

ニュートリノ質量分光に向けた コヒーレンスによる二光子対超放射の観測

岡山大学 極限量子研究コア

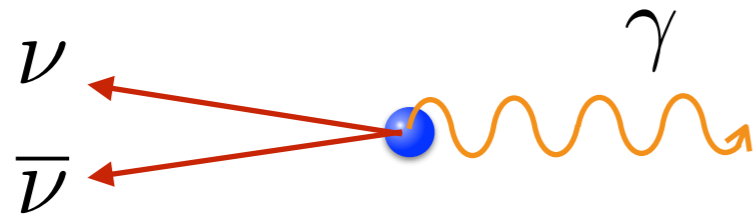
増田孝彦 on behalf of the SPAN collaboration



OKAYAMA UNIV.

- *Motivation* :
Neutrino absolute mass measurement using atomic de-excitation.

REN (Radiative Emission of Neutrino Pair)



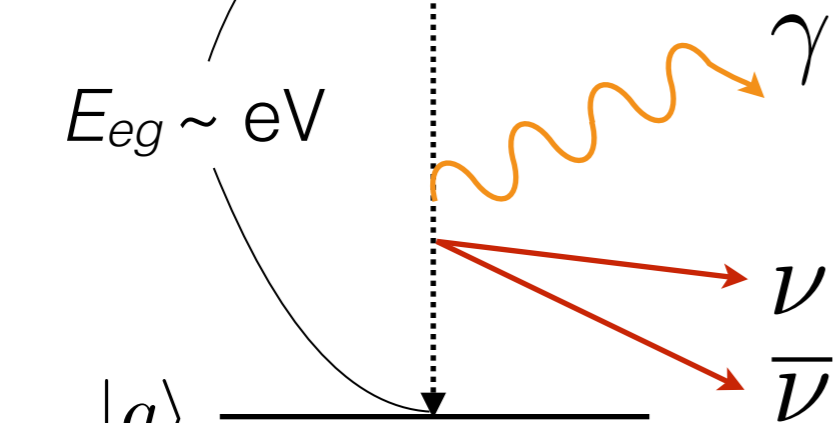
atomic excited state

$|e\rangle$

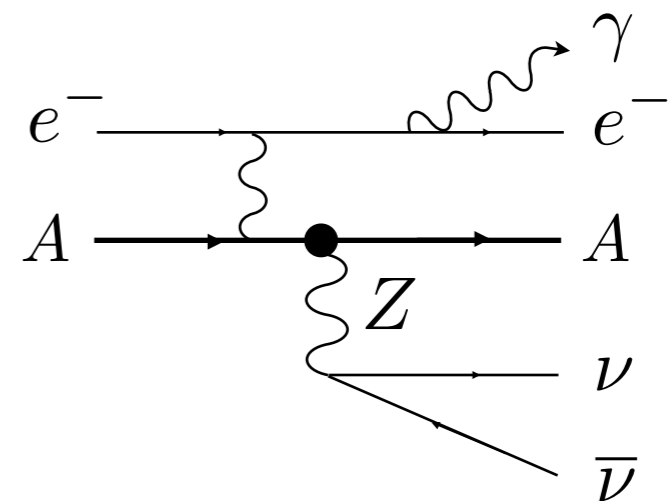
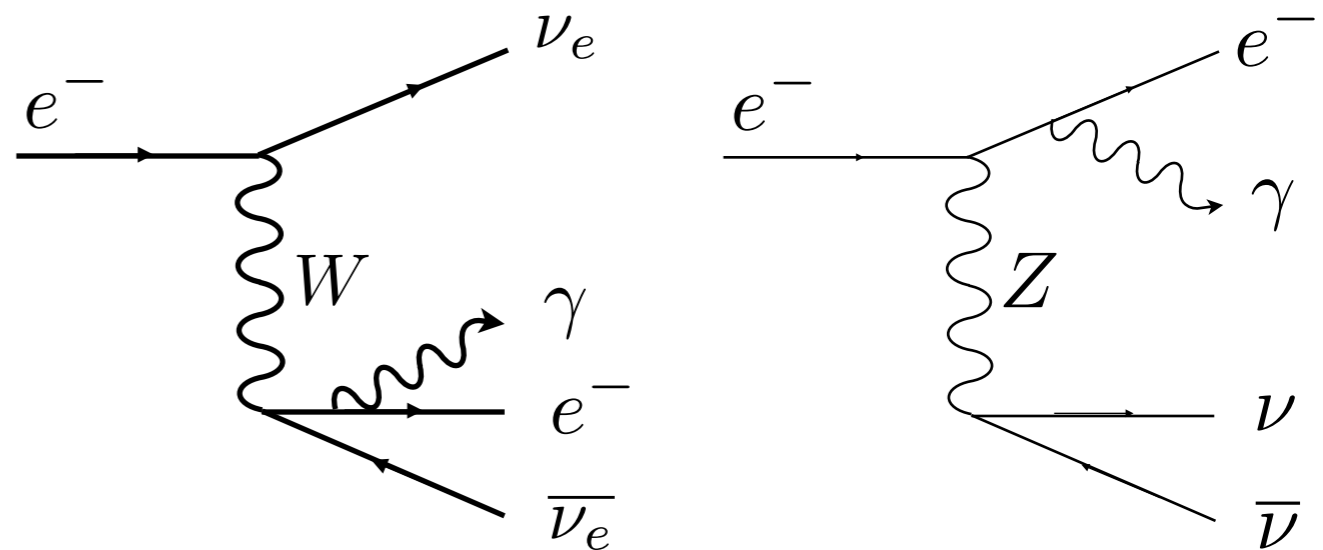
$E_{eg} \sim \text{eV}$

$|g\rangle$

atomic ground state

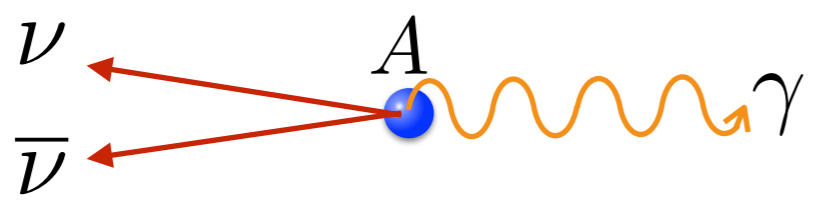


Diagrams of RENP

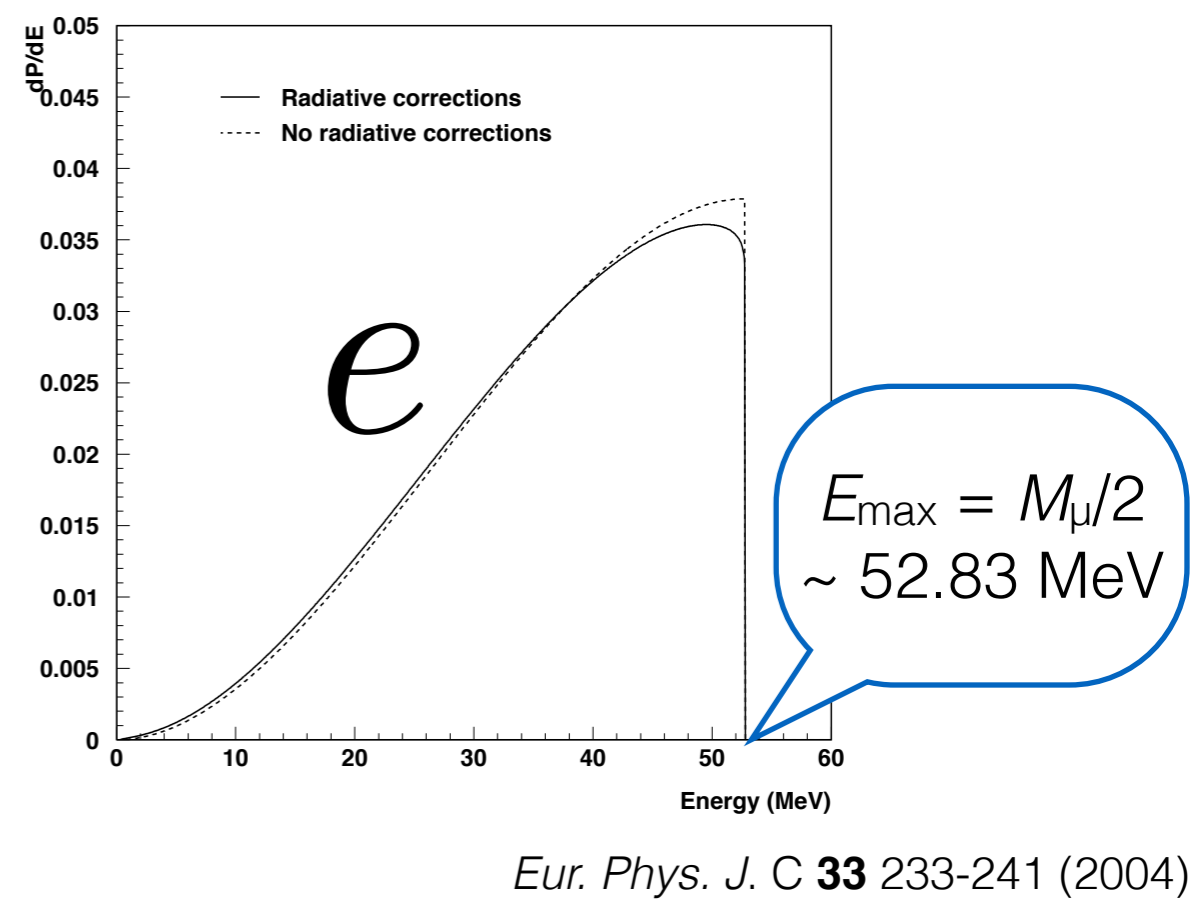
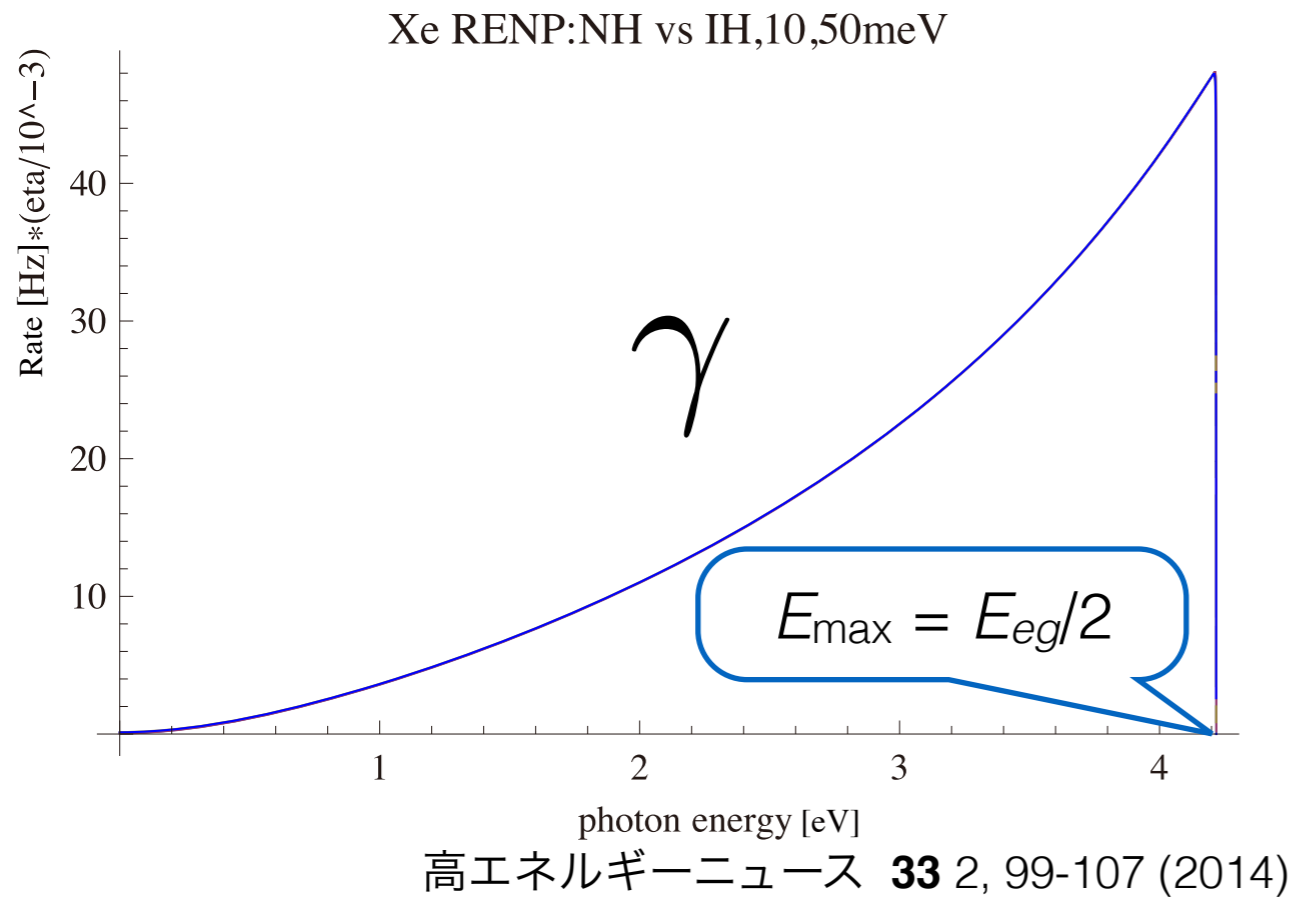
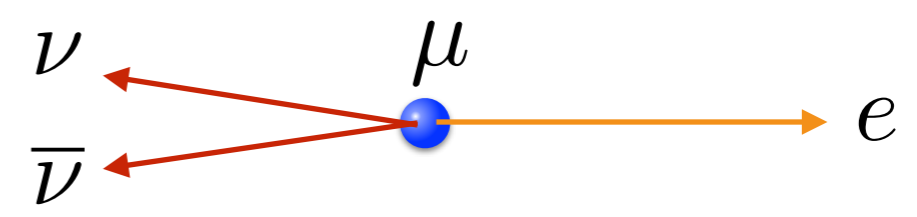


Energy spectrum of emitted photon

$$E_{eg} \rightarrow \gamma \nu_i \bar{\nu}_j$$



$$\mu^- \rightarrow e^- \nu_\mu \bar{\nu}_e$$



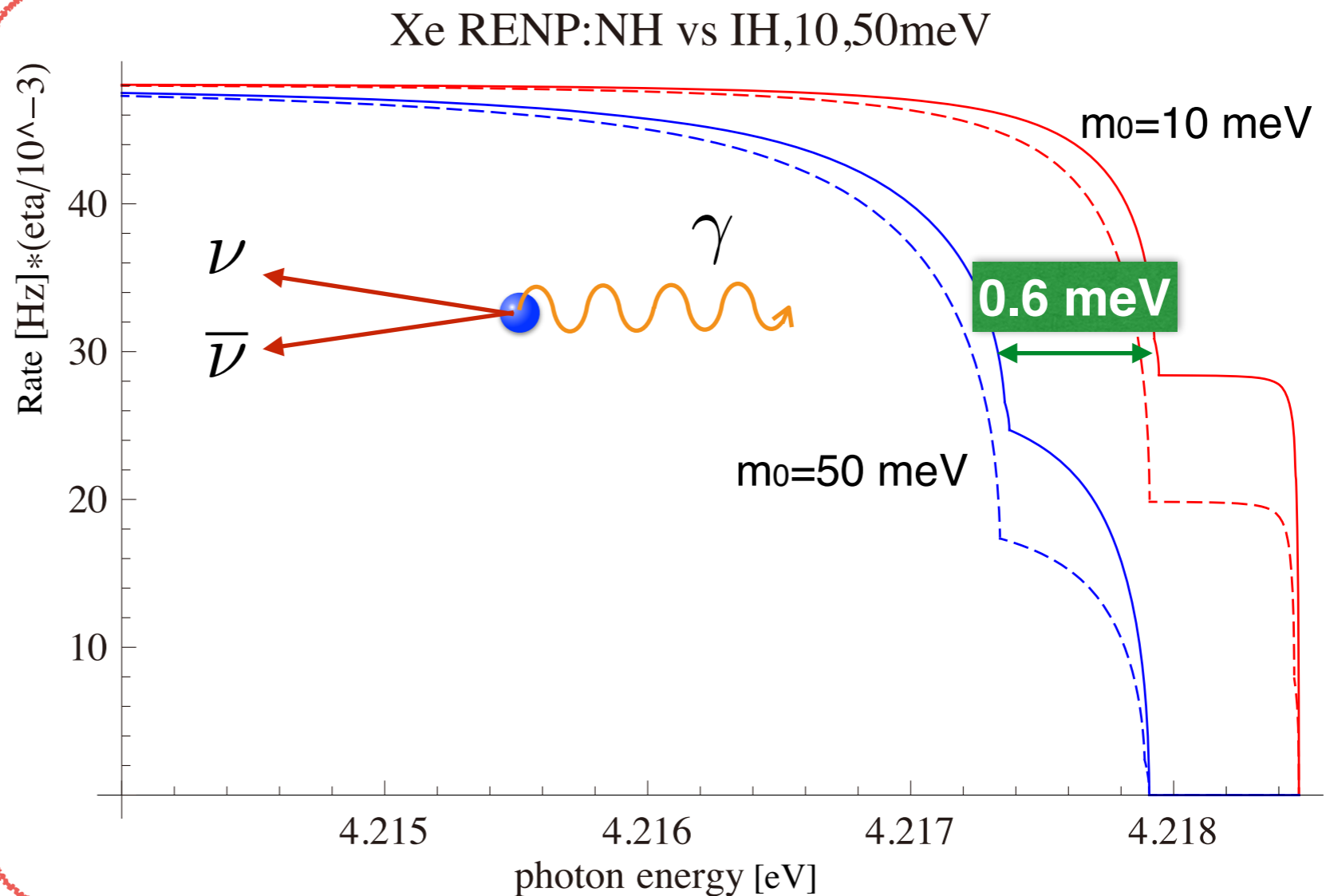
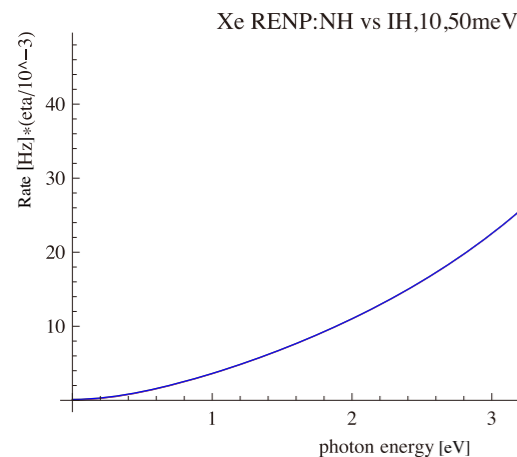
- Photon energy spectrum is similar to Michel Spectrum
 → Maximum value is q-value/2 based on **the massless assumption**

Energy spectrum of emitted gamma

End-points contain the mass information : $\omega_{ij} = \frac{\omega_{eg}}{2} - \frac{(m_i + m_j)^2}{2\omega_{eg}}$

- ☑ $\omega_{eg} \sim m_\nu$ → **Atomic transition**
- ☑ sub meV energy resolution → **Laser spectroscopy**

高エネルギーニューズ 33 2, 99-107 (2014)

