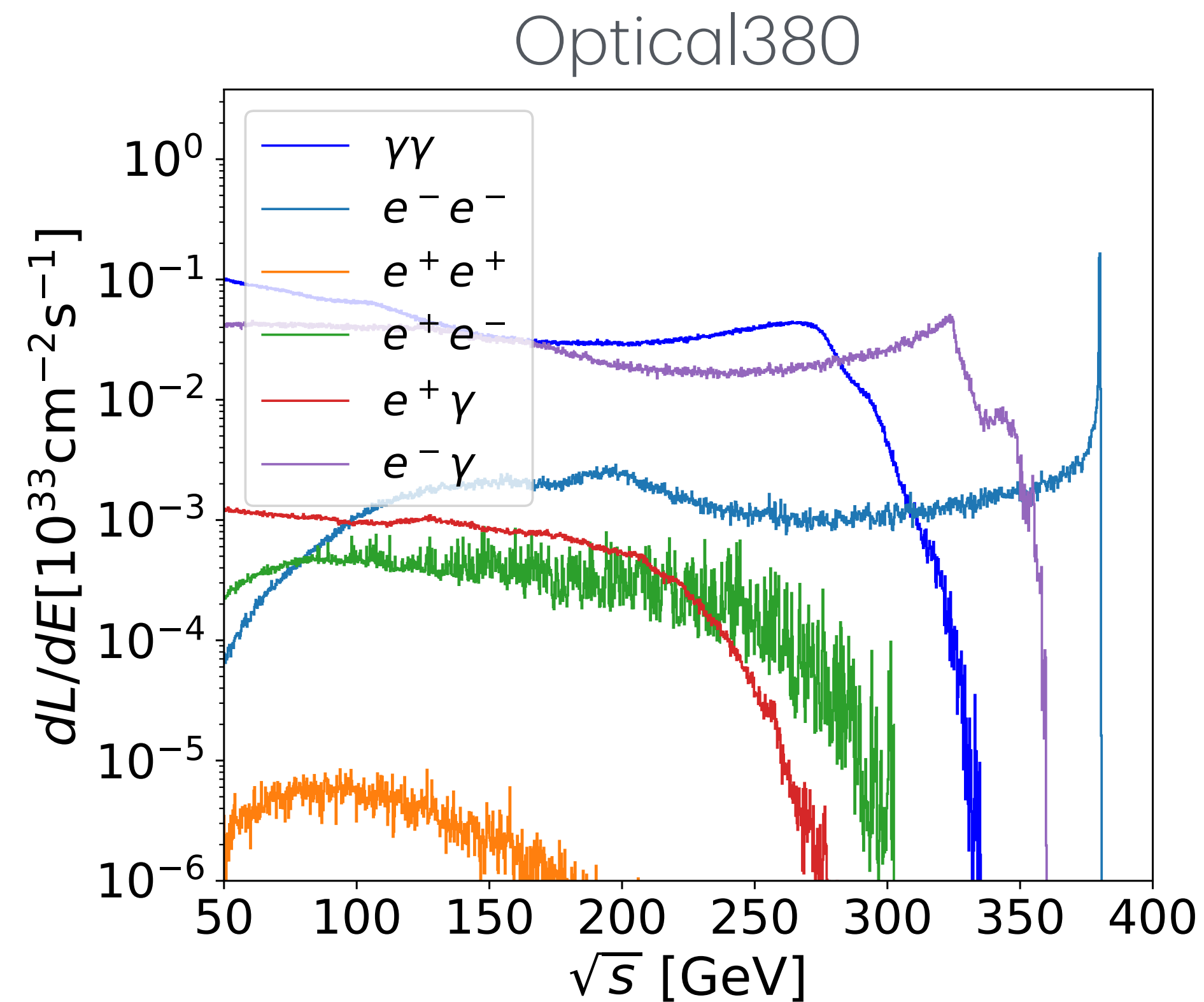
Photon collider - Optical background

- Not all electrons are converted, leading to different processes in the same interaction point



