

Simulation Study for Electron
Beam in the Main Linac of
the Linear Collider

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Introduction

In the main linac, **emittance-blowup** can be produced by

- **non-linear magnetic field** in quadratic magnets
- **beam-scattering** on residual gas and thermal photon

➔ We have a simulation for the **GLC main linac** and estimate the effect in detail.

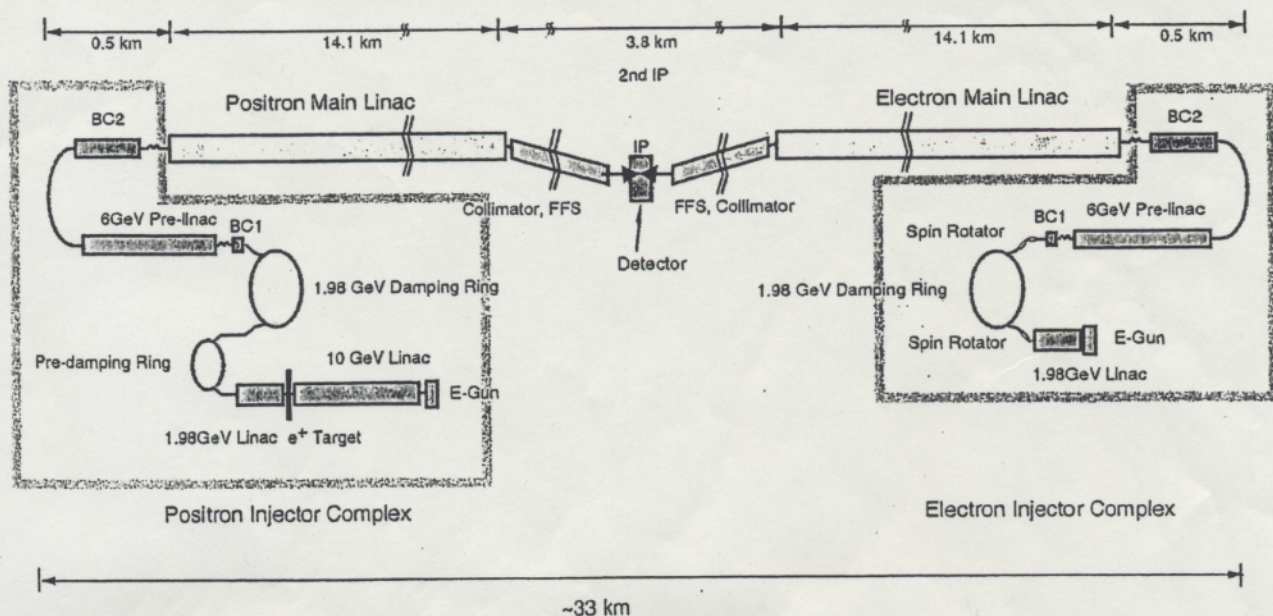


Fig. Schematic view of GLC

Tracking Simulation

©Tracking code

★SAD(Strategic Accelerator Design)

- It calculates the orbit of electrons by matrix method.
- Lattice design is to be input.
- It has been developed by KEK.

©Monte Carlo simulation for the beam-scattering is performed by our own program.

Tracking Simulation(2)

© Parameters for accelerator and beam

We mainly use the parameters designed for the **stage- II of GLC** ($\sqrt{s} = 1\text{TeV}$).

© Core or Halo ?

Only core is studied in this simulation.

Parameters

©Parameters for GLC

initial energy ... 8GeV

final energy ... 500GeV

total length ... 12.8km

number of electrons ... 7.5×10^9 /bunch

RF frequency ... 11.424GHz (X-band)

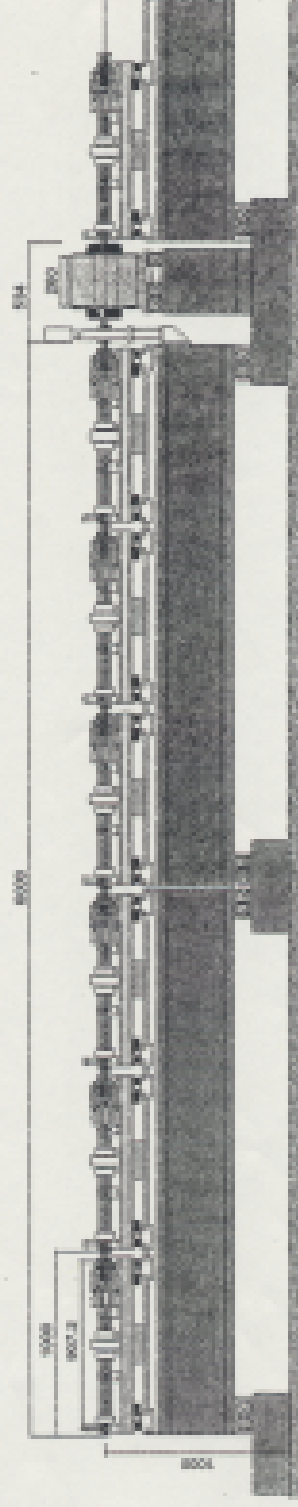


Fig: GLC main linac

Parameters(2)