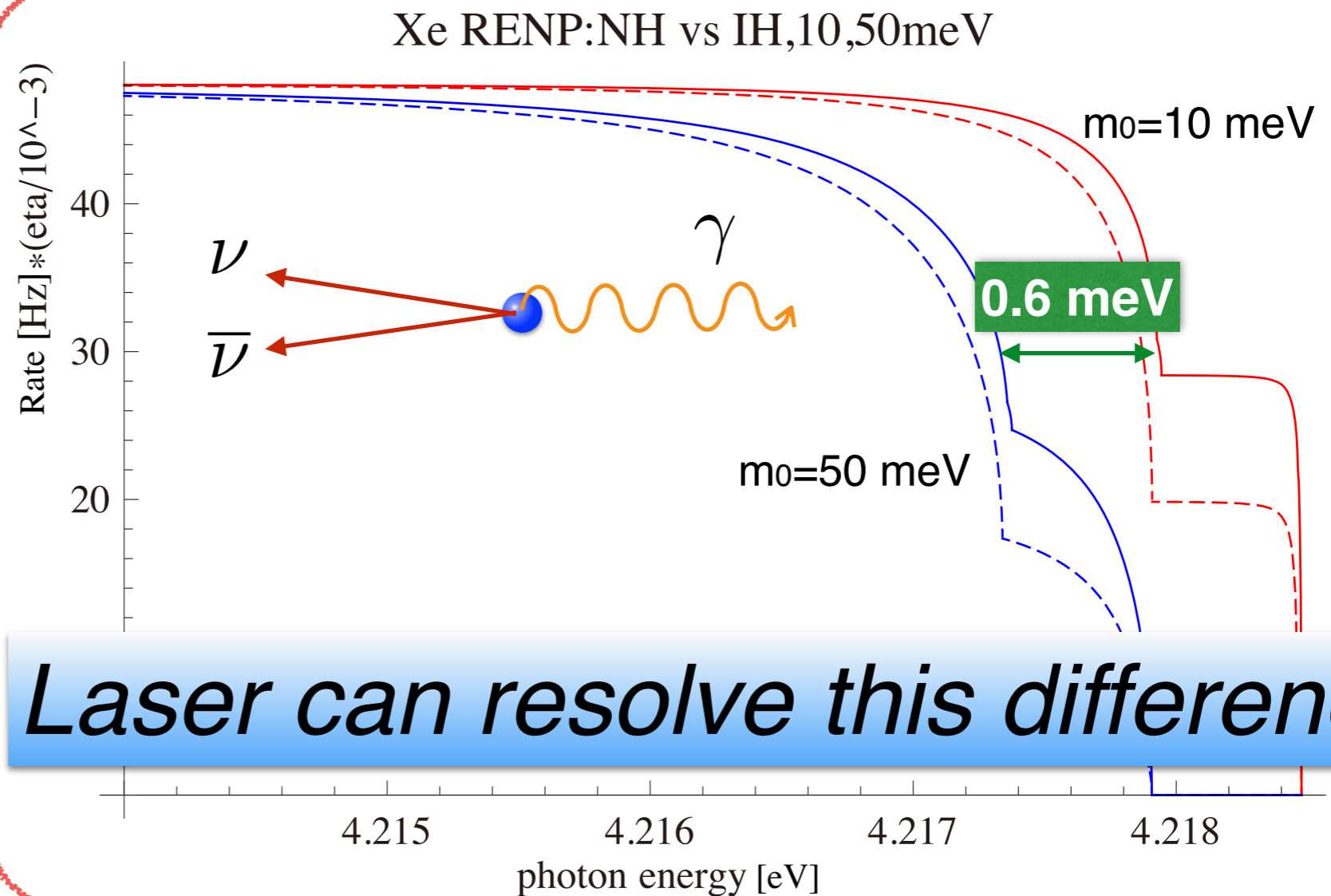
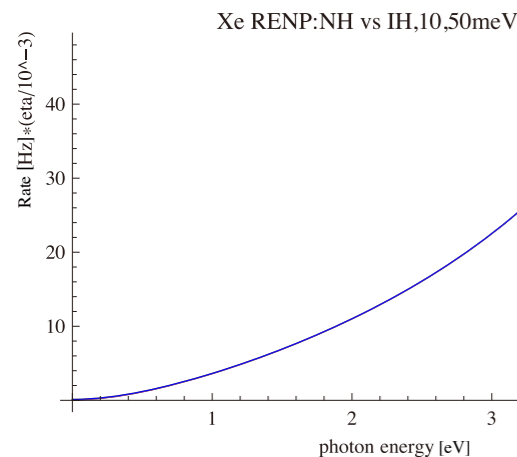


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高エネルギーニュース 33 2, 99-107 (2014)



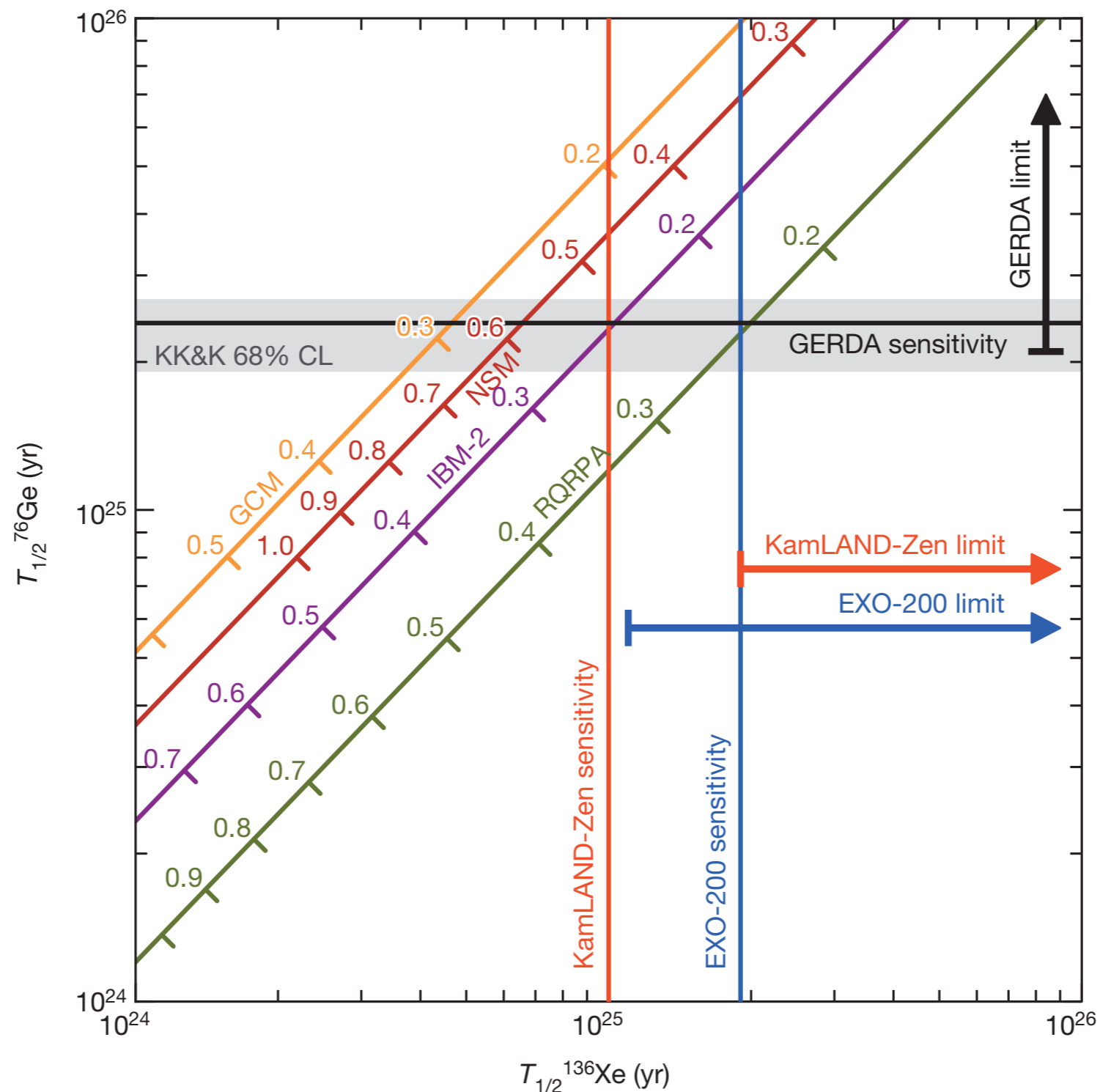
Laser can resolve this difference!

Problem

same as the current sensitivity of $0\nu\beta\beta$

- rare process

i.e. : $E_{eg}=1$ eV $\Rightarrow \Gamma_{\text{REN P}} \sim 1/(10^{26}$ year) for single atom



Nature **510** 229-234 (2014)

Problem

same as the current
sensitivity of $0\nu\beta\beta$

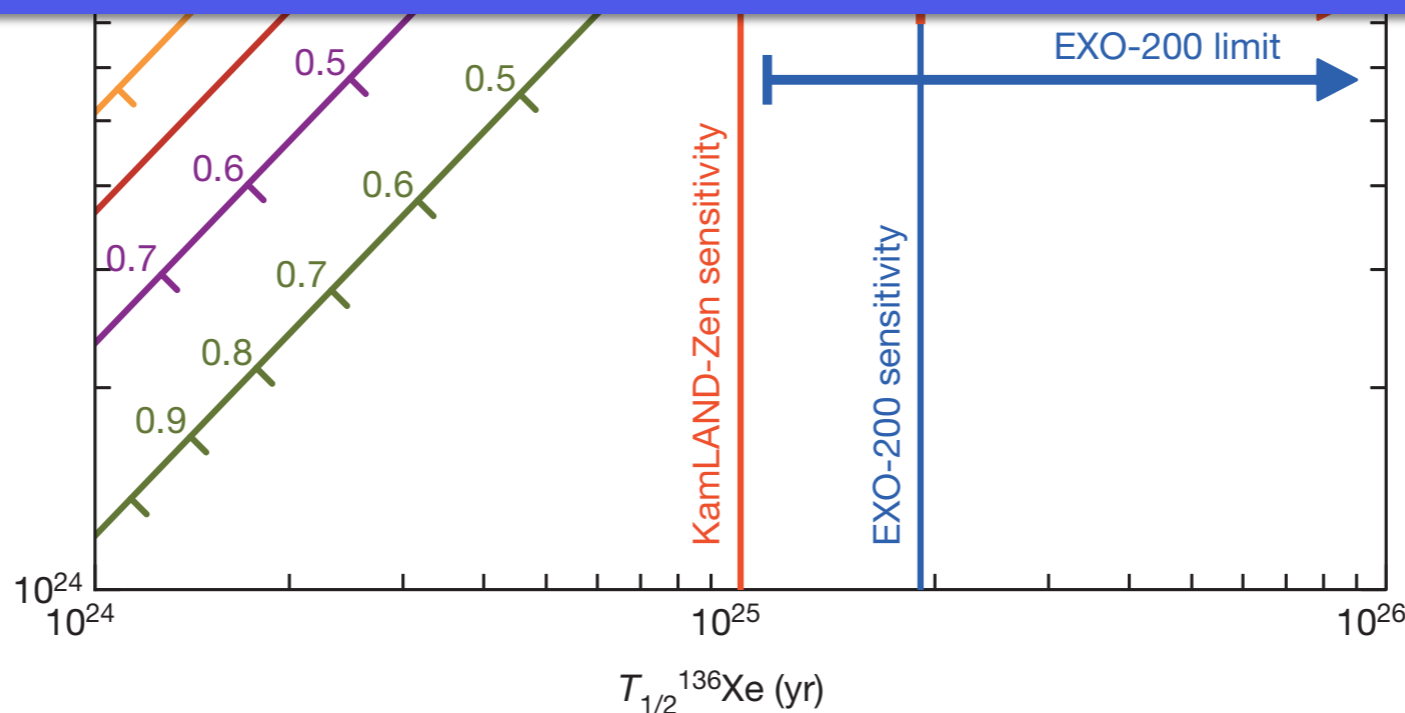
- rare process

i.e. : $E_{eg}=1$ eV $\Rightarrow \Gamma_{\text{REN P}} \sim 1/(10^{26}$ year) for single atom



Should we prepare a huge detector?

*No. Target atoms have to be excited.
We need other method.*

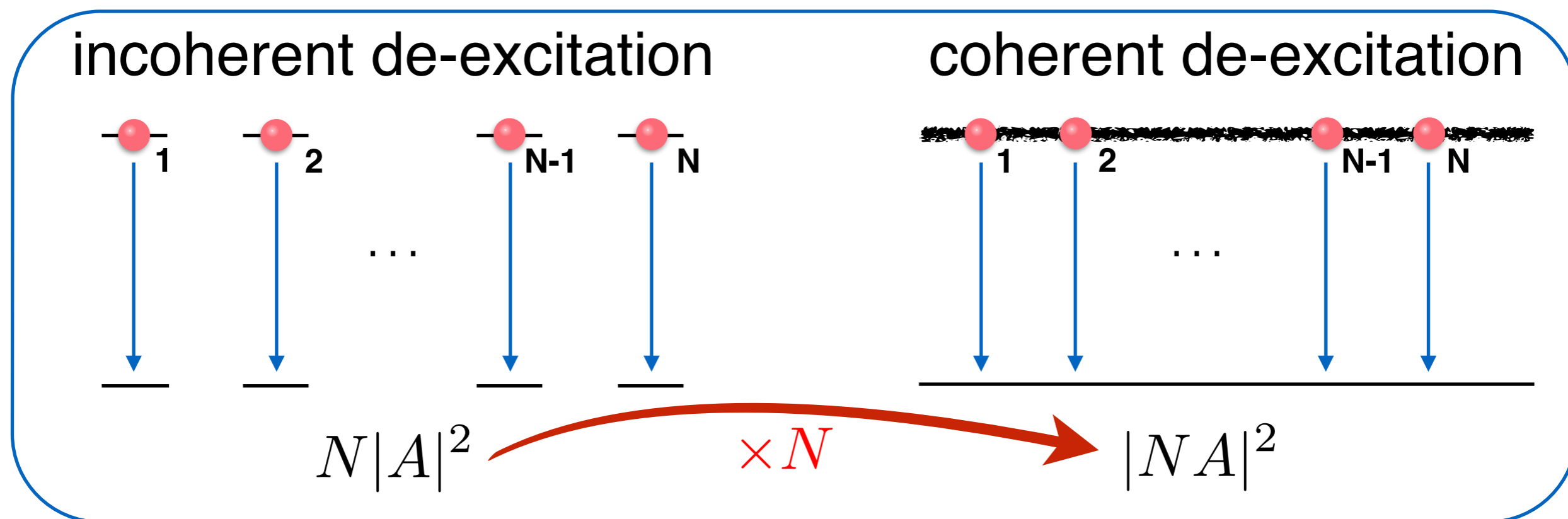


Nature **510** 229-234 (2014)

Macro-coherent Amplification

- **Macro-Coherent Amplification**

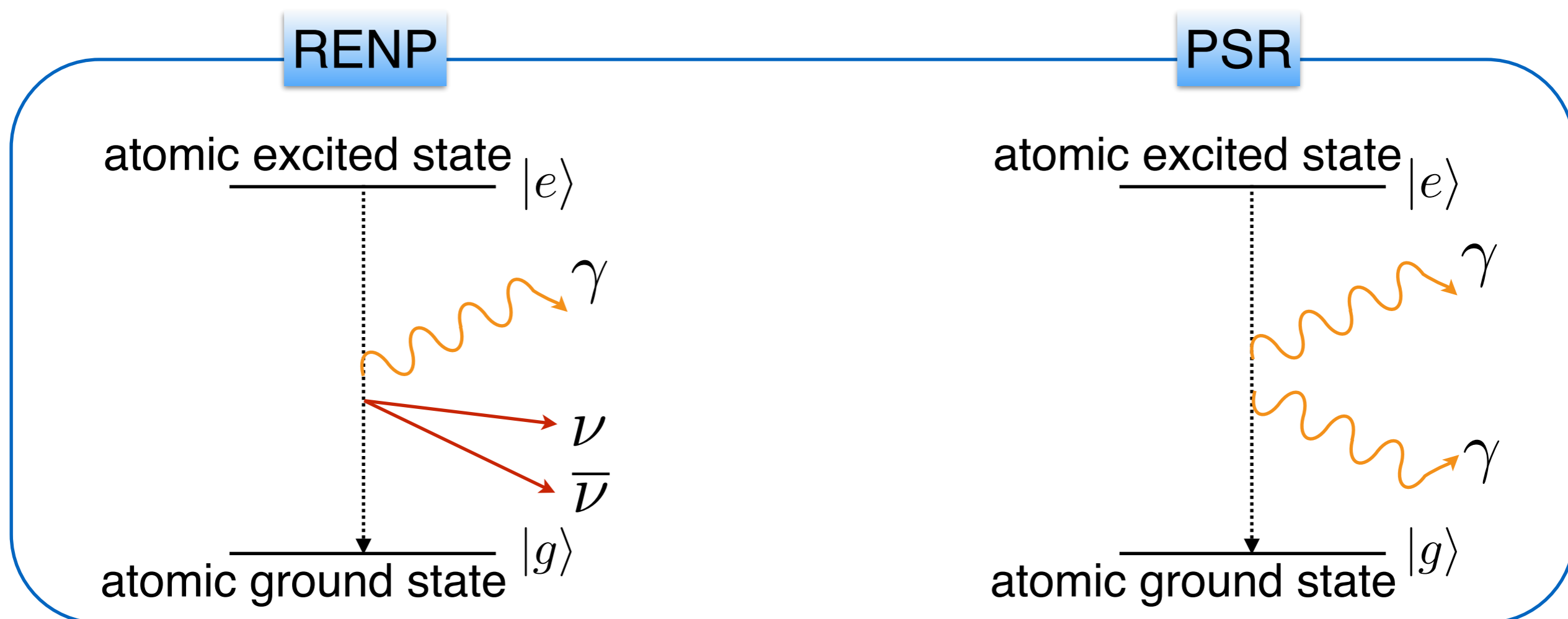
- Coherence in macroscopic size can *enhance* the rate of de-excitation.
 - *Not enlarge the amount of target atoms, but enlarge its transition rate.*



- The mutual phase among N atoms have to be same.
 - Momentum of emitted particles should be conserved.
- $$\left| \sum_a^N e^{-i \sum_i k_i r} \mathcal{M}_a \right|^2$$

Demonstration on the Macro-coherent Amplification

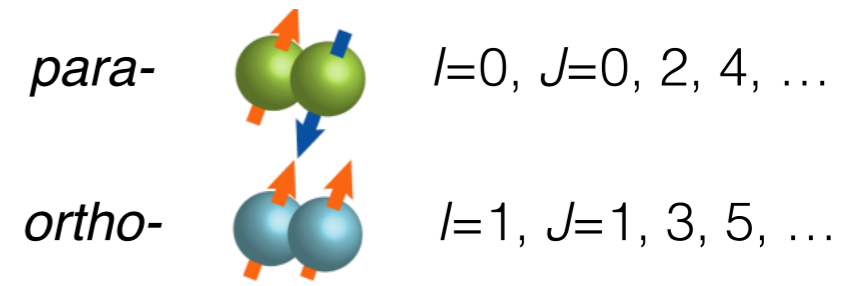
- Macro-coherent Amplification should be demonstrated experimentally.
→ Two photon emission version, named Paired Super-radiance (*PSR*), is a suitable example for this phenomenon.
- Rare QED process. Macro-coherent amplification is needed to observe it.
- Easier than RENP.



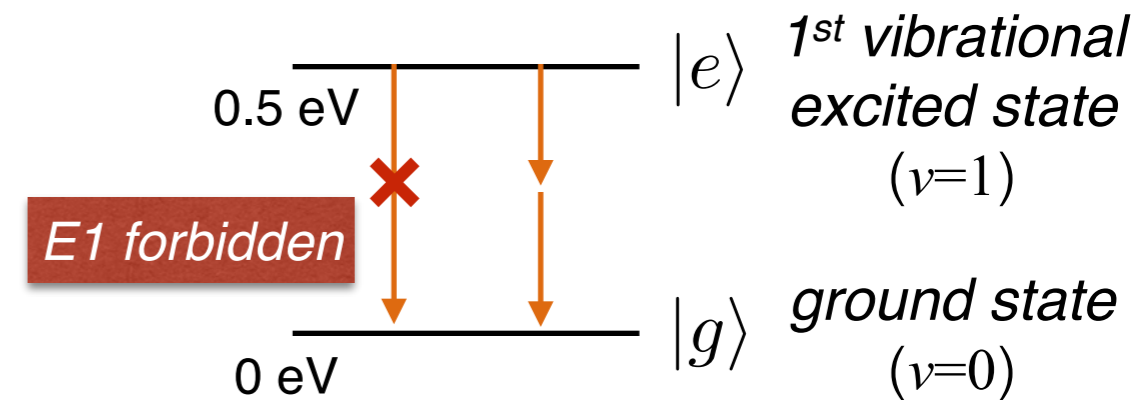
PSR experiment using $p\text{H}_2$ molecule

- Target : para- H_2 gas
@ 77 K, 60 kPa ($5.6 \times 10^{19} \text{ cm}^{-3}$)

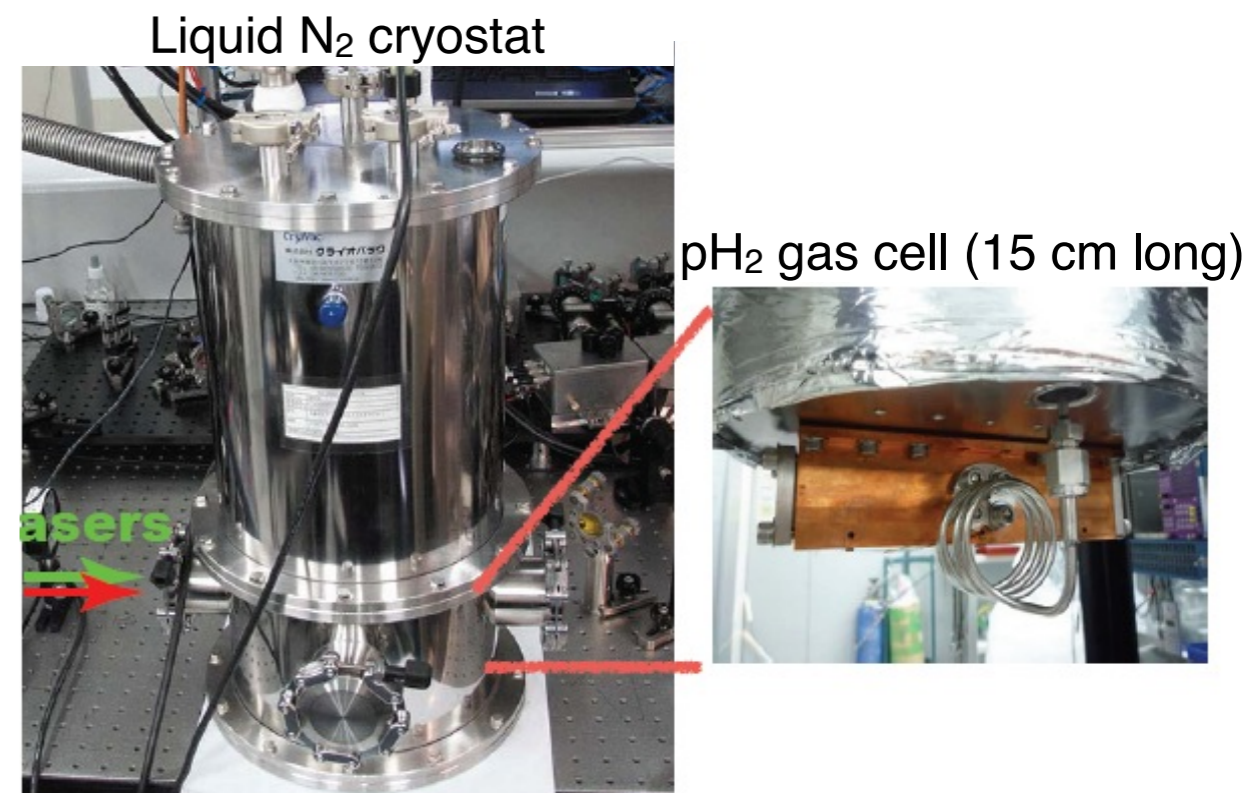
para- & ortho- hydrogen



- Two photon emission can be expected
 - One photon (E1) transition is forbidden
 - Two photon (E1×E1) transition is allowed
* Rate $\sim 5 \times 10^{-12} \text{ Hz}$