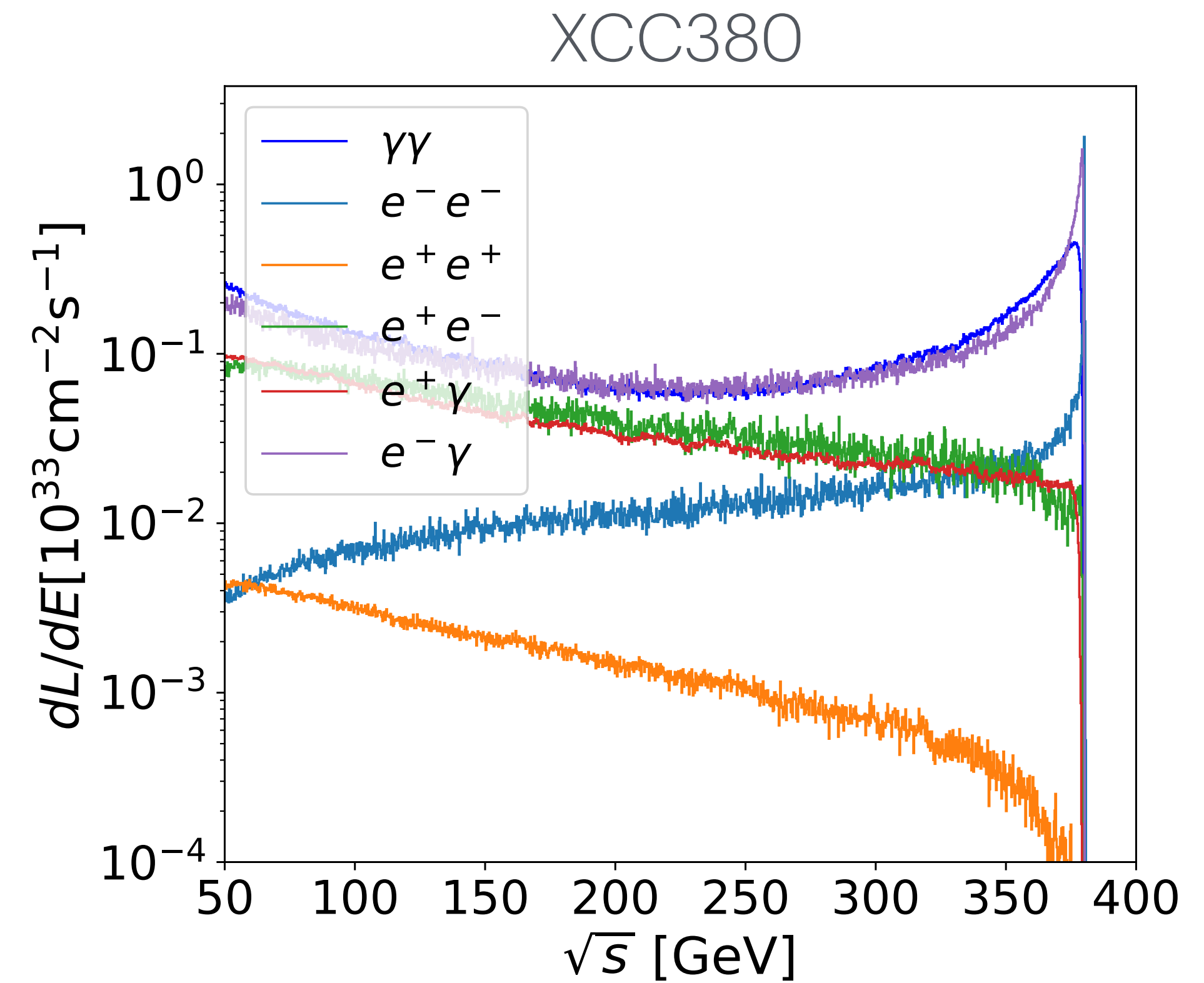
