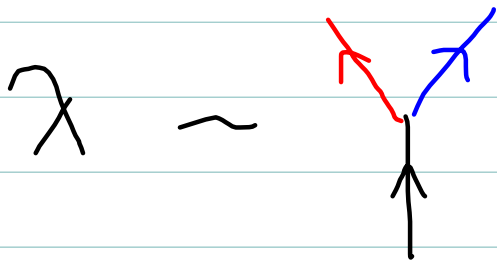


QFT



choose types of particles in a theory.

choose how they interact.



STRING

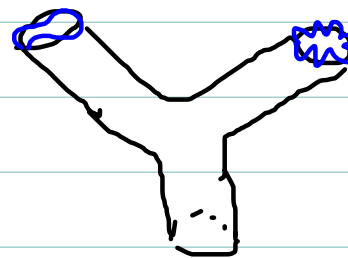


choose a type of string.

study how they vibrate.



interaction is then determined.



There are **two** known consistent types of strings.



Type II string



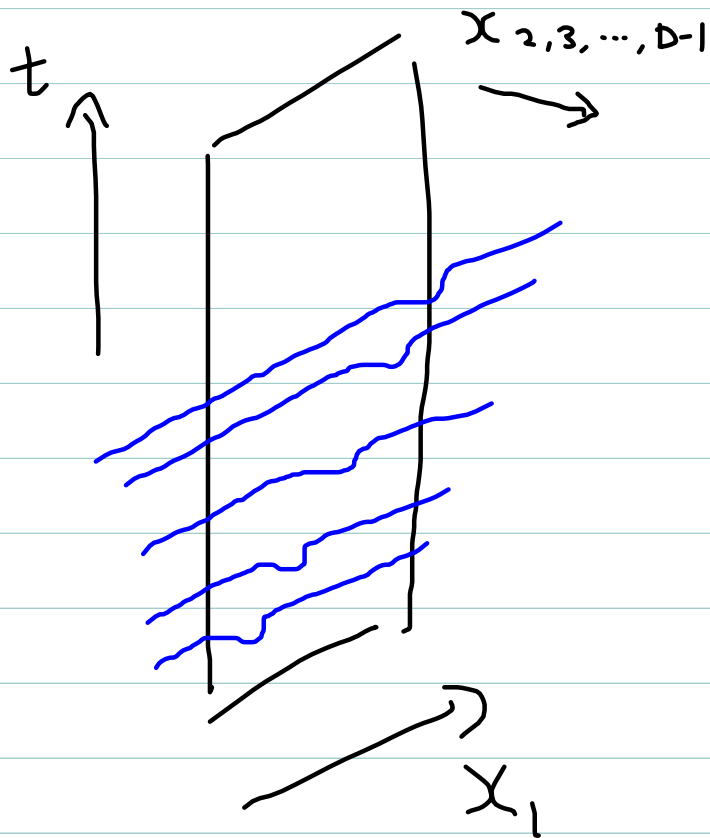
Heterotic string

Both require $1+9$ dimensional spacetime.
have supersymmetry.
Constrains quantized gravity.

All these follow from the consistency of
Lorentz inv. & **Positivity of probability.**

I'd like to give you some flavor of the derivation.

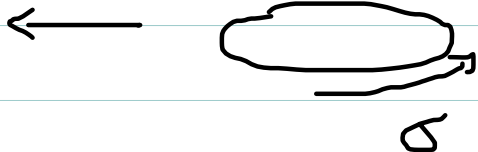
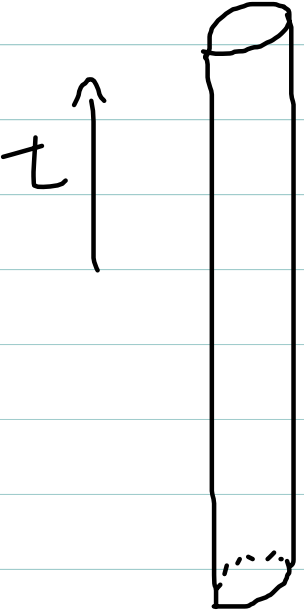
Consider a string moving in D -dim'l spacetime.