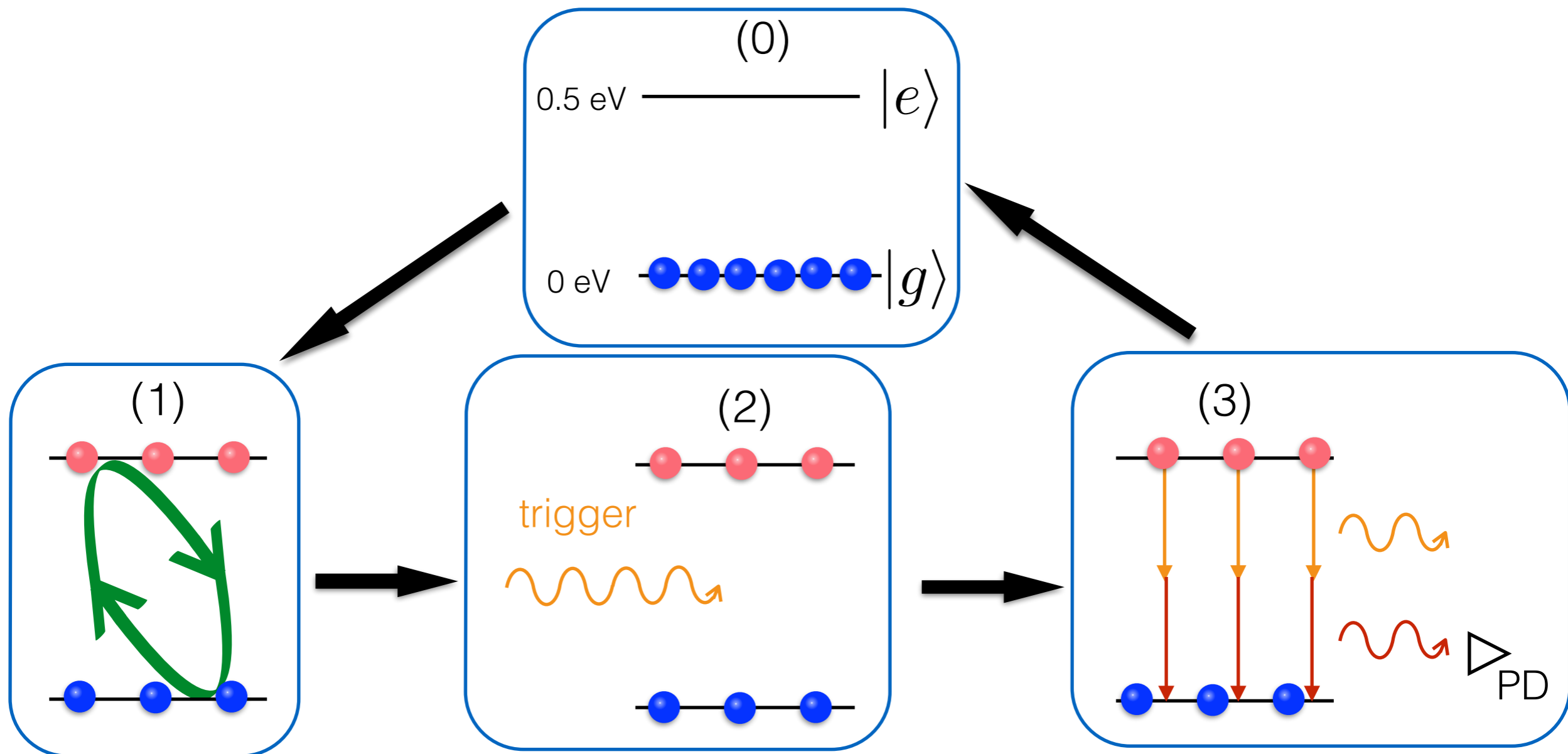


- Initial coherence can be generated
 - Previous researches have studied coherence generation methods.
→ Adiabatic Raman Excitation



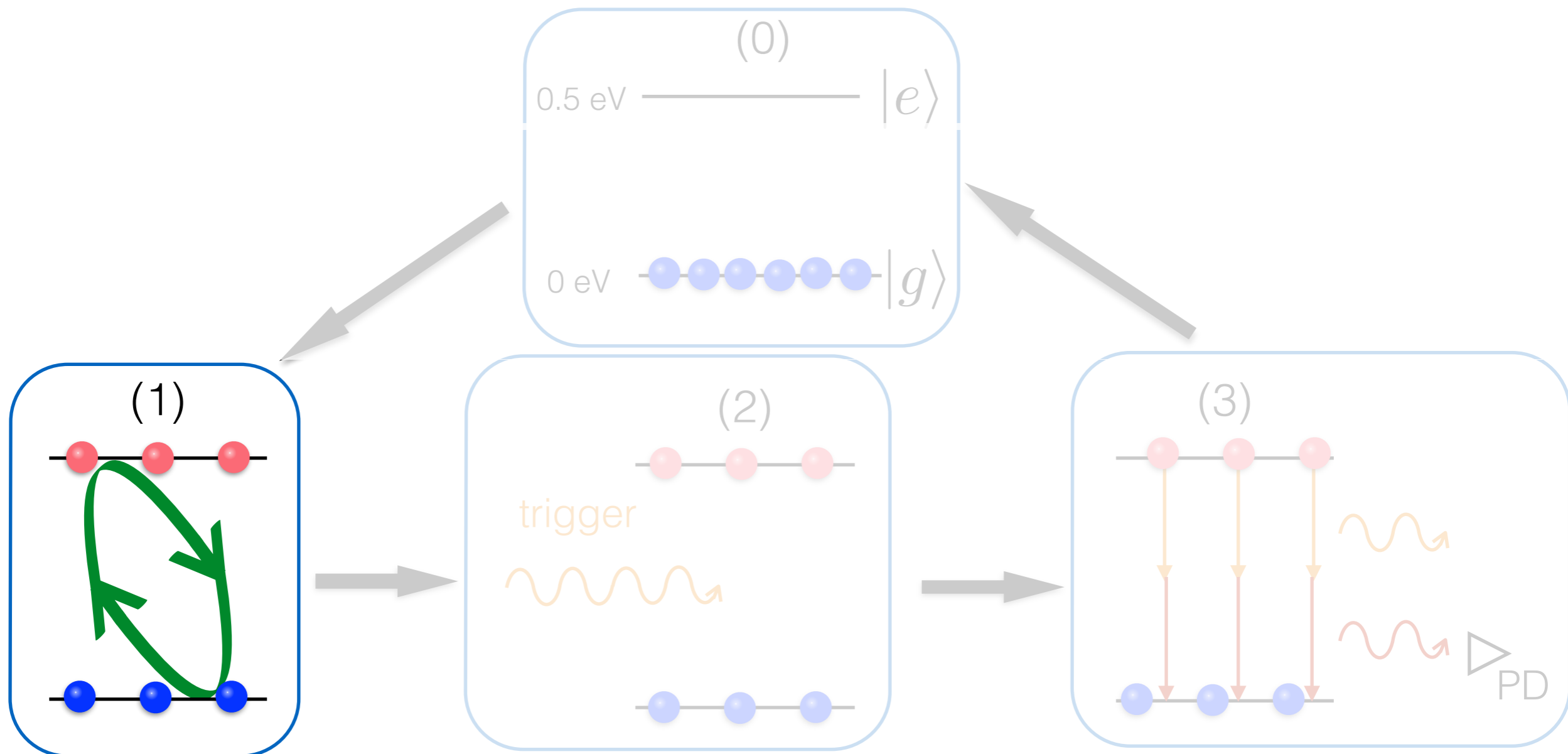
Experiment procedure

- (1) Generate initial coherence in $p\text{H}_2$ medium using **Adiabatic Raman Excitation**
- (2) Stimulate two photon emission by injecting **trigger laser**
- (3) Measure the **intensity of the other wavelength** from two photon de-excitation
- (4) Repeat (1)~(3) by 10 Hz



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(1) Coherence preparation

State w/ Laser Fields : $\psi(t) = \cos \frac{\theta(t)}{2} |g\rangle + \sin \frac{\theta(t)}{2} e^{i\phi(t)} |e\rangle$

Mixing angle : $\tan \theta(t) \simeq \frac{2|\Omega_{ge}(t)|}{\delta}$

Coherence : $\rho_{ge}(t) = [|\psi(t)\rangle\langle\psi(t)|]_{ge} = \frac{1}{2} \sin \theta(t)$

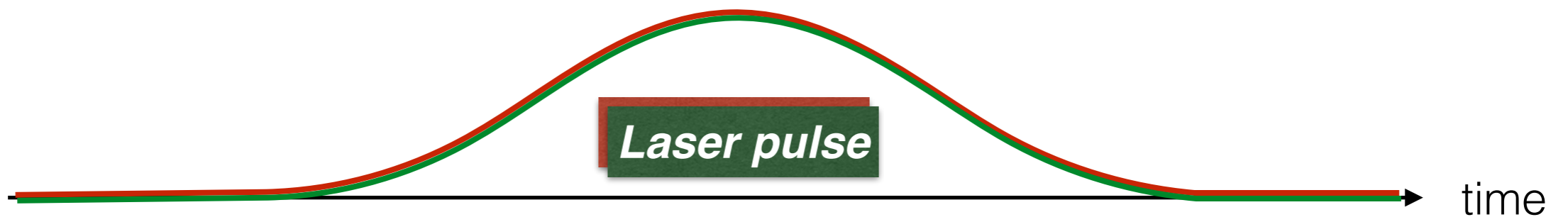
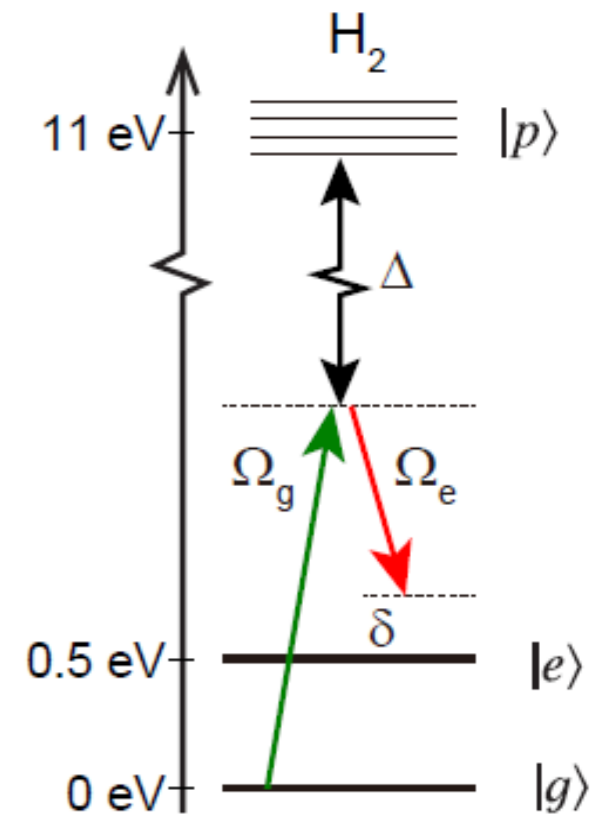
$$\Omega_{ge}(t) \equiv \frac{\Omega_g(t)\Omega_e(t)}{\Delta}$$

$$= \frac{(\vec{d}_{gp} \cdot \vec{E}_g(t))(\vec{d}_{pe} \cdot \vec{E}_e(t))}{\Delta}$$

$$e^{i\phi(t)} \equiv \frac{\Omega_{ge}(t)}{|\Omega_{ge}(t)|}$$

Theory: S.E. Harris & A.V. Sokolov, Phys. Rev. A 55, R4019 (1997)

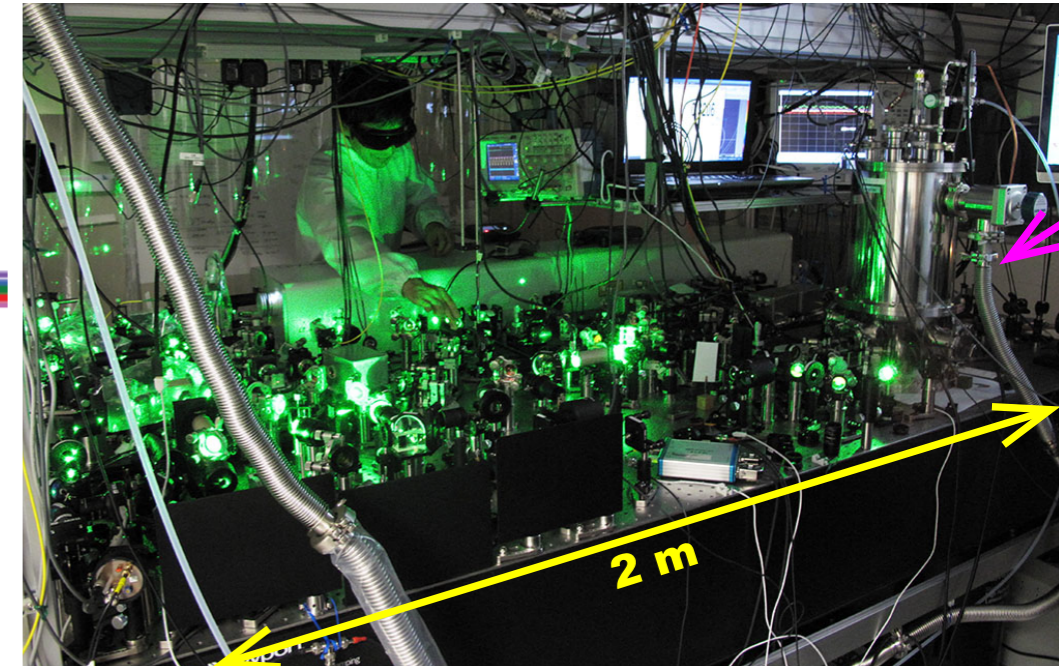
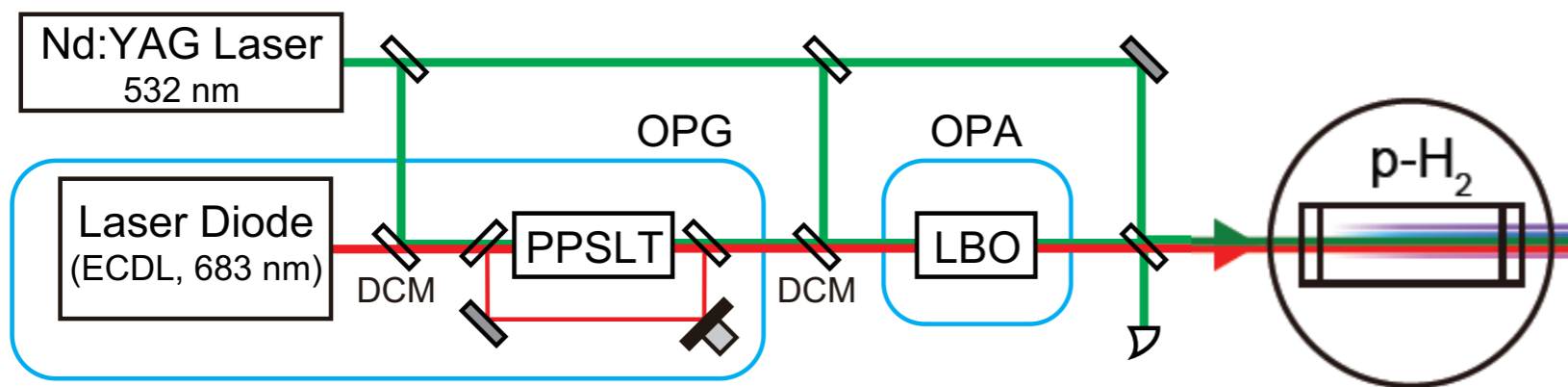
Experiments: M. Katsuragawa et al., CLEO/QELS (1999), A.V. Sokolov et al., PRL 85, 562 (2000)



state	$ g\rangle$ ($\theta = 0$)	$\frac{1}{\sqrt{2}} g\rangle + \frac{1}{\sqrt{2}} e\rangle$ ($\theta = \frac{\pi}{2}$)	$ g\rangle$ ($\theta = 0$)
coherence	0	0.5	0

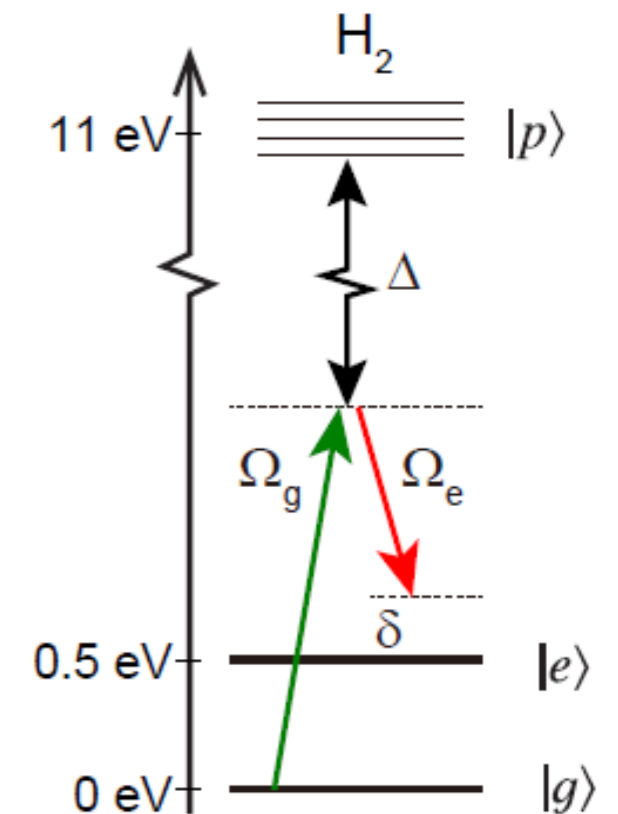
Strong two color lasers generate the initial coherence in the pH_2 medium

Driving laser system

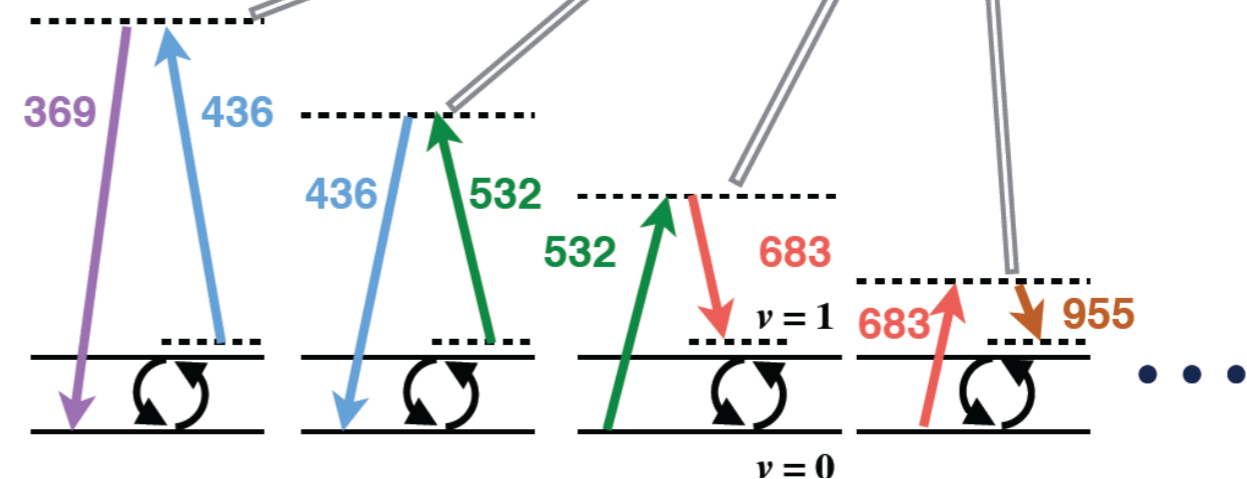
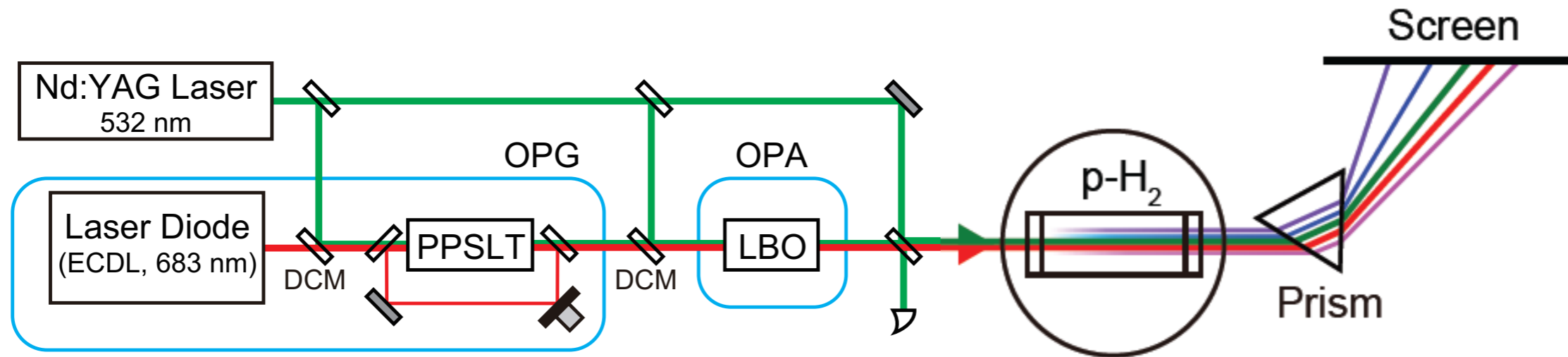