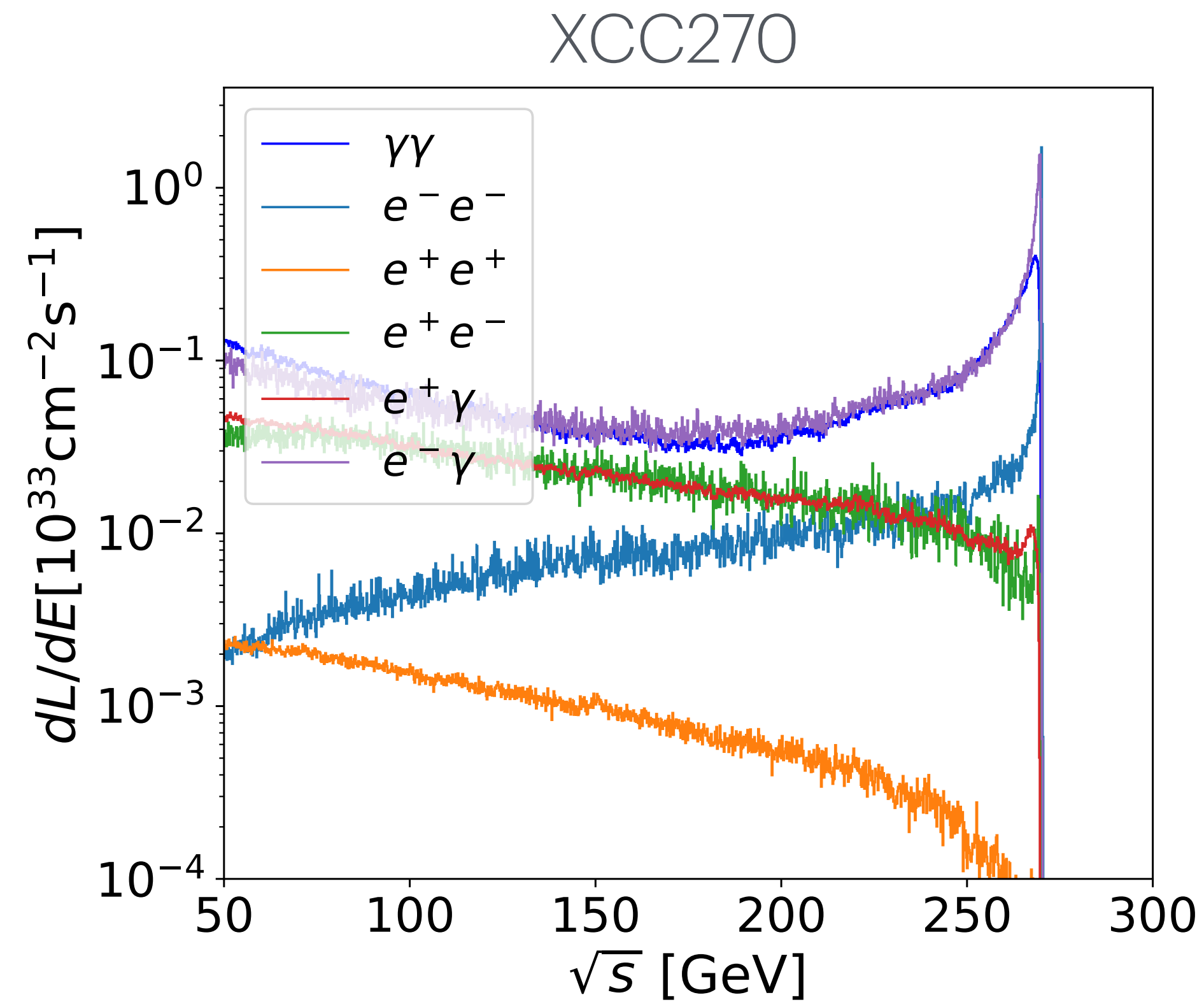
Photon collider - Positron background