The config. of a string is specified by giving

$$x_2(t, x_1), x_3(t, x_1), \dots, x_{D-2}(t, x_1), x_{D-1}(t, x_1)$$

We have D-2 scalar field on the 2d space (t, x_1) .



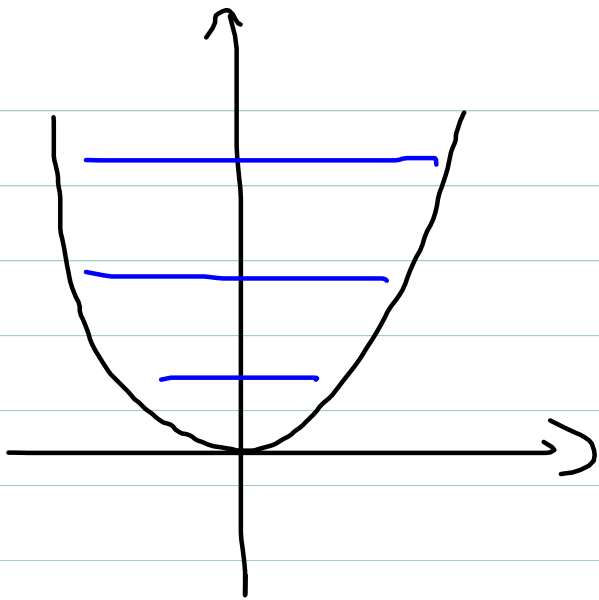
$$X_i(\sigma) = \sum a_{i,k} e^{2\pi i k \sigma} + a_{i,k}^\dagger e^{-2\pi i k \sigma}$$

$$\begin{cases} i = 2, 3, \dots, D-2 \\ k = 1, 2, 3, \dots \end{cases}$$

Each $a_{i,k}, a_{i,k}^\dagger$ is a harmonic oscillator.
 $|0\rangle, a_{1,1}^\dagger |0\rangle, (a_{1,1}^\dagger)^2 |0\rangle, \dots$



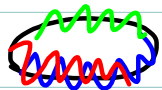
The same mode
excited twice!



$$\leftarrow \frac{5}{2}\hbar\omega$$

$$\leftarrow \frac{3}{2}\hbar\omega$$

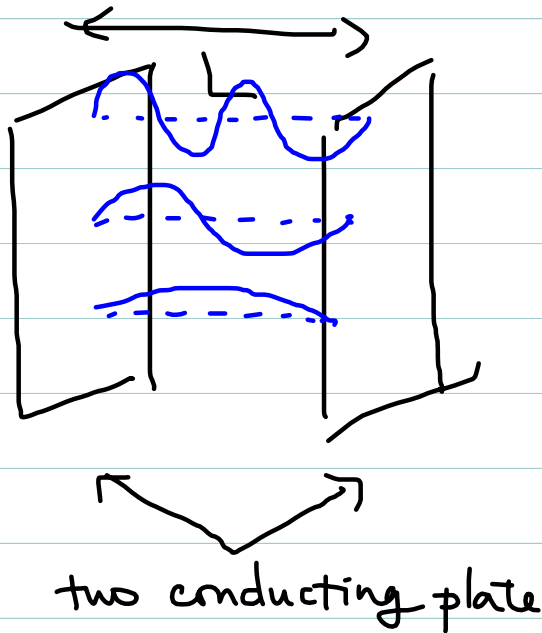
$$\leftarrow \frac{1}{2}\hbar\omega : \text{zero point energy.}$$



$$a_k^\dagger a_{k'}^\dagger \dots a^\dagger \dots a_{k''}^\dagger |0\rangle$$

$$\text{has } m^2 = \frac{1}{c_s^2} \left[(k+k'+\dots+k''') + \text{zero pt energy.} \right]$$

ASIDE the Casimir effect



The sum of the zero pt energy / area

$$\sim \frac{1}{L^3} (1^3 + 2^3 + 3^3 + \dots)$$


$$= - \frac{\pi^2}{720} \frac{1}{L^3} \quad \text{after renormalization.}$$

This is MEASURED.