©Other assumptions

Pressure in the linac ... 10nTorr

Temperature in the linac ... 300K

©Misalignment and wakefield are not considered in this study.

Initial State of the Beam Core

Distribution in phase space

... elliptic in $x-x'(y-y')$ plane

$$\left[\begin{array}{l} x' = dx/ds, \quad y' = dy/ds \\ s; \text{ coordinate along the beam line} \end{array} \right]$$

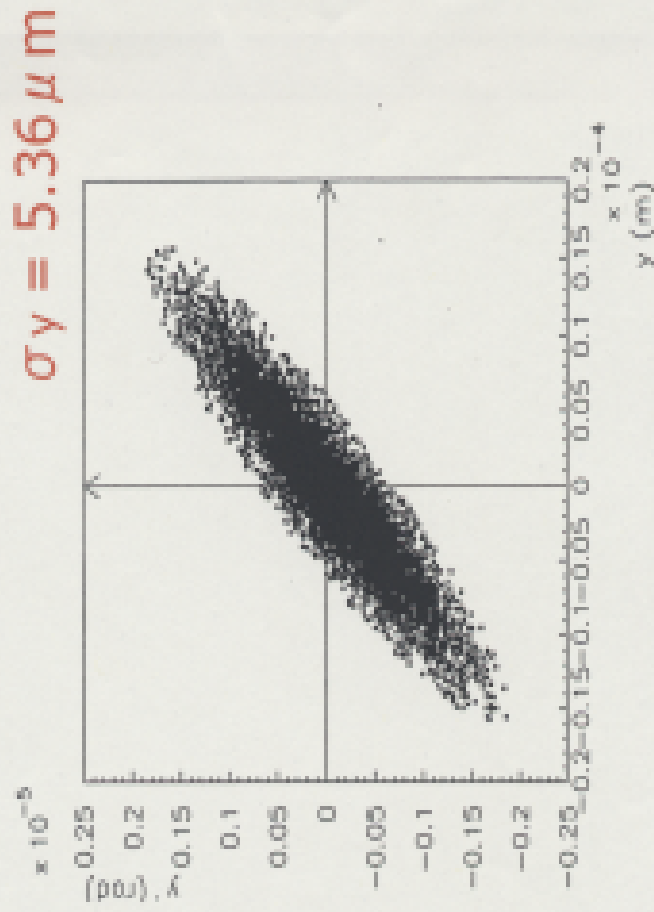
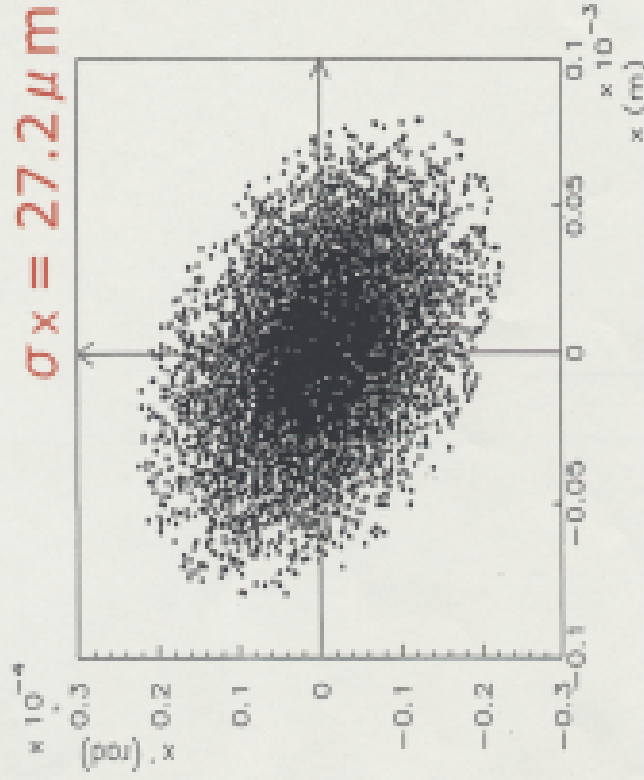