- Not all electrons are converted, leading to different processes in the same interaction point



Photon collider - XCC background

- Not all electrons are converted, leading to different processes in the same interaction point

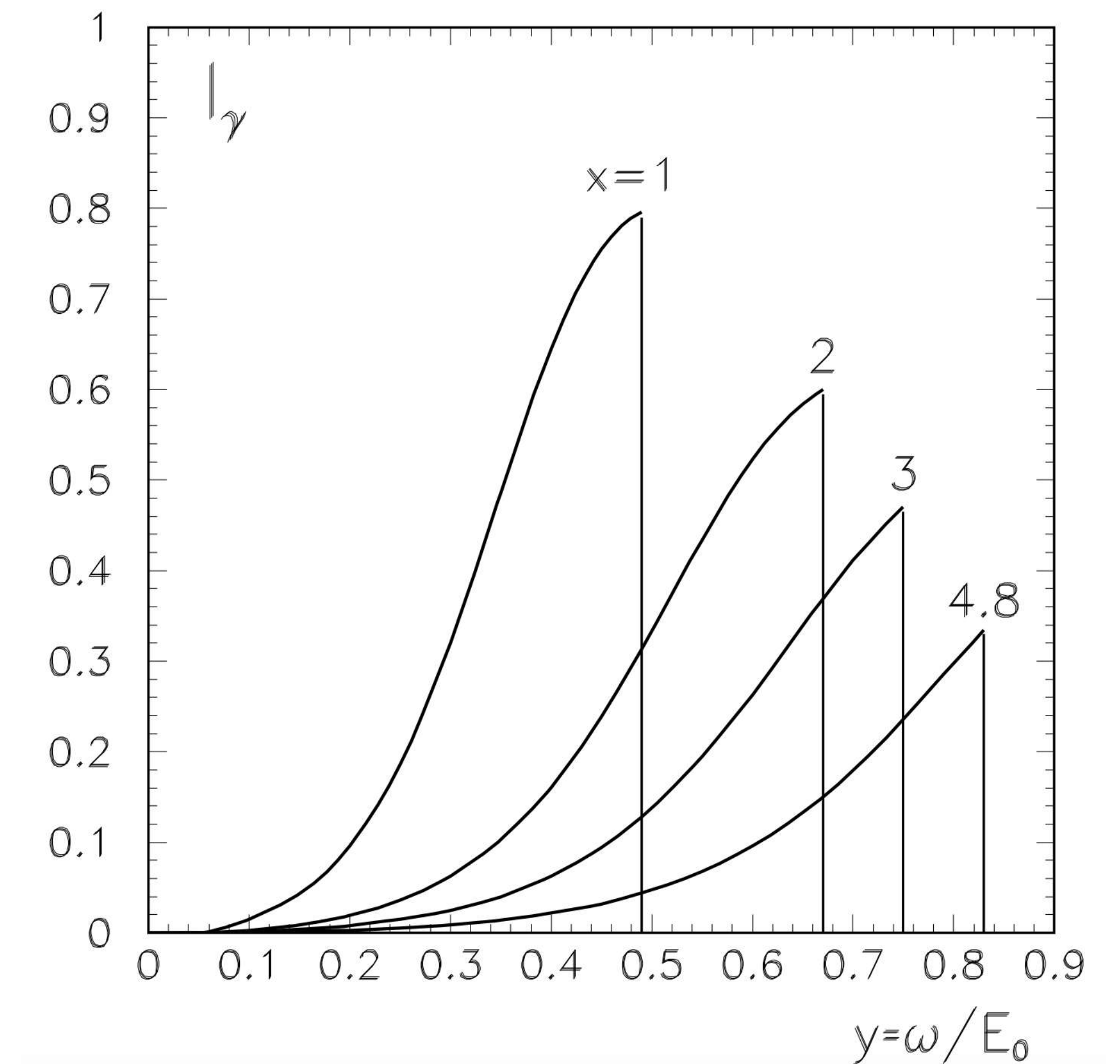


Photon collider - Linear polarisation

- Scattered photons can also carry linear polarisation

$$\langle l_\gamma \rangle = \frac{2r^2 P_l}{(1-y)^{-1} + 1 - y - 4r(1-r) - 2\lambda_e P_c x r(2-y)(2r-1)}$$