- Two driving Lasers

Light source	Nd:YAG	ECDL + Non-linear crystals
Wavelength	532 nm (green) ~ 2.3 eV	683 nm (red) ~ 1.8 eV
Pulse duration	8 ns	6 ns
Pulse energy*	4.3 mJ	4.3 mJ



Initial Coherence Generation

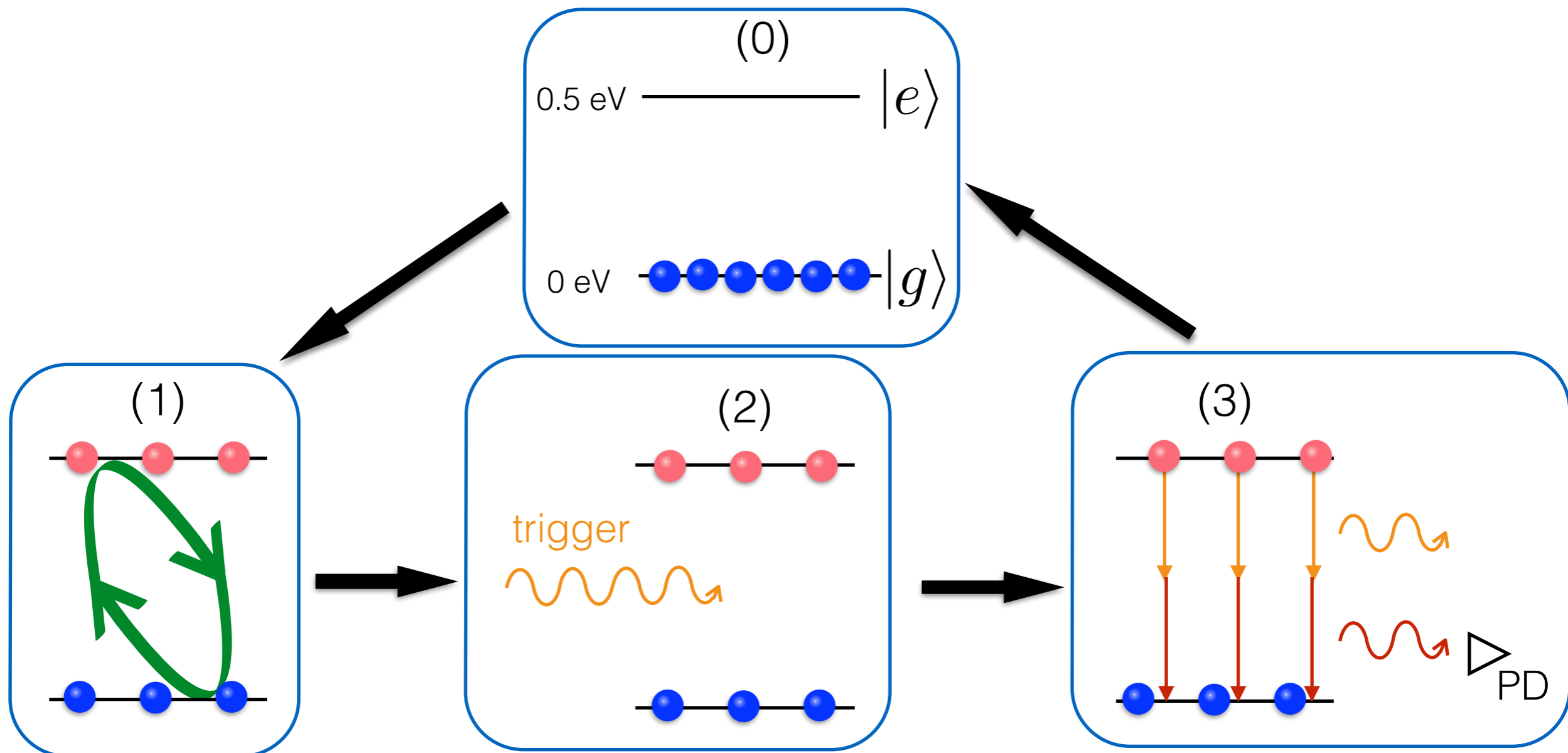


- The Raman side band in higher orders were observed.
→ Initial coherence was generated.

($\rho_{eg}=0.032$ based on a numerical simulation)

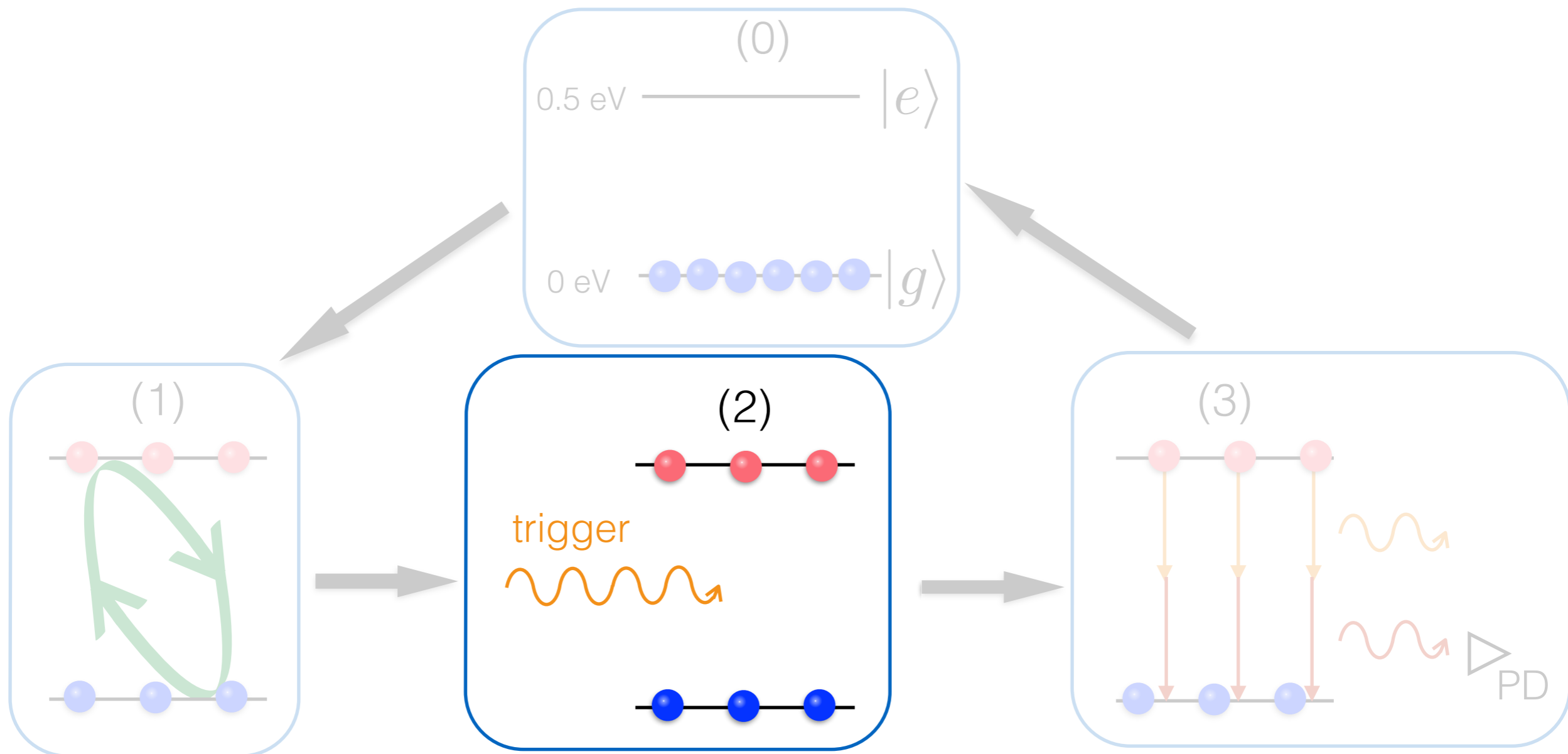
Experiment procedure

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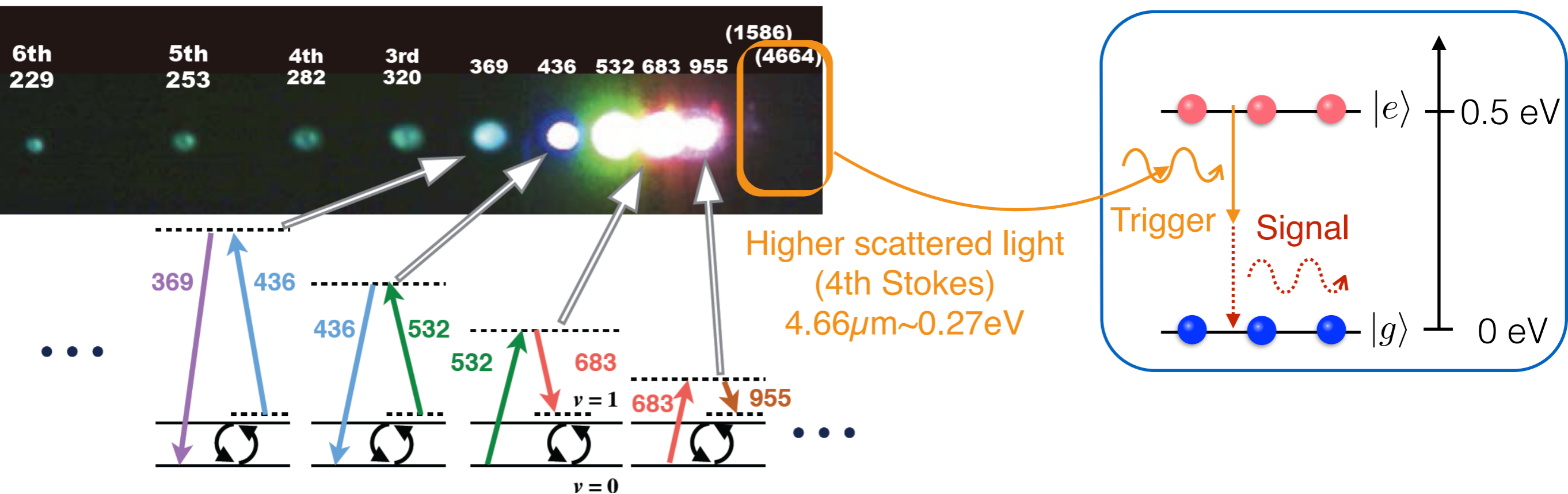


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(2) Trigger

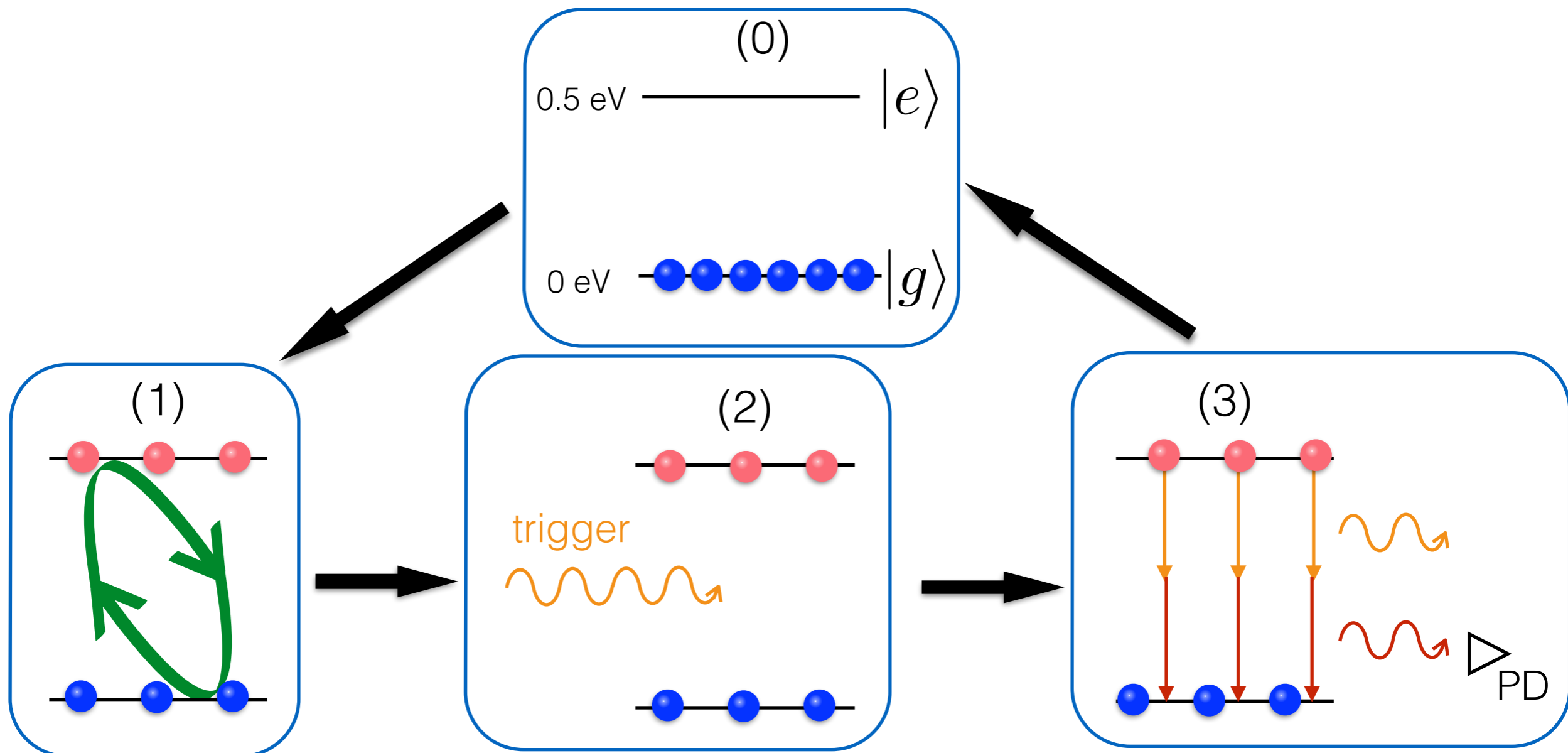


- The energy of 4th Stokes light (4.66 μm = 0.27 eV) is almost a half of the transition energy (0.5 eV)
 → 4th Stokes can be a trigger.

Trigger	4.66 μm (Mid Infra Red)	0.27 eV
Signal	4.96 μm (Mid Infra Red)	0.25 eV

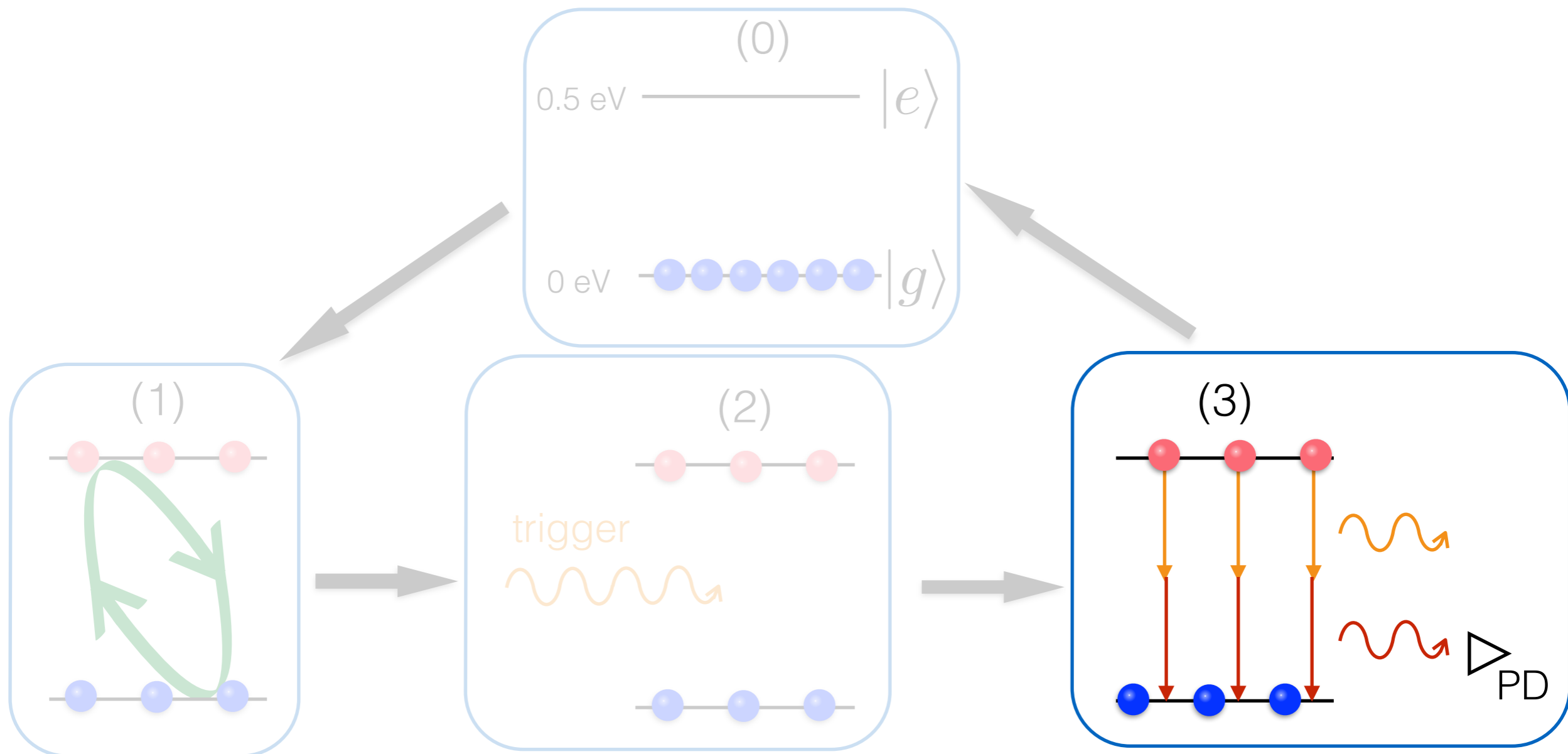
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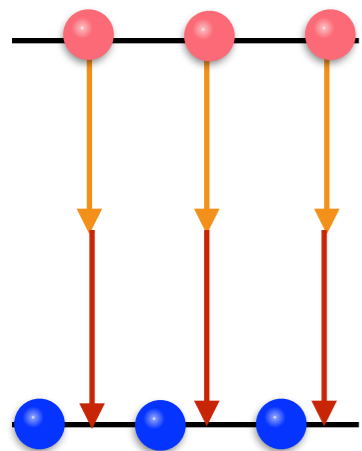
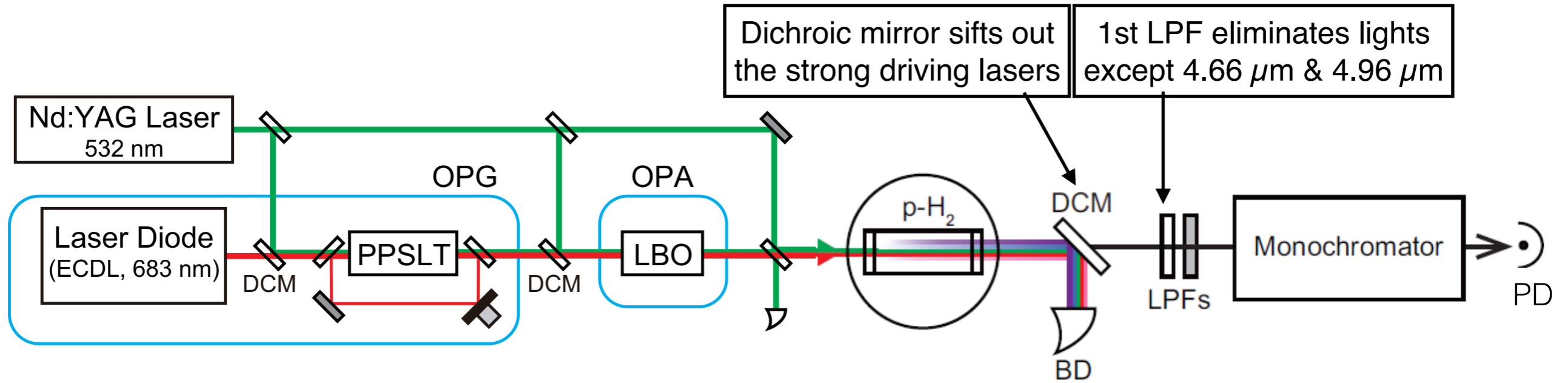