(Lamoreaux, PRL 78 ('97) 5)

So, there's Casimir energy on the string:

$|0\rangle$



without
any excitation

$$\dots m^2 = \frac{1}{\alpha_s^2} \cdot L \left(-\frac{1}{24L} \right) \times (D-2)$$

\uparrow Casimir energy
of one $X_i(t, x_i)$

\uparrow $i = 2, 3, 4, \dots$
 $D-1$.


$a_{i,1}^\dagger |0\rangle$


with one
excitation of
the lowest mode

$$\dots m^2 = \frac{1}{\alpha_s^2} \left(-\frac{D-2}{24} + 1 \right).$$

Note that there are states
for $i = 2, 3, \dots, D-1$

It's a vector boson.

 gives $D-2$ polarizations of a vector boson
with $mass^2 = \frac{1}{\alpha_s^2} \left(-\frac{D-2}{24} + 1 \right)$.

This makes sense only when $m^2=0 \rightarrow \underline{D=26}$

Recall: in 4d, massless vector: 2 polarizations
massive vector: 3 polarizations.

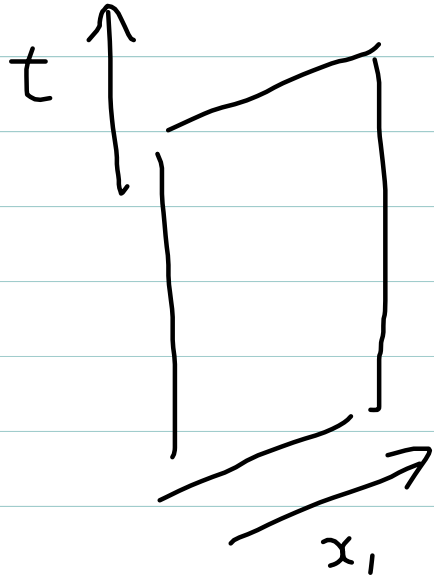
In general in D dim, massless $\rightarrow D-2$ pol.
massive $\rightarrow D-1$ pol.

But, even when $D=26$,

The lowest mode

 : $m^2 = -\frac{1}{\alpha_s^2}$: extremely tachyonic!

The cure add supersymmetry.



$X_2(t, x_1) \dots X_{D-2}(t, x_1)$
spacetime fluctuation

$\psi_2(t, x_1)$
fermionic fluctuation.

Casimir energy:

$$(D-2) \sum (1+2+\dots) - (\text{\# fermion components}) \sum (1+2+\dots)$$

This can be 0, removing the tachyon.

$$\rightsquigarrow \underline{\underline{D=10}}$$

Summary

Quantizing particles can be done in any D ,

Quantizing strings can only be done in

$$D=10.$$

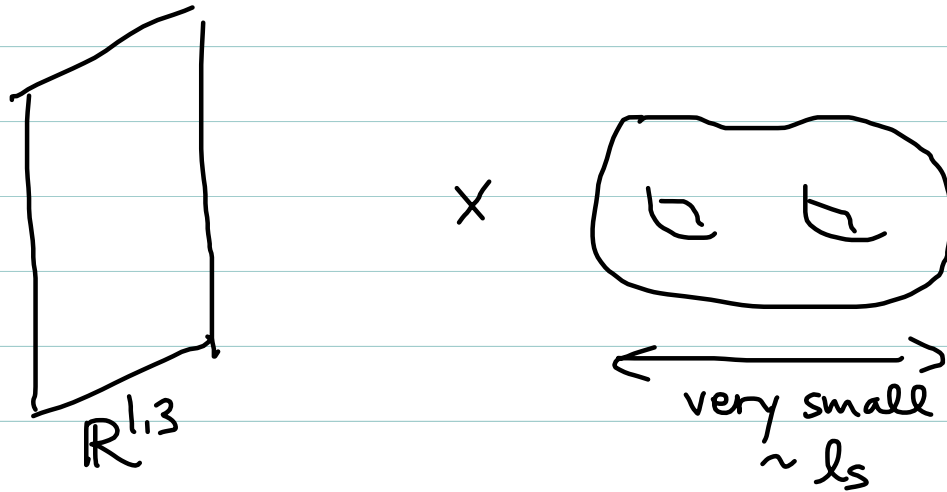
You say, wait, our spacetime is $1+3$ d.