- Very useful for CP studies of the Higgs
- Could increase the reach of ALPs searches



Photon collider - How to get the Beams

- Compton backscattering $\sigma_c = \sigma_c^{np} + \lambda_e P_c \sigma_c^p$, where

$$\sigma_c^{np} = \frac{2\pi\alpha^2}{xm_e^2} \left[\left(1 - \frac{4}{x} - \frac{8}{x^2} \right) \ln(x+1) + \frac{1}{2} + \frac{8}{x} - \frac{1}{2(x+1)^2} \right],$$

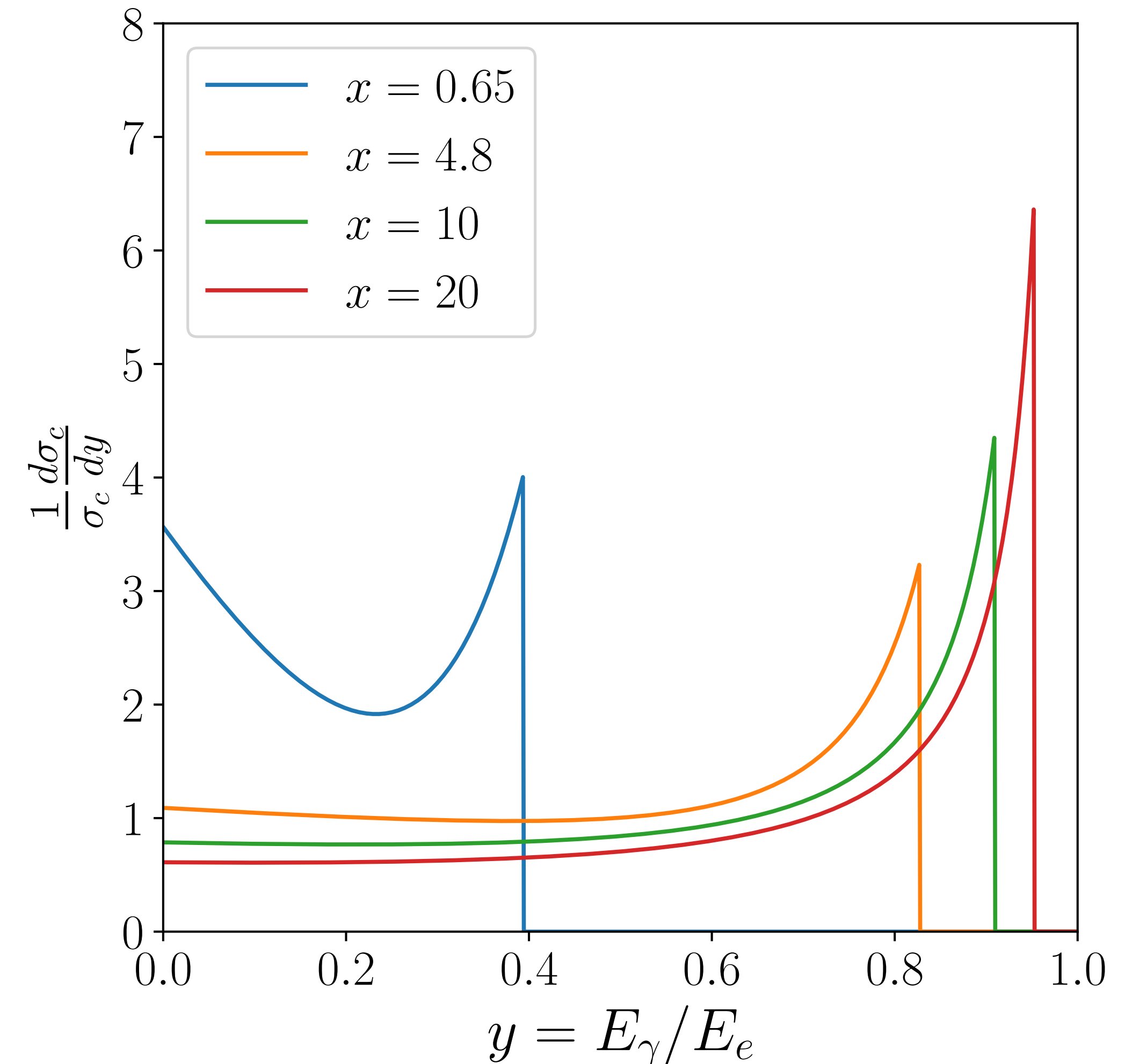
$$\sigma_c^p = \frac{2\pi\alpha^2}{xm_e^2} \left[\left(1 + \frac{2}{x} \right) \ln x + 1 - \frac{5}{2} + \frac{1}{x+1} - \frac{1}{2(x+1)^2} \right]$$