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(3) Observation of Emission



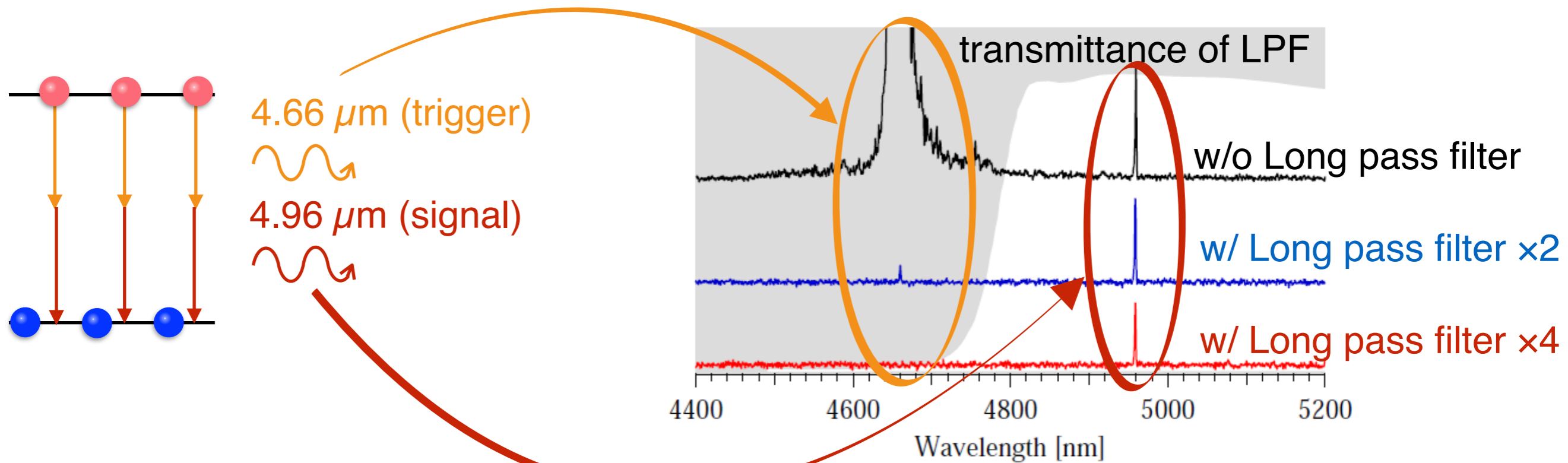
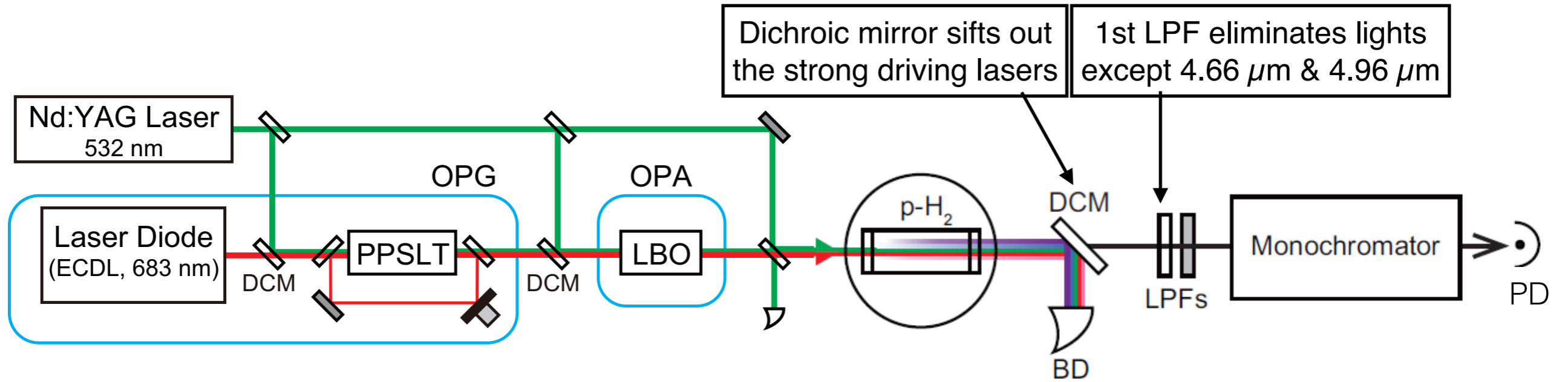
4.66 μm (trigger)



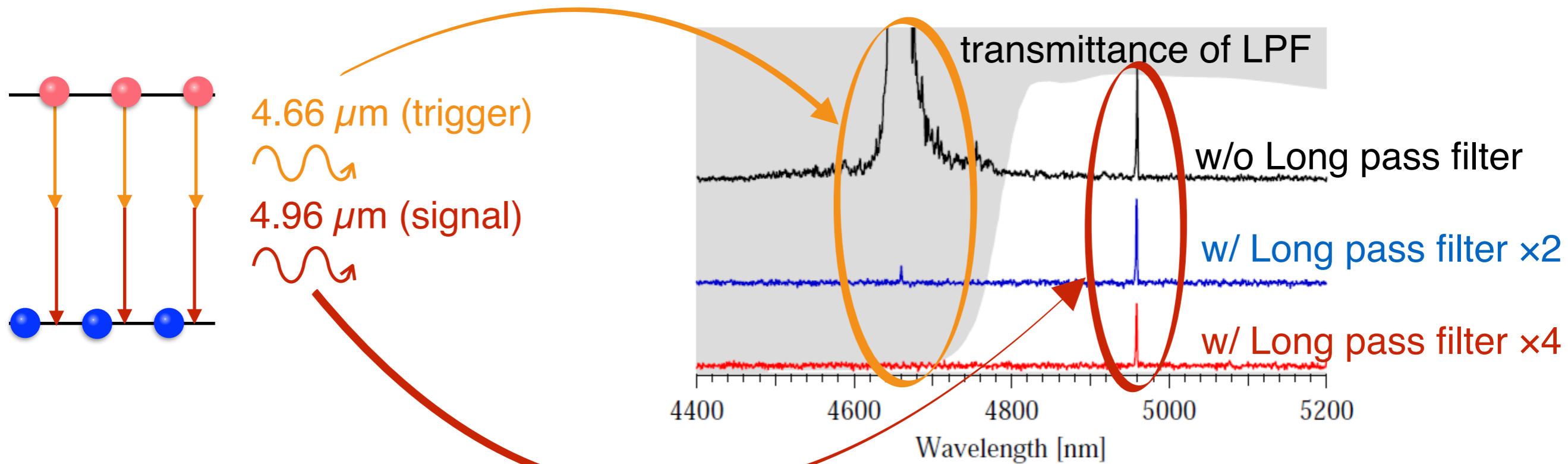
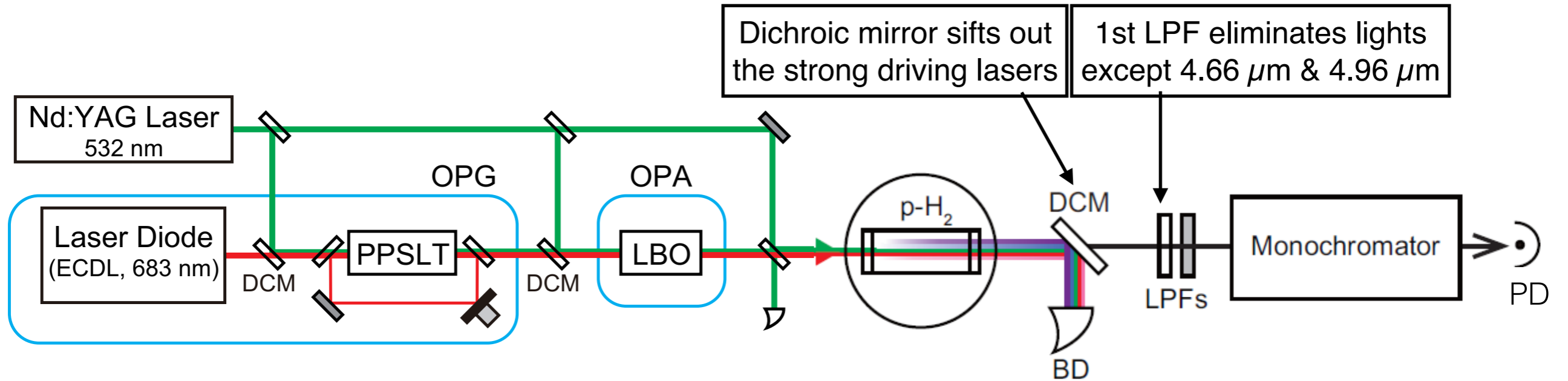
4.96 μm (signal)



(3) Observation of Emission

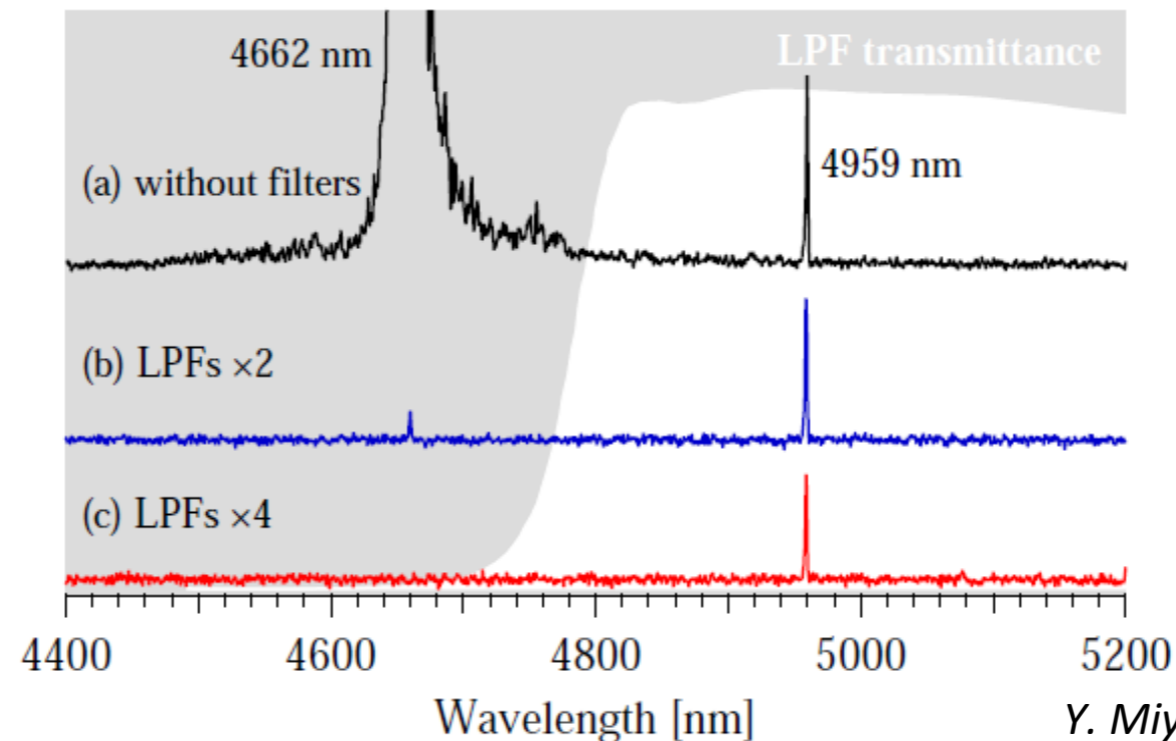


(3) Observation of Emission



Clear peak exists!!

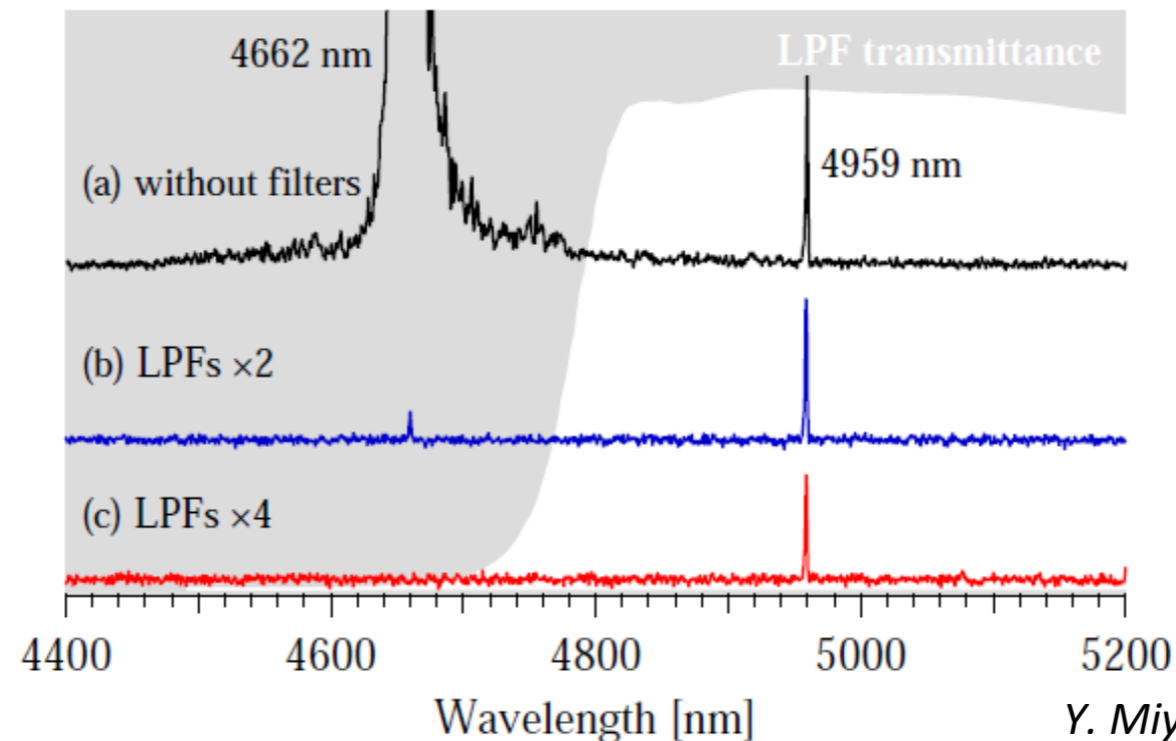
Enhancement Factor



- Observed number of photons : $> 4.4 \times 10^7$ photon/pulse
- Expected number of photons : $< 1.6 \times 10^{-8}$ photon/pulse

Factor	Value
Spontaneous decay rate : dA/dz ($z=\omega/\omega$)	$3.2 \times 10^{-11} \text{ s}^{-1}$ ($z=1/2$)
Energy band width (monochromator) : Δz	4.9×10^{-3}
measurement time : Δt	80 ns
detector solid angle fraction : $\Delta\Omega/(4\pi)$	1.2×10^{-4}
the maximum number of excited molecules	$< 1.5 \times 10^{16}$
The expected number of photons	$< 1.6 \times 10^{-8}$ photon/pulse

Enhancement Factor



Y. Miyamoto et. al. , arXiv:1406.2198v2

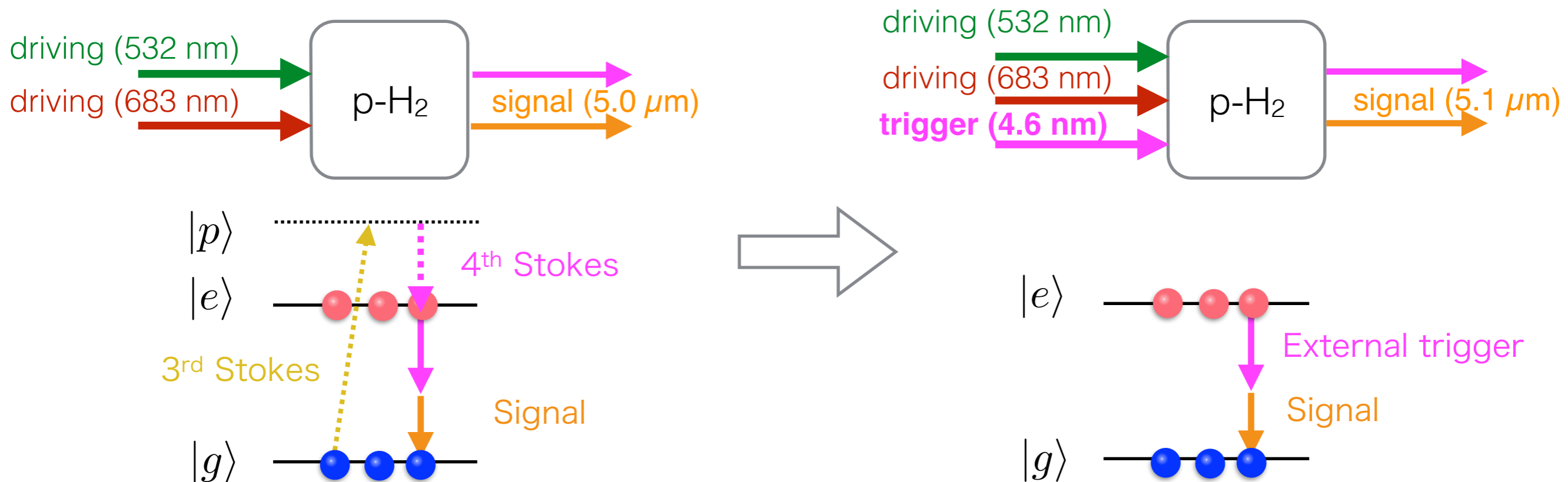
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The expected number of photons	$< 1.6 \times 10^{-8}$ photon/pulse

Huge enhancement ($> 10^{15}$) has been obtained!

Next experiments

- PSR with External trigger



Independent control

coherence generation : driving power, p-H₂ pressure, etc.

PSR stimulation : power, timing, etc.

⇒ **Quantitative discussion**

- *RENP has potential to determine the neutrino absolute mass.*
 - Macro-coherent Amplification is essential to realize it.
- *Macro-coherent Amplification has been demonstrated using two photon emission*

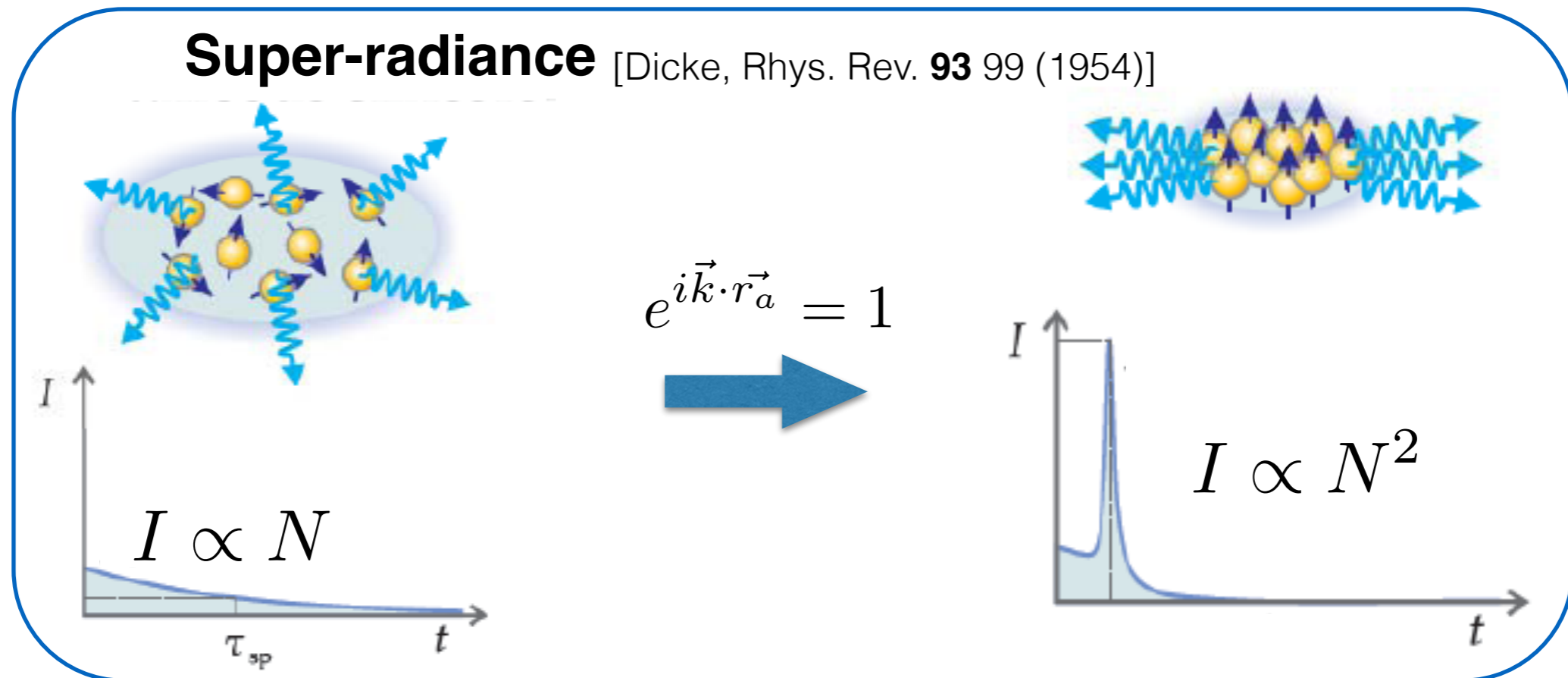
***4.96 μm emission has been observed.
Two photon emission has been enhanced by 10^{15}***

- *The next experiment is on-going*
 - PSR experiment with External trigger
 - Separate the coherence generation and the PSR stimulation.
 - Enable us to discuss more quantitatively.