\rightsquigarrow We need to make 6 out of 10 dimensions
very small so that
they are not observed.

(NOTE THAT if we lived in 3D dimensions, string theory
would have been already experimentally falsified.)

Quantizing strings \rightarrow Only fully consistent in
10 d and with SUSY.

To say anything about 4d physics, we need to compactify
the extra 6d.



Then it's important to study the geometry of the 6d internal space.
That's why a lot of math is necessary.

Before going there, let's study 10d physics first.



Strings have **right-moving** & **left-moving** excitations.

Yesterday I only talked about one of them.

$a_{j,1}^\dagger a_{i,1}^\dagger |0\rangle$ has two spacetime indices, i & j .
 $i, j = 2, 3, \dots, D-1=9$

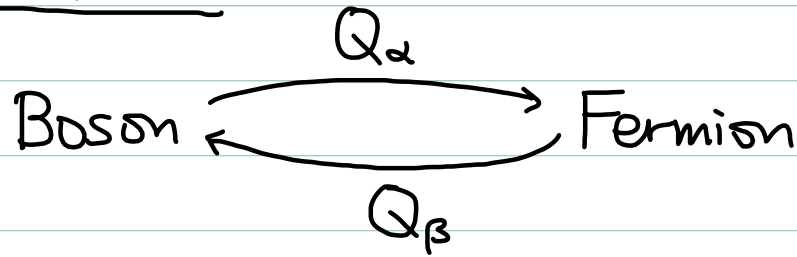
\rightsquigarrow contains graviton!

(cf. $g_{\mu\nu}$ has two indices)

The system also has supersymmetry.

\rightsquigarrow 10d **SUPERGRAVITY**.

SUPERSYMMETRY



$$\{Q_\alpha, Q_\beta^\dagger\} = \gamma_{\alpha\beta}^\mu P_\mu \quad \rightarrow \quad Q_\alpha \text{ is like } \sqrt{\text{translation}}$$

translation

squark \leftrightarrow quark
gauge boson \leftrightarrow gaugino
graviton \leftrightarrow gravitino
⋮

This is possible only in 11d or less !

BOSON \leftrightarrow FERMION

\rightsquigarrow # of boson degrees of freedom
= # of fermion d.o.f.

$g_{\mu\nu}, A_\mu, \dots$: D, D^2 components, roughly speaking.

ψ_a : How many components? γ^M needs to act on it.

In 4d, γ^M is a 4×4 matrix. $\{\gamma^M, \gamma^N\} = \eta^{MN}$.

In D dimensions, let $b_1 = \gamma^0 + \gamma^1$, $b_1^\dagger = \gamma^0 - \gamma^1$
 $b_2 = \gamma^2 + i\gamma^3$, $b_2^\dagger = \gamma^2 - i\gamma^3$

\vdots

$b_{D/2} = \gamma^{D-1} + i\gamma^D$, $b_{D/2}^\dagger = \gamma^{D-1} - i\gamma^D$

They satisfy standard fermion commutator: $\{b_i, b_j^\dagger\} = \delta_{ij}$.