- The energy spectrum of the scattered photon beam

$$\frac{1}{\sigma_c} \frac{d\sigma_c}{dy} \equiv f(x, y)$$

$$= \frac{2\pi\alpha^2}{\sigma_c xm_e^2} \left[\frac{1}{1-y} + 1 - y - 4r(1-r) - \lambda_e P_c r x (2r-1)(2-y) \right]$$

$$r = \frac{y}{x(1-y)}$$



Photon collider - Additional parameters

- Non-linear QED effects, leads to additional harmonics

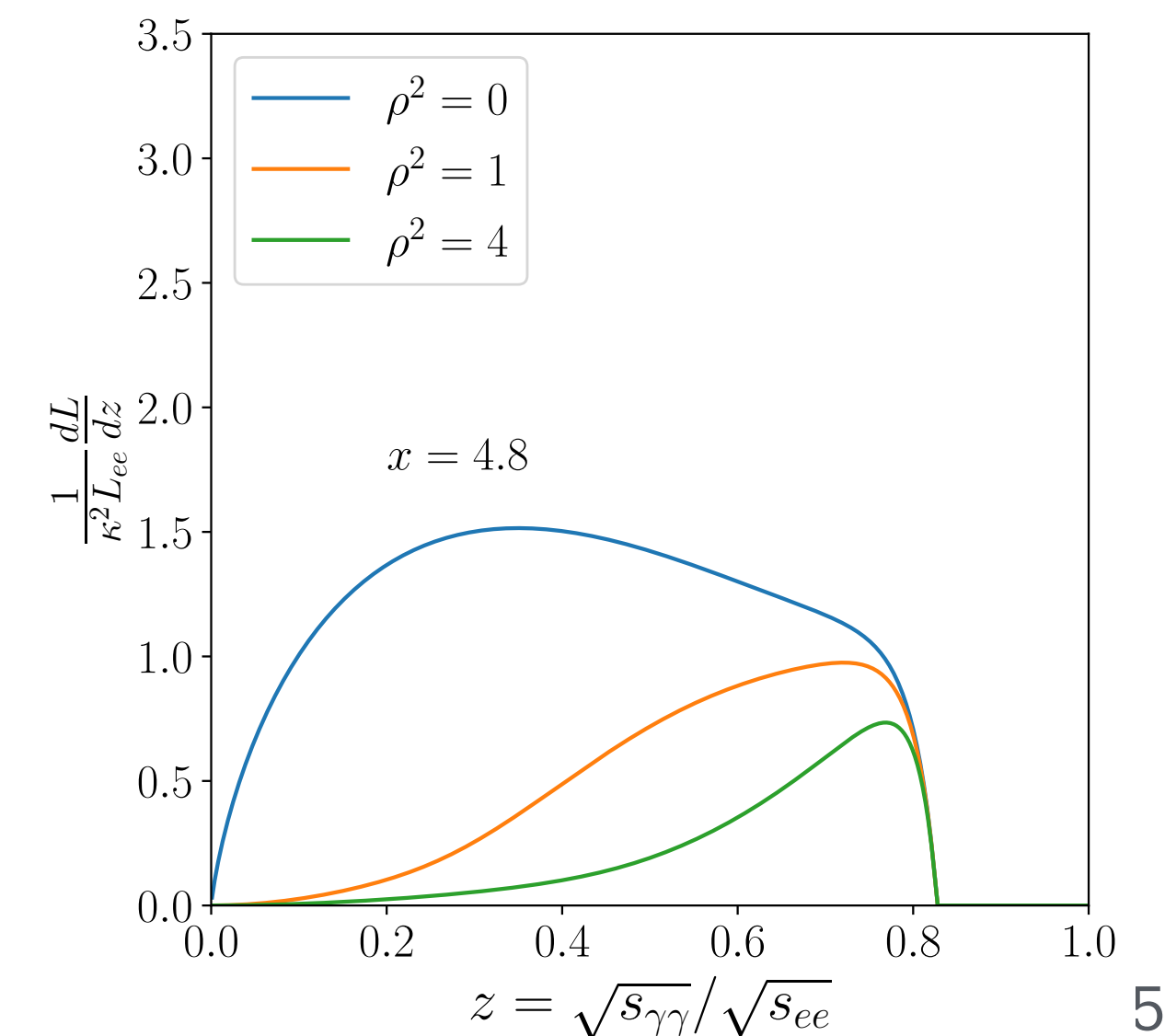
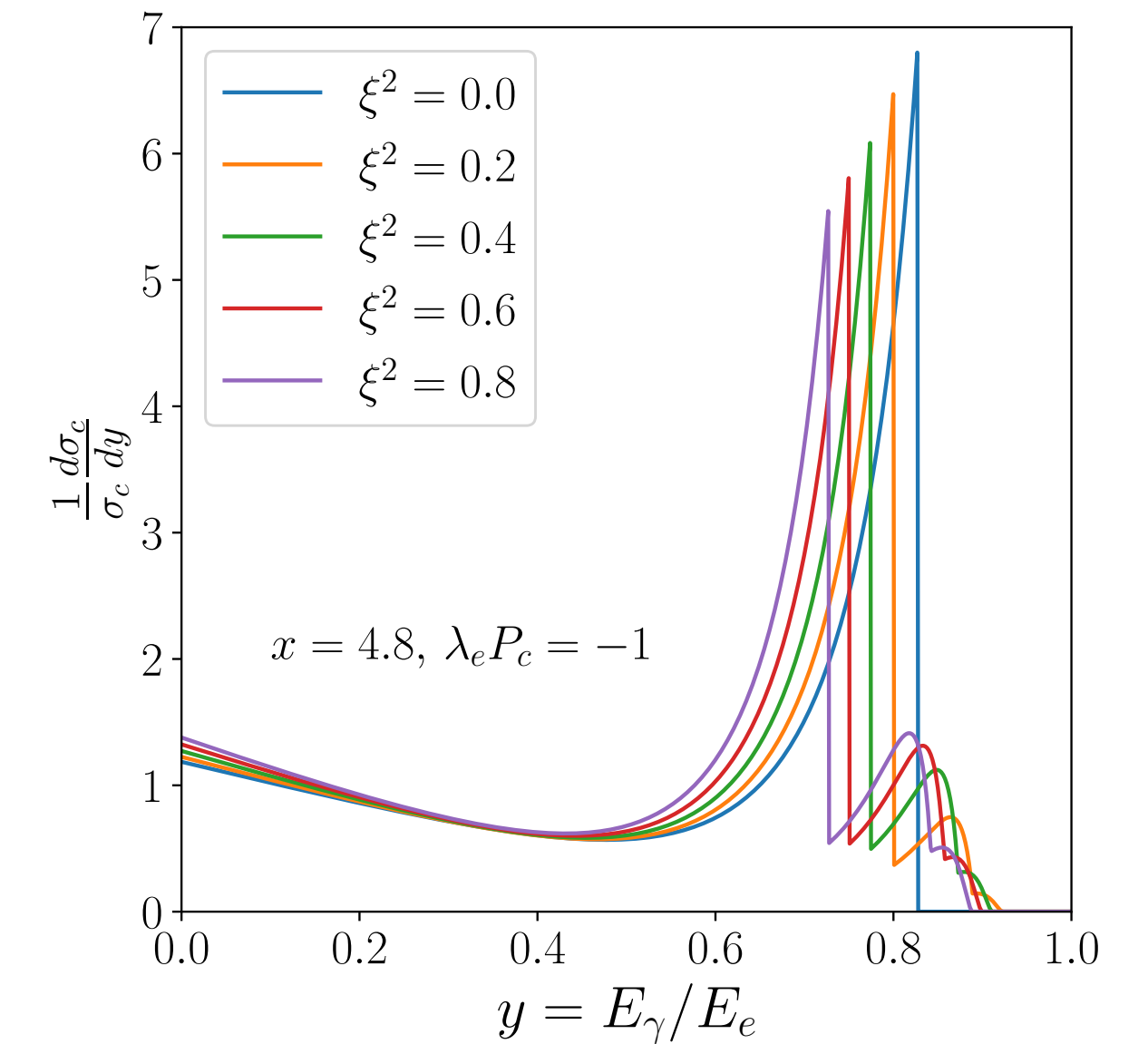
$$\xi^2 = \frac{e^2 \bar{F}^2 \hbar^2}{m_e^2 c^2 \omega_0^2} = \frac{2n_\gamma r_e^2 \lambda}{\alpha} = \left(\frac{\lambda}{2\pi m_e} \right)^2 \mu_0 c P$$