Back up

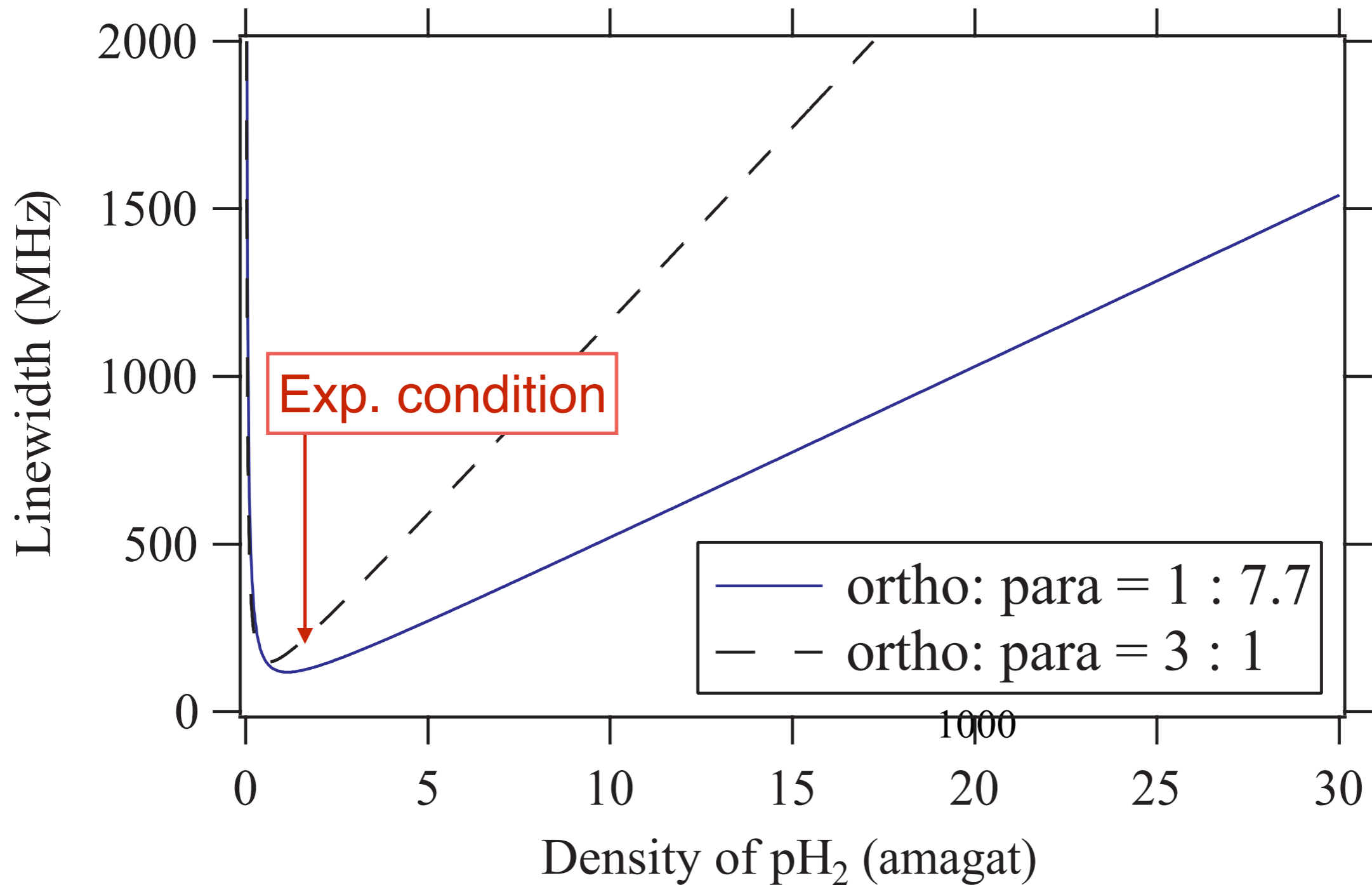
Enhancement by Coherence

- Example of the amplification by coherence : **Single particle emission**



- The maximal coherent region is *limited by the wavelength* ($\sim 1 \mu\text{m}$)
 $\rightarrow N < 10^8$ (for 10^{20} cm^{-3})
- Macro-Coherent Amplification : **Plural particles emission**
 - Thanks to the momentum conservation ($e^{i(k_1+k_2+k_3)r_i} = 1$),
the wavelength limitation can be removed.
 \rightarrow *N can reach the Avogadro number.*

pH₂ line width



$$1 \text{ amagat} \sim 2.7 \times 10^{19} \text{ cm}^{-3}$$

Adiabatic Raman Excitation

- State of atoms

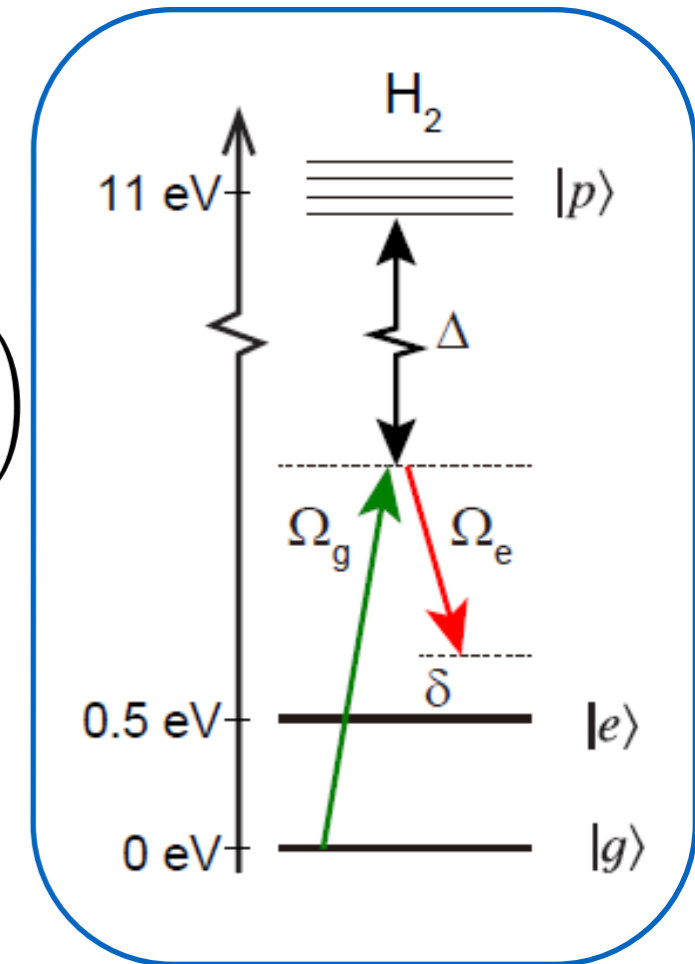
- Wave function $\psi = c_g|g\rangle + c_e|e\rangle$

- Density matrix : $\rho = \begin{pmatrix} \rho_{gg} & \rho_{ge}e^{-i\phi} \\ \rho_{eg}e^{i\phi} & \rho_{ee} \end{pmatrix} = \begin{pmatrix} |c_g|^2 & c_g c_e^* \\ c_e c_g^* & |c_e|^2 \end{pmatrix}$

- Inject two color lasers(E_g, E_e) simultaneously

- Effective hamiltonian

$$i\hbar \begin{pmatrix} \frac{\partial c_g}{\partial t} \\ \frac{\partial c_e}{\partial t} \end{pmatrix} = \mathcal{H}_{eff} \begin{pmatrix} c_g \\ c_e \end{pmatrix} = -\hbar \begin{pmatrix} \Omega_{gg} & \Omega_{ge} \\ \Omega_{eg} & \Omega_{ee} - \delta \end{pmatrix} \begin{pmatrix} c_g \\ c_e \end{pmatrix}$$



- Two color lasers change the eigenstates of the system from $|g\rangle$ and $|e\rangle$ to the superposition of them.

Eigenstates : $|\pm\rangle = \cos \frac{\theta}{2} |g\rangle \pm \sin \frac{\theta}{2} e^{-i\phi} |e\rangle$

Mixing angle : $\tan \theta \simeq \frac{2|\Omega_{ge}|}{\delta}$

Coherence : $\rho_{ge} = \frac{1}{2} \sin \theta e^{i\phi}$

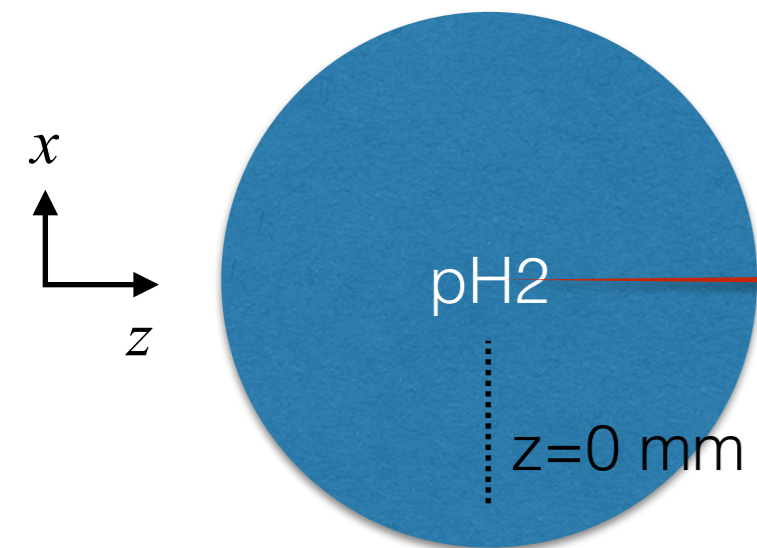
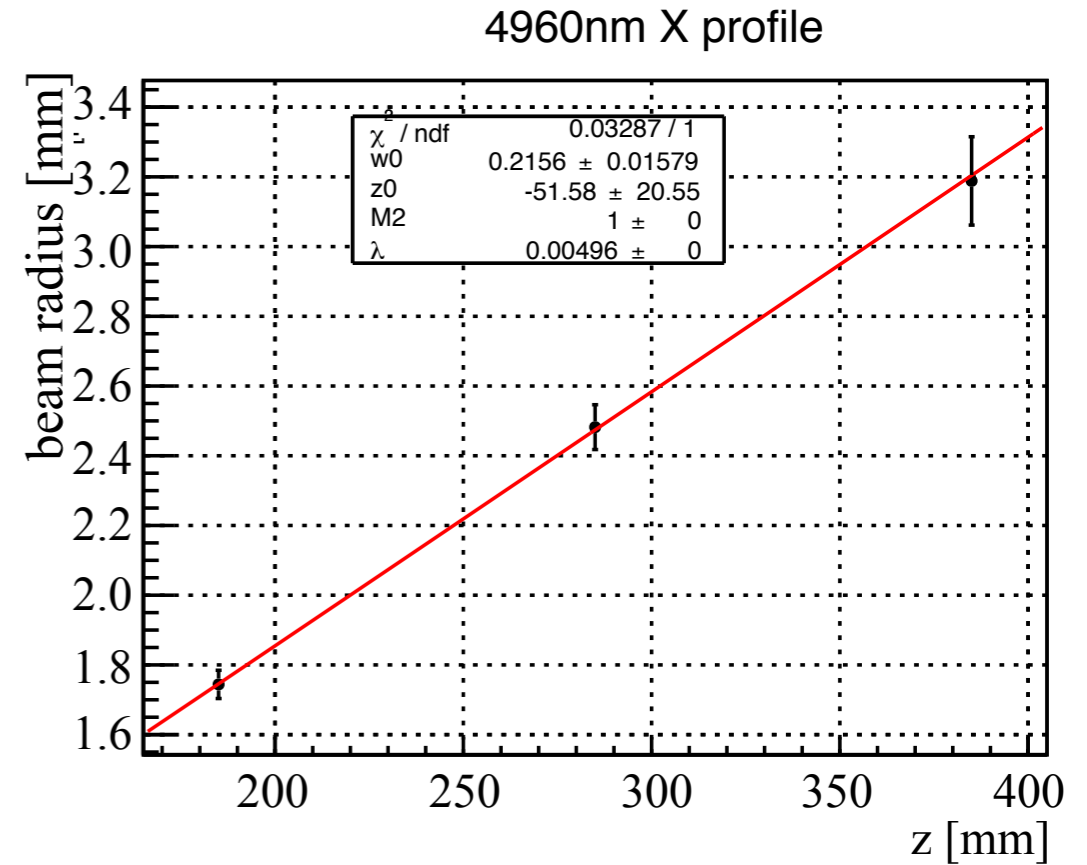
$$\Omega_{ij} \equiv \frac{\Omega_i \Omega_j}{\Delta} = \frac{(d_{ip} E_i)(d_{pj} E_j)}{\Delta}$$

$$\left(e^{i\phi} \equiv \frac{\Omega_{ge}}{|\Omega_{ge}|} \right)$$

4.96 μm spatial profile

- Knife edge method

	X [mrad]	Y [mrad]
4.96 μm	7.3	9.7
532 nm	1.8	2.0
683 nm	2.0	1.9



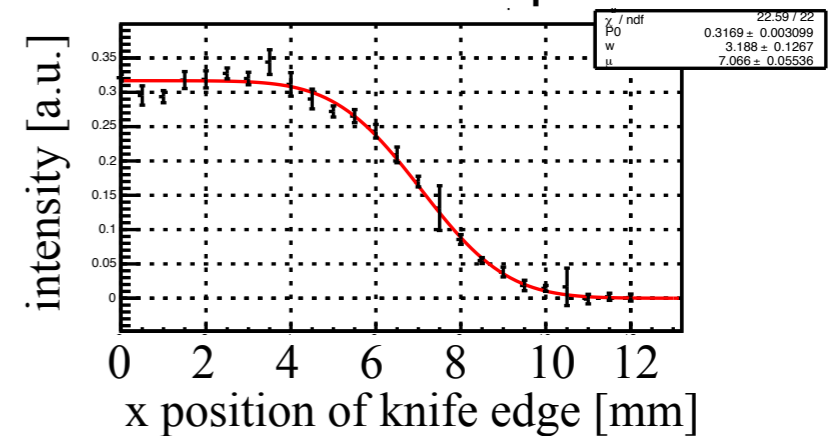
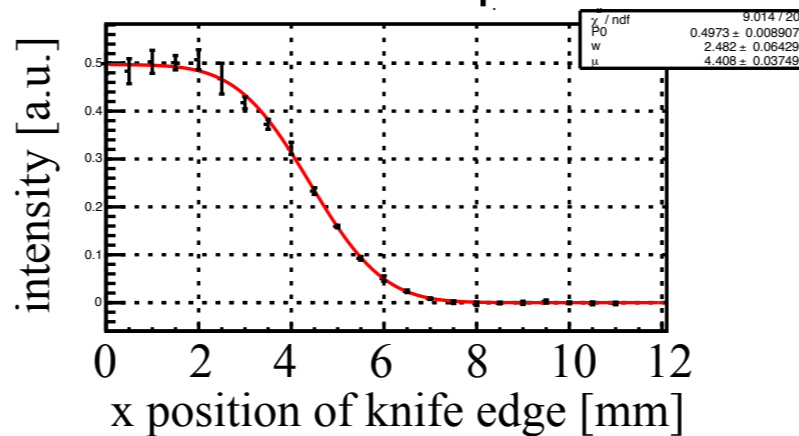
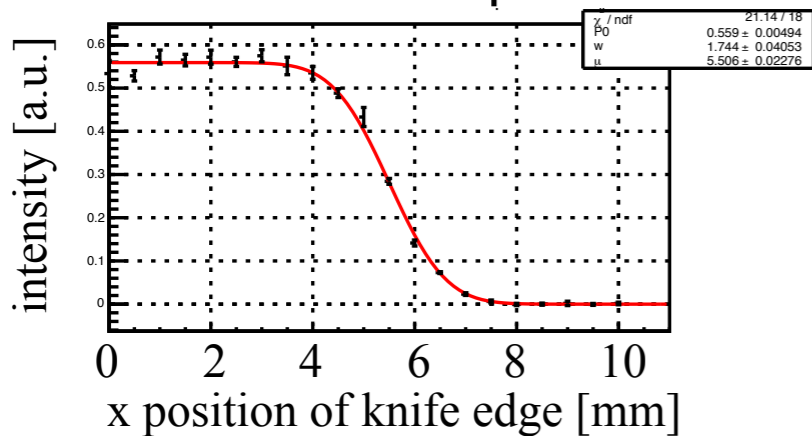
Monochromator & Photo Detector

$$P(a) = \frac{P_0}{2} \operatorname{Erfc} \left(\frac{\sqrt{2}(a - \mu)}{w} \right)$$