So, a Dirac spinor in D spacetime dimensions is like having $D/2$ fermions. eg when $D=4$, $D/2=2$.

$$\begin{array}{cc} |00\rangle & |01\rangle \\ |10\rangle & |11\rangle \end{array} \rightsquigarrow 4 \text{ states.}$$

This was why γ^M is 4×4 . $4 = 2^2$

In general, a Dirac spinor has $2^{D/2}$ components.

You can impose Weyl or Majorana conditions ...

Boson : $\phi, A_\mu, g_{\mu\nu}$: 1, D, D^2 components



Fermion $\psi_\alpha, \psi_{\alpha\mu}$: $2^{D/2}, D 2^{D/2}$ components

grows much faster
than bosons.

SUSY \rightarrow #boson d.o.f = # fermion d.o.f
 \rightarrow D can't be very big.

Maximum is $D=11$. There's only one theory there.

11d supergravity

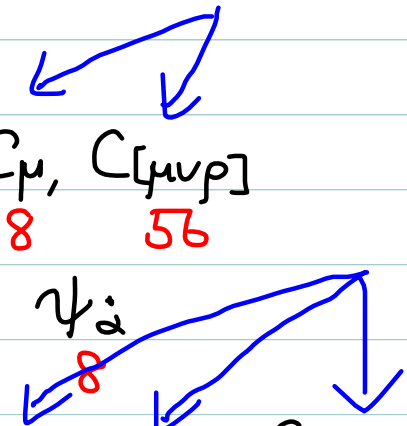
	dof	dof
graviton $\rightarrow g_{\mu\nu}$	44	
$\nearrow C_{[\mu\nu\rho]}$	84	
C-field		
		$\psi_{\alpha\mu}$
		\uparrow
		gravitino
		128

That's it. You can't add gauge field, scalar, ...

If we lived in 11d, and had SUSY, the theory is
uniquely defined.

10d supergravity : 3 types.

Type IIA	boson	$g_{\mu\nu}$	$B_{\mu\nu}$	ϕ	C_μ	$C_{[\mu\nu\rho]}$	
	128	35	28	1	8	56	
	fermion	$\psi_{\mu\alpha}$	$\psi_{\mu\dot{\alpha}}$	ψ_α	$\psi_{\dot{\alpha}}$		
	128	56	56	8	8		
Type IIB	boson	$g_{\mu\nu}$	$B_{\mu\nu}$	ϕ	C	$C_{[\mu\nu]}$	$C_{[\mu\nu\rho\sigma]}$
	128	35	28	1	1	28	35
	fermion	$\psi_{\mu\alpha}$	$\psi'_{\mu\alpha}$	ψ_α	ψ'_α		
	128	56	56	8	8		



Again, you can't add gauge fields for type IIA, IIB supergravity.

But they have D-branes.

Type I = Heterotic

boson	$g_{\mu\nu}, B_{[\mu\nu]}, \phi$
64	35 28 1
fermion	$\psi_{\mu\alpha}, \psi_{\beta}$
64	56 8