- The distance b between the conversion and interaction point

$$\rho^2 = \left(\frac{b}{\gamma \sigma_x} \right)^2 + \left(\frac{b}{\gamma \sigma_y} \right)^2, \text{ with } \sigma_{x/y} \text{ the transverse electron beam size}$$

$$\frac{1}{L_{\gamma\gamma}} \frac{dL_{\gamma\gamma}}{dz} = \frac{1}{k^2 L_{ee}} \frac{dL_{\gamma\gamma}}{dz} = 2z \int_{z^2/y_{\max}}^{y_{\max}} \frac{dy}{y} f(x, y) f\left(x, \frac{z^2}{y}\right) * F(\rho),$$

$$F(\rho) = I_0 \left(\rho^2 \sqrt{\left(\frac{y_{\max}}{y} - 1 \right) \left(\frac{y_{\max} y}{z^2} - 1 \right)} \right) \exp \left[-\frac{\rho^2}{2} \left(\frac{y_{\max}}{y} + \frac{y_{\max} y}{z^2} - 2 \right) \right]$$



Photon collider - Flat-top laser

- New photon collider designs use flat-top laser beams, providing a higher effective conversion efficiency compared to Gaussian beam

	Flat-top	Gaussian
$I(r, z = 0)$	$I_0 \left[\frac{2J_1(r/s)}{r/s} \right]^2$	$I_0 \exp \left(-2 \frac{r^2}{\omega^2(z)} \right)$
	$s = \frac{f_{\#} \lambda}{\pi}$	$\omega(z) = \omega_0 \left[1 + \left(\frac{z}{z_R} \right)^2 \right]^{1/2}$
$I(r = 0, z)$	$I_0 \left[\frac{\sin(z/z_0)}{z/z_0} \right]^2$	$I_0 \left[1 + \left(\frac{z}{z_R} \right)^2 \right]^{-1}$
	$z_0 = \frac{8}{\pi} f_{\#}^2 \lambda$	$z_R = \frac{4}{\pi} f_{\#}^2 \lambda$
Effective area	$\frac{f_{\#}^2 \lambda^2}{\pi}$	$\frac{\pi \omega_0^2}{2} = \frac{2 f_{\#}^2 \lambda^2}{\pi}$