z=185 mm X profile

z=285 mm X profile

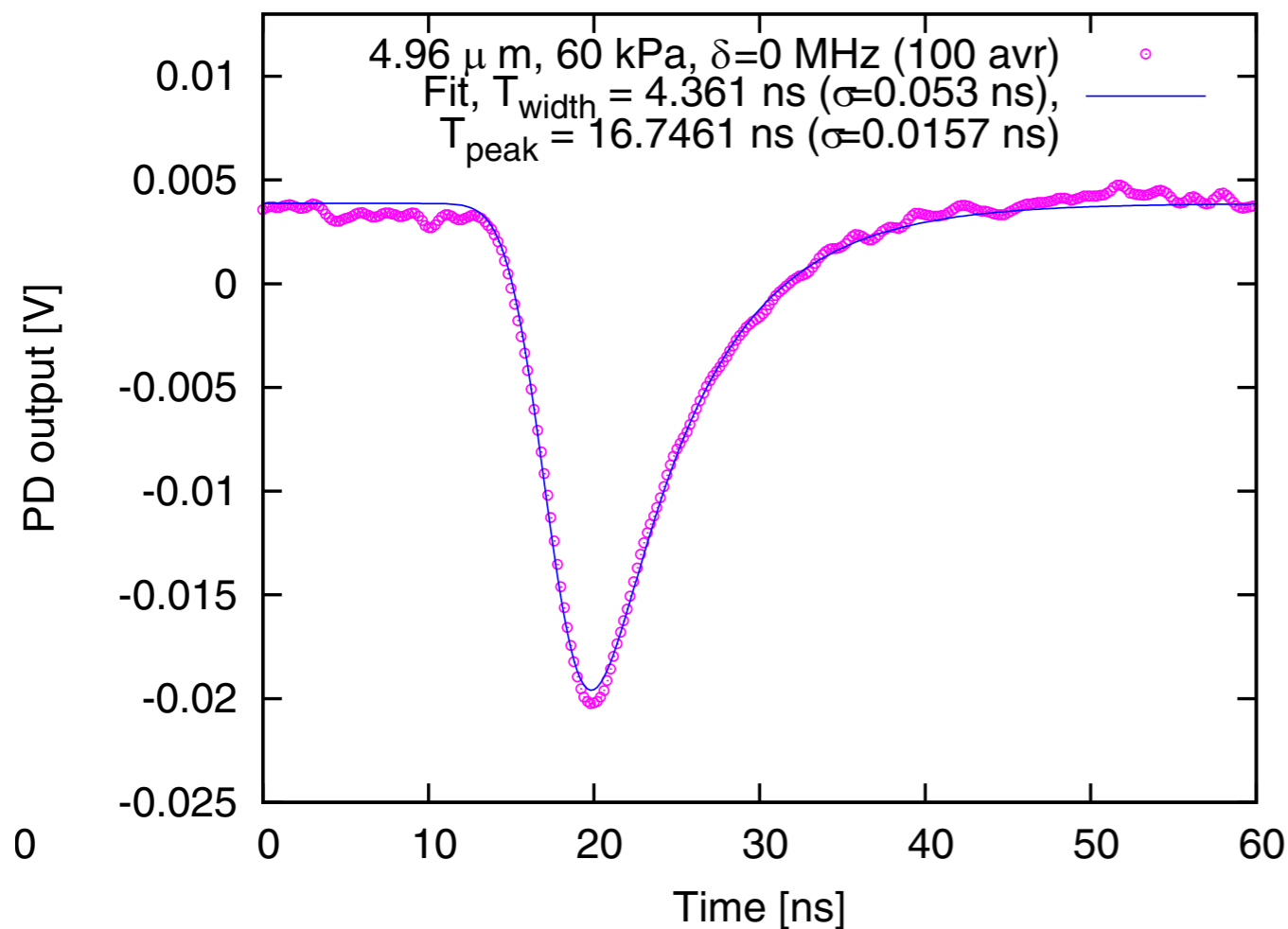
z=385 mm X profile



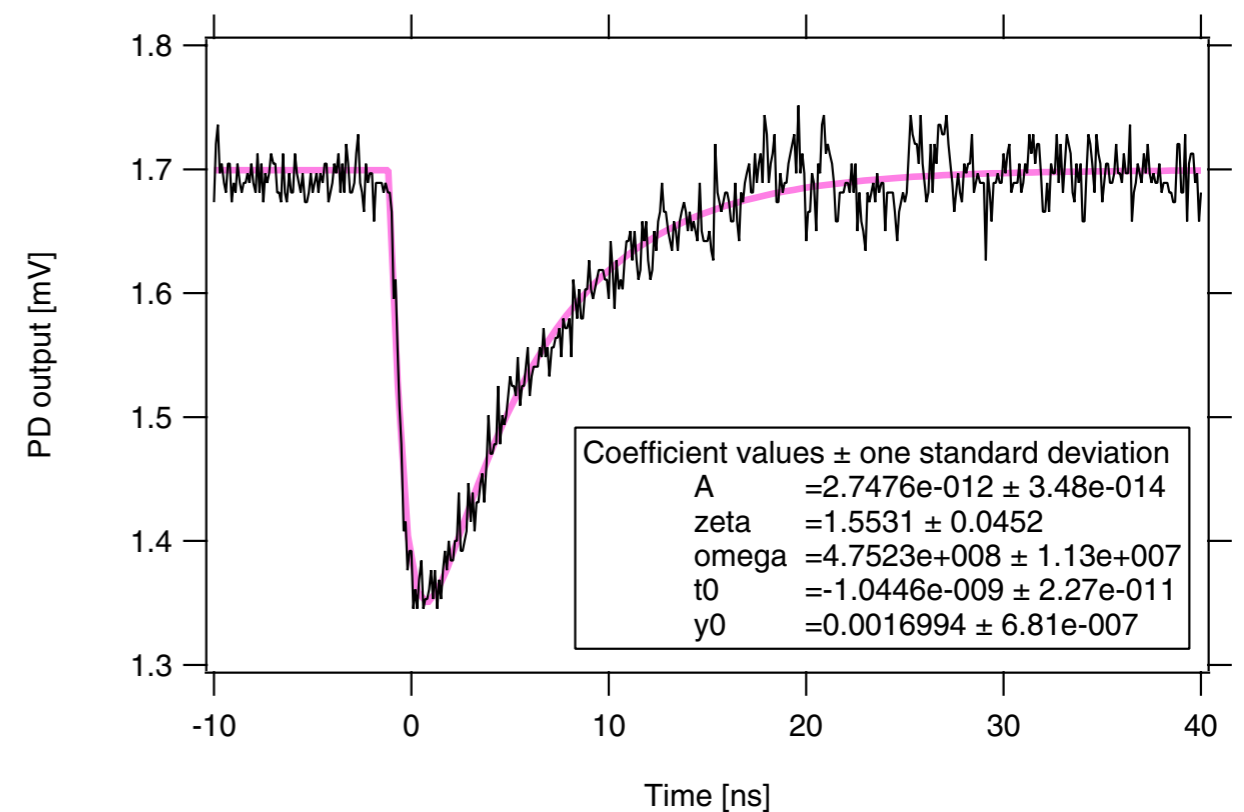
4.96 μm time profile

- Pulse duration of 4.36 ± 0.05 ns is slightly shorter than driving lasers.

4.96 μm pulse output

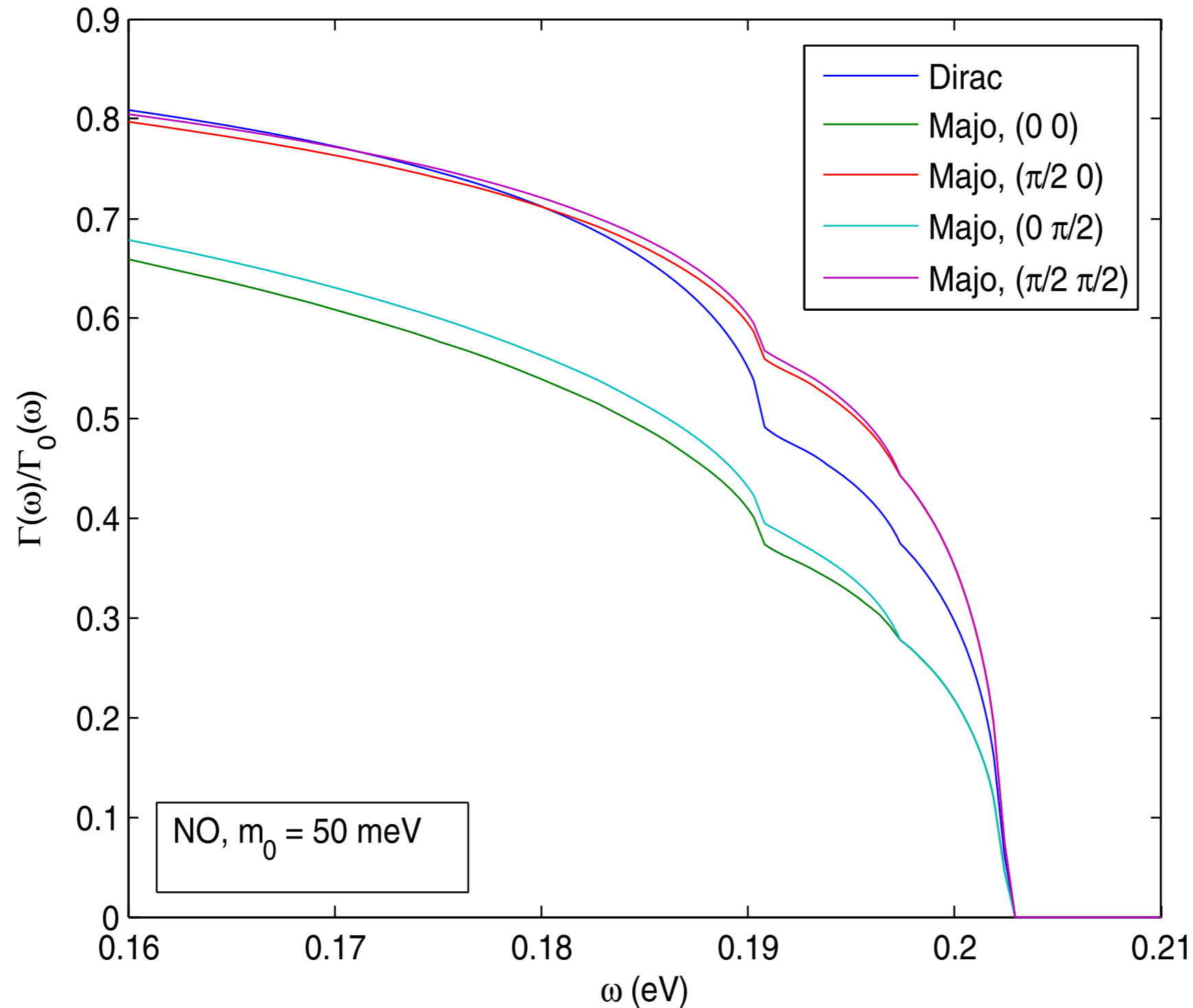


Detector response



$$g(t) = \frac{\omega_n}{\sqrt{\zeta^2 - 1}} e^{-\zeta\omega_n t} \sinh\left(\sqrt{\zeta^2 - 1}\omega_n t\right) \quad (\zeta > 1)$$

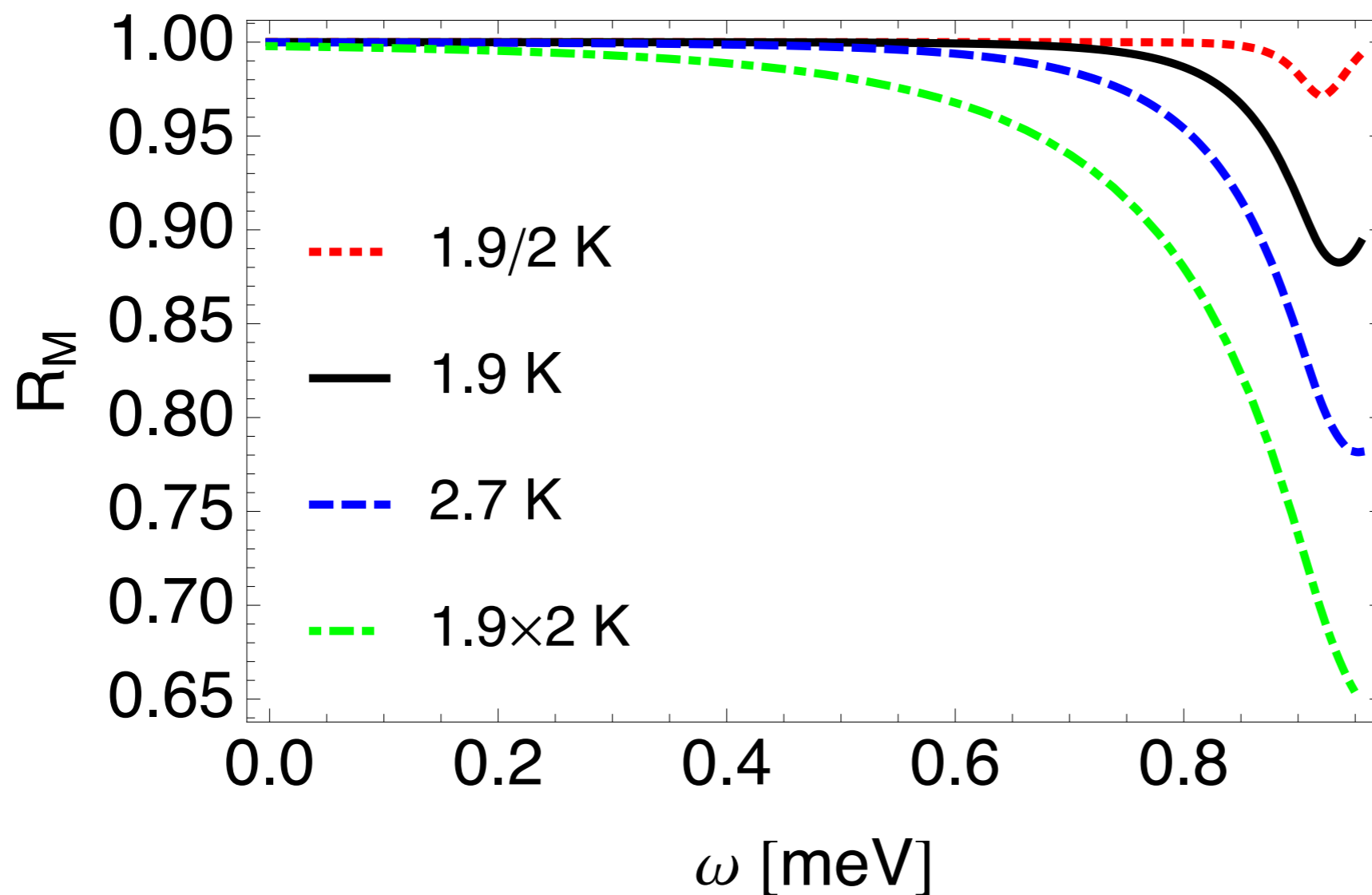
Majorana/Dirac, CPV phases



Normal hierarchy
 $m_1 = 50$ meV
 $\omega_{eg} = 0.43$ eV

D.N. Dinh et al., Phys. Lett. **B719** (2013)

Cosmic neutrino background



$m_{\text{smallest}} = 5 \text{ meV}$
 $\omega_{eg} = 11 \text{ meV}$
 $\mu = 0 \text{ eV}$

Time Development of the System

- Time-development of the system is obtained by numerically solving the Maxwell-Bloch equation

► Bloch equation

$$\frac{\partial \rho_{gg}}{\partial t} = i(\Omega_{ge}\rho_{eg} - \Omega_{eg}\rho_{ge}) + \gamma_1\rho_{ee}$$

$$\frac{\partial \rho_{ee}}{\partial t} = i(\Omega_{eg}\rho_{ge} - \Omega_{ge}\rho_{eg}) - \gamma_1\rho_{ee}$$

$$\frac{\partial \rho_{ge}}{\partial t} = i(\Omega_{gg} - \Omega_{ee} + \delta)\rho_{ge} + i\Omega_{ge}(\rho_{ee} - \rho_{gg}) - \gamma_2\rho_{ge}$$

Response of the molecules

► Maxwell equation

$$\frac{\partial \mathbf{E}_q}{\partial \xi} = \frac{i\omega_q \mathcal{N}}{2c} \left\{ (\rho_{gg}\alpha_{gg}^{(q)} + \rho_{ee}\alpha_{ee}^{(q)})\mathbf{E}_q + \rho_{eg}\alpha_{eg}^{(q-1)}\mathbf{E}_{q-1} + \rho_{ge}\alpha_{ge}^{(q)}\mathbf{E}_{q+1} \right\}$$

$$\frac{\partial \mathbf{E}_p}{\partial \xi} = \frac{i\omega_q \mathcal{N}}{2c} \left\{ (\rho_{gg}\alpha_{gg}^{(p)} + \rho_{ee}\alpha_{ee}^{(p)})\mathbf{E}_p + \rho_{eg}\alpha_{ge}^{(p\bar{p})}\mathbf{E}_{\bar{p}}^* \right\}$$

Wave propagation

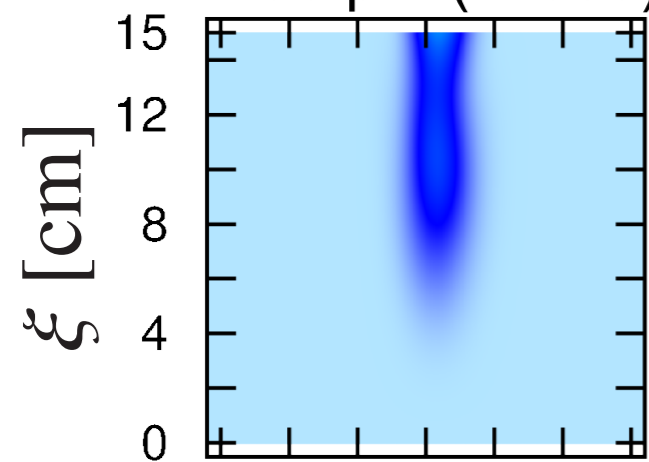
Stark Shift $\Omega_{aa} = \frac{\epsilon_0}{4\hbar} \sum_{m=q,p,\bar{p}} \alpha_{aa}^{(m)} |\mathbf{E}_m|^2 \quad (a = g, e)$

Two-photon Rabi frequency $\Omega_{ge} = \Omega_{eg}^* = \frac{\epsilon_0}{4\hbar} \left\{ \sum_q \alpha_{ge}^{(q)} \mathbf{E}_q \mathbf{E}_{q+1}^* + \alpha_{ge}^{(p\bar{p})} \mathbf{E}_p^* \mathbf{E}_{\bar{p}}^* \right\}$

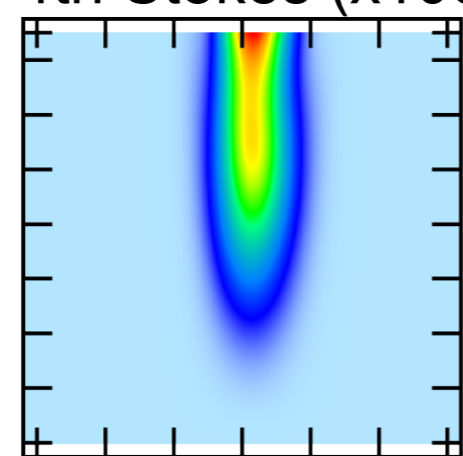
Simulation Results

電場 $|E_q|$

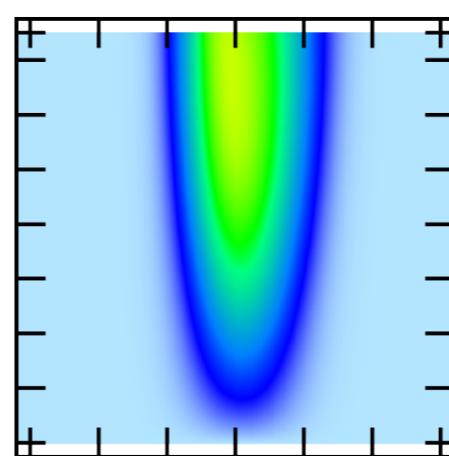
4.96 μm (x1000)



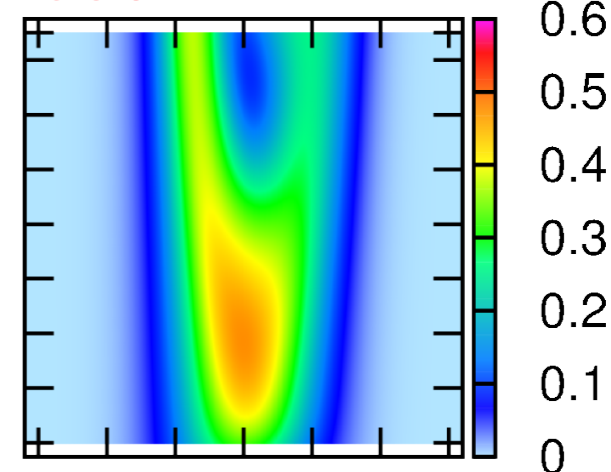
Trigger
4th Stokes (x100)



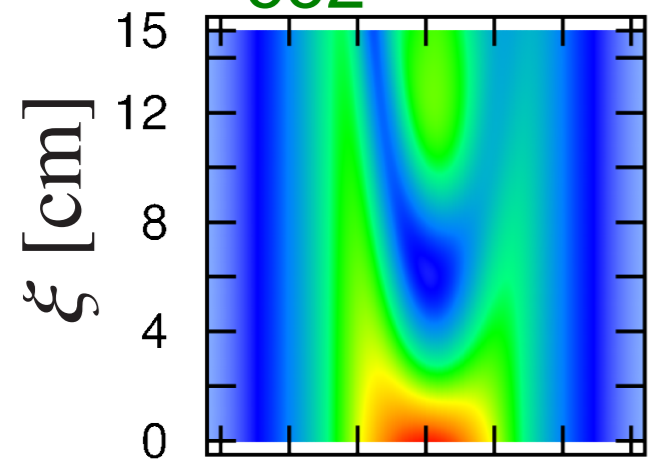
2nd Stokes



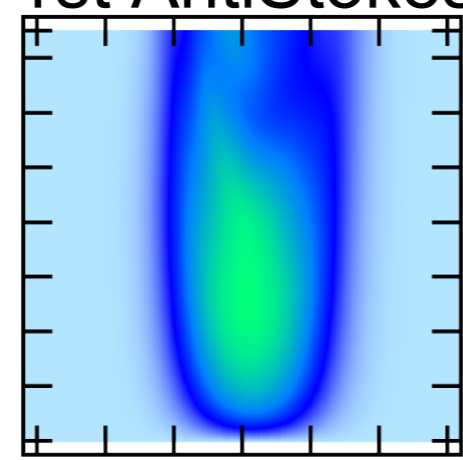
683



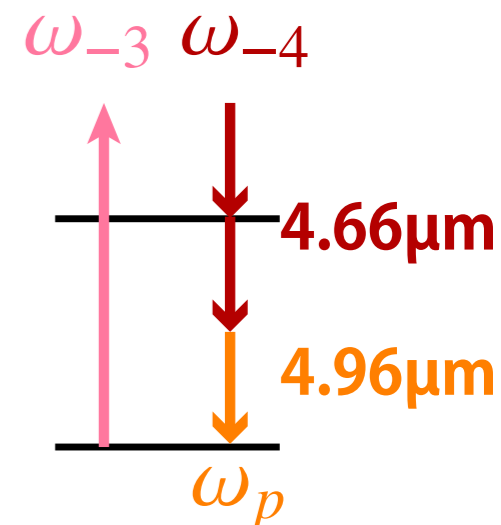
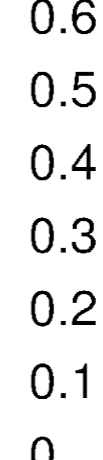
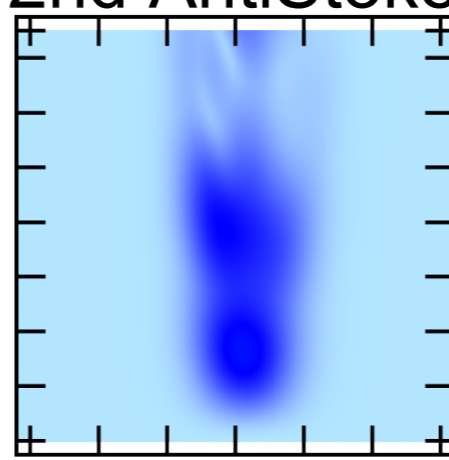
532