← superpartners of gravitons

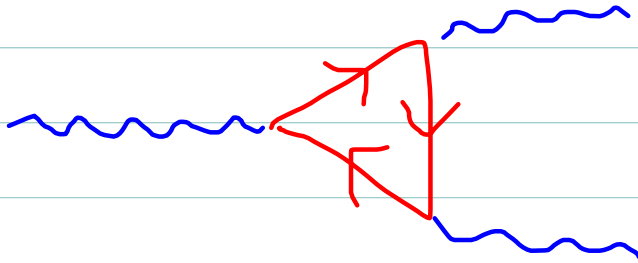
boson	A_{μ}^a
8	8
fermion	λ_{β}^a
8	8

← superpartners of gauge fields

Now we can have gauge fields. But only when

$$G = SO(32) \quad \text{or} \quad E_8 \times E_8.$$

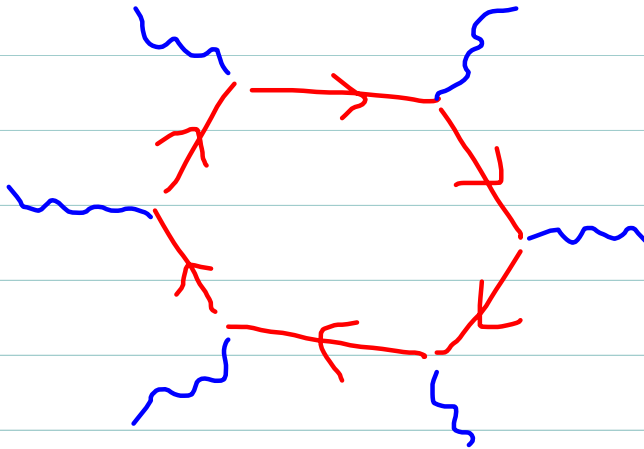
In 4d, with chiral fermions,



caused trouble when

$$\sum g_L^3 - \sum g_R^3 \neq 0.$$

In 10d type I, fermions are always chiral.



causes trouble ...

unless $G = SO(32)$
or $E_8 \times E_8$.

So, in 10d + SUSY, there're only 4 possibilities:

Type IIA,

Type IIB

Heterotic SO(32),

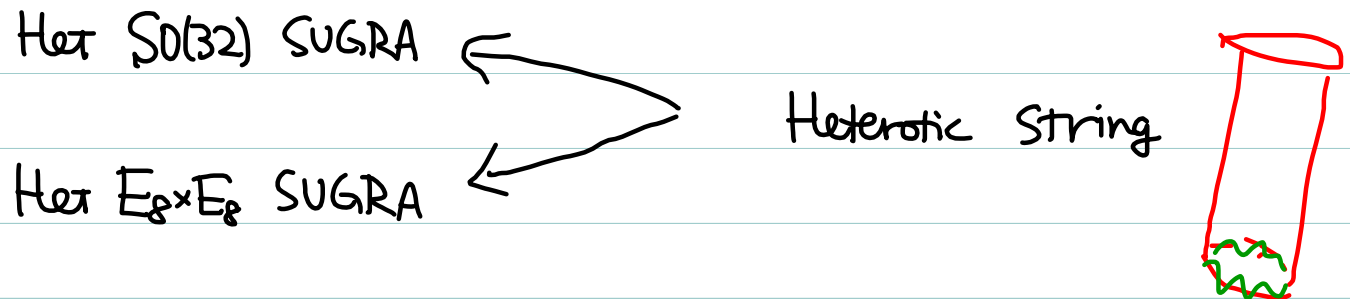
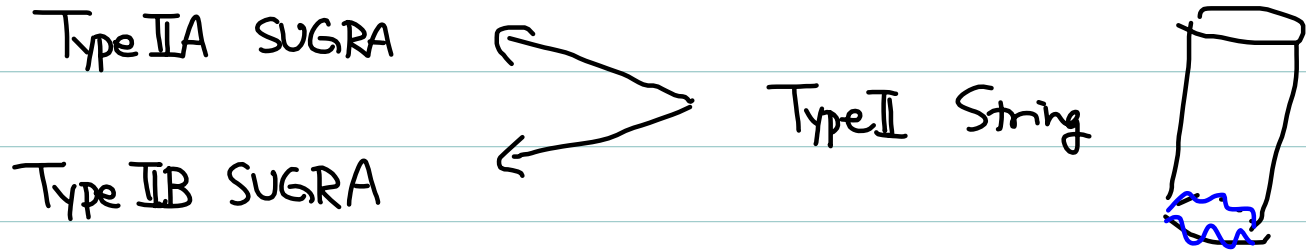
Heterotic $E_8 \times E_8$.

Then, every interaction vertex is fixed.

It's not like in 4d, where we can choose
gauge group, fermions, scalars, Yukawa couplings...

These theories are very NON-RENORMALIZABLE.
in the standard QFT sense. (contains gravity!)

But in fact, there're corresponding string theories



10d supergravities are successfully quantized by strings!

This is one reason why string theorists hope that similar things might work in 4d ...