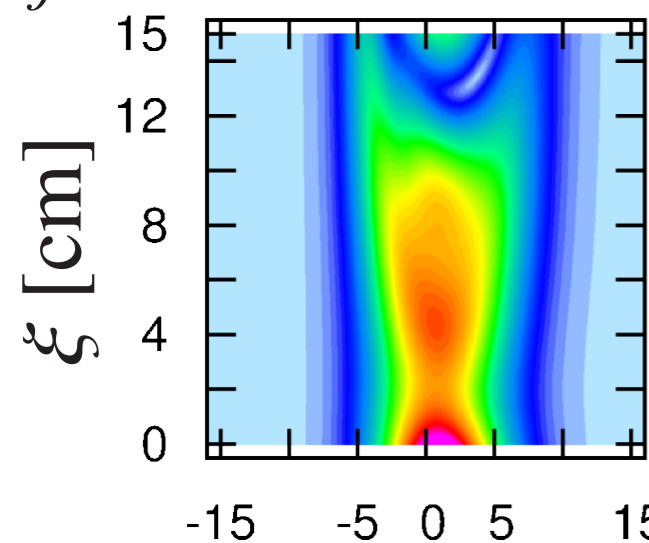
1st AntiStokes



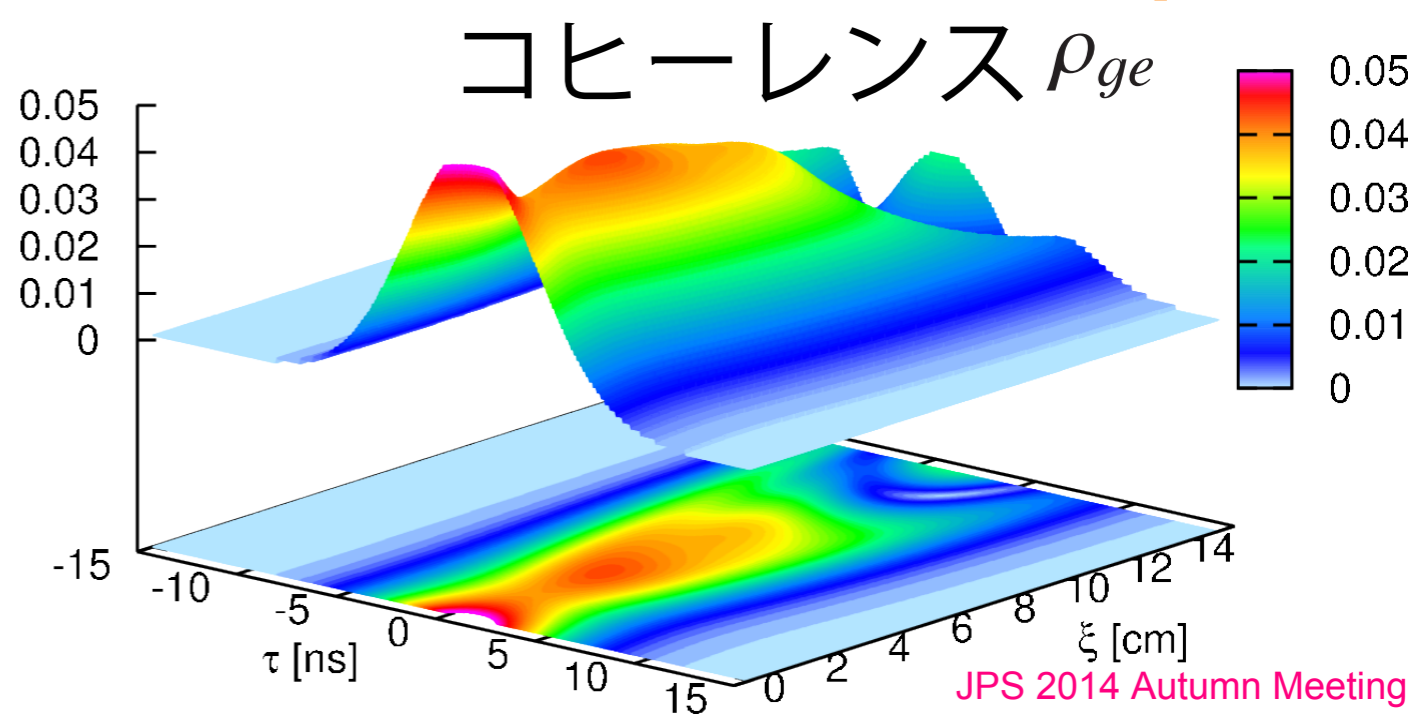
2nd AntiStokes



ρ_{ge}



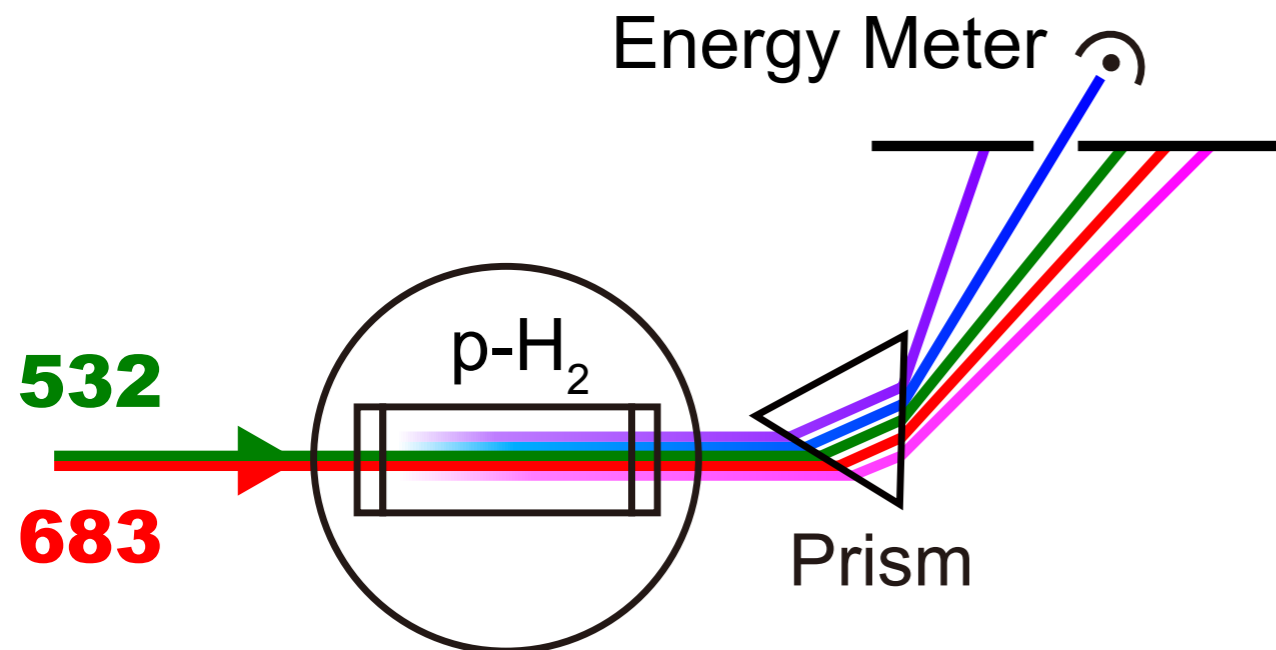
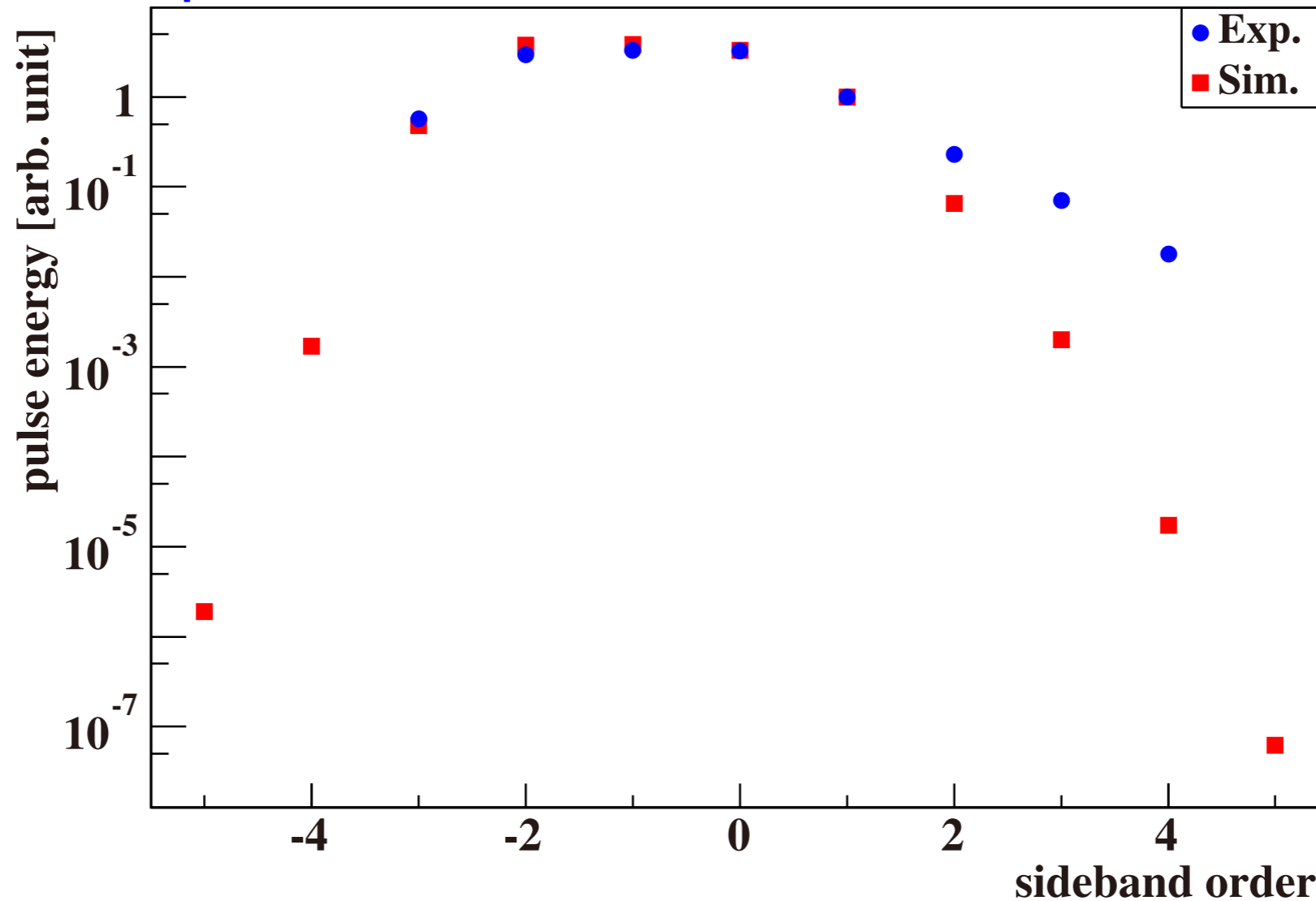
τ [ns]



コヒーレンス ρ_{ge}

Raman Coherence Estimation

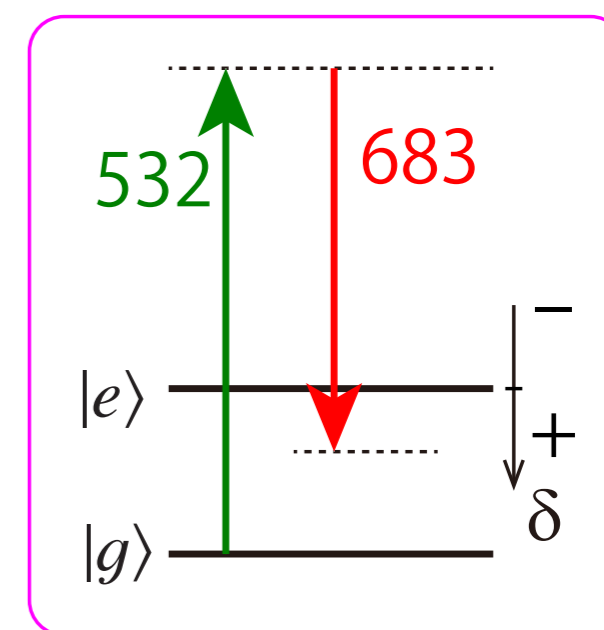
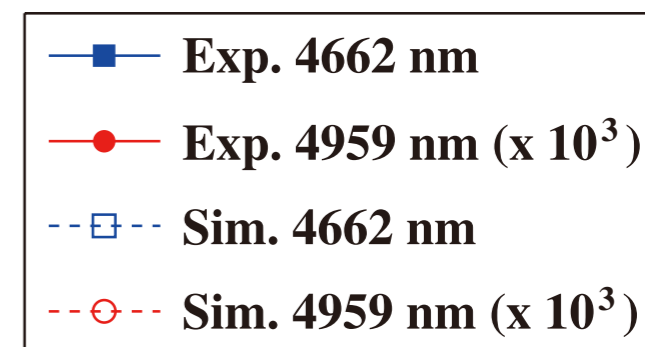
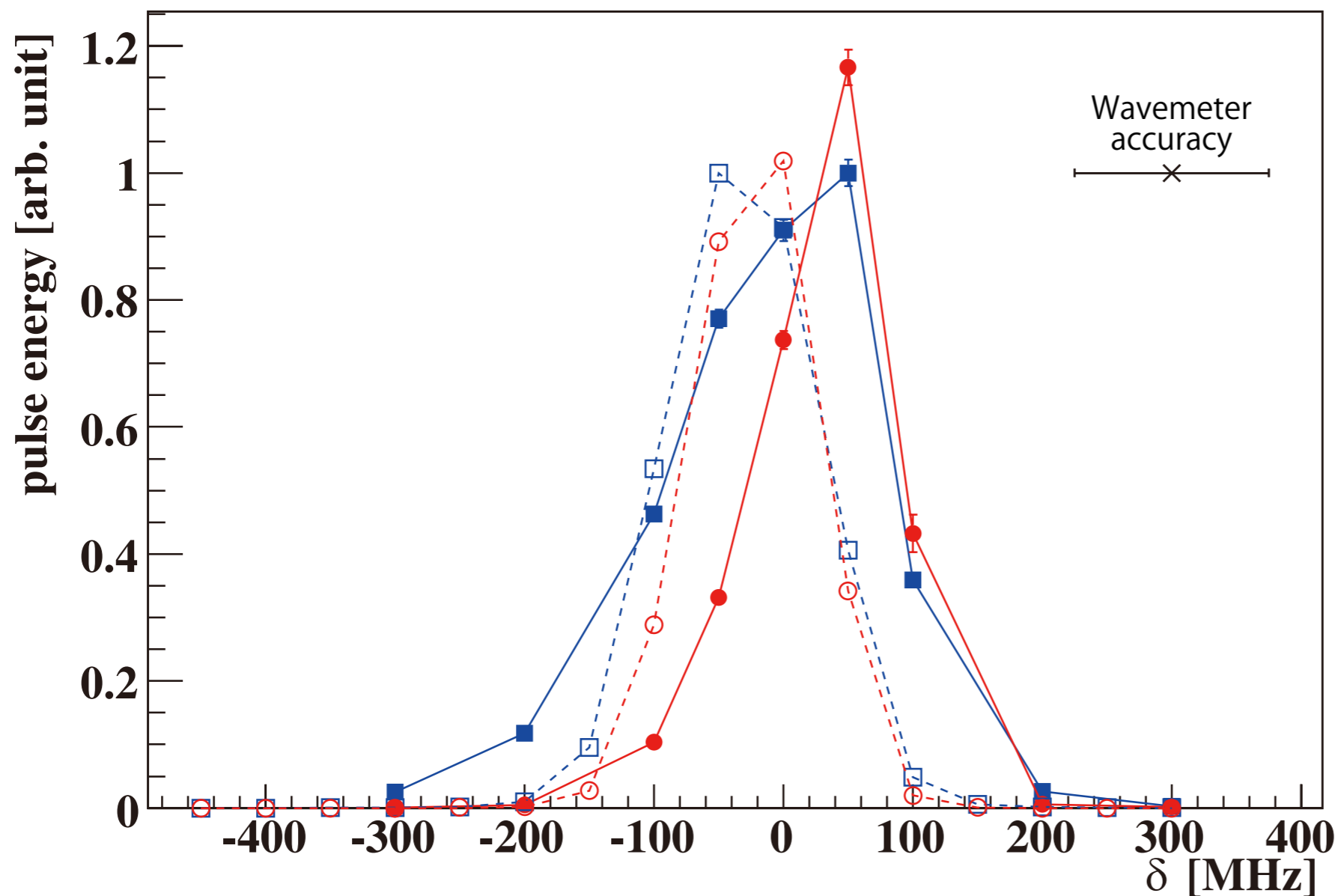
● Experimental v.s. Simulation result



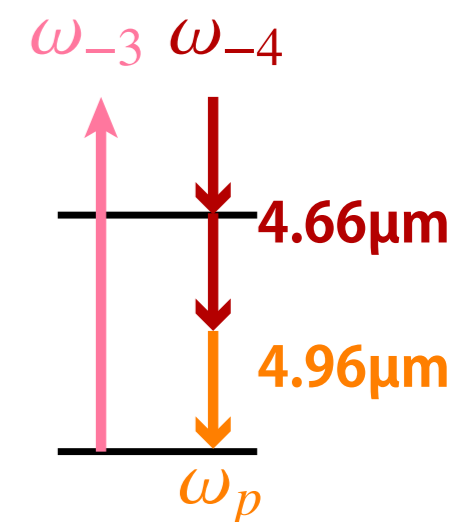
- 全体的な傾向は再現
- $q \geq 3$ の再現性は悪い

平均コヒーレンス 0.032

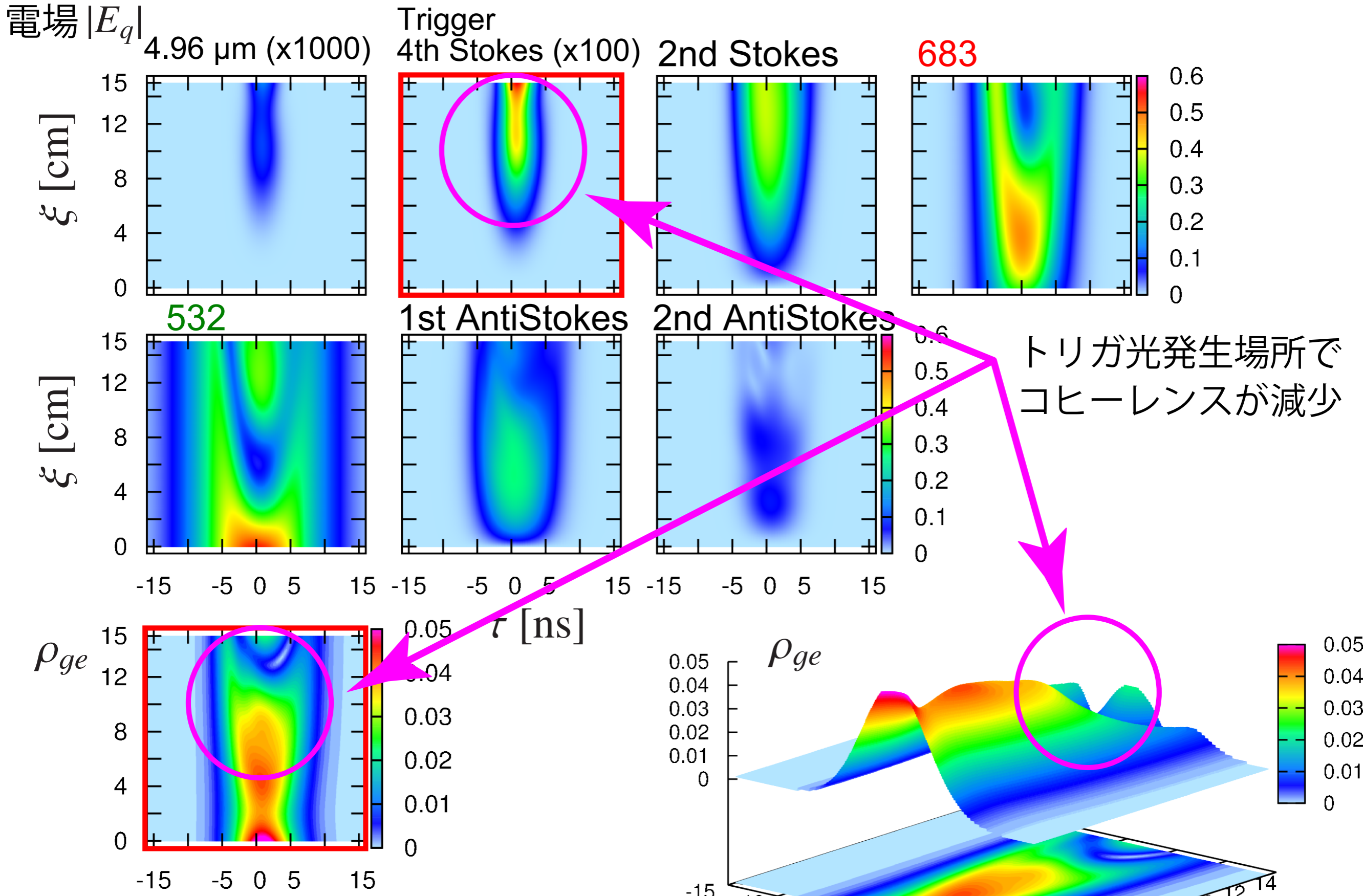
Detuning Curve of Two-Photon Emission



- Detuning curve の形状は一致 (測定誤差範囲内)
- $E_{4.96} / E_{4.66} = 10^{-3}$ two-photon pair の強度比は一致



Simulation モデルの妥当性を確認



- 全体の平均値: $\rho_{eg} = 0.032$ ($\tau = 0$)
- トリガ光周辺の平均値: $\rho_{eg} = 0.022$

External trigger exp. setup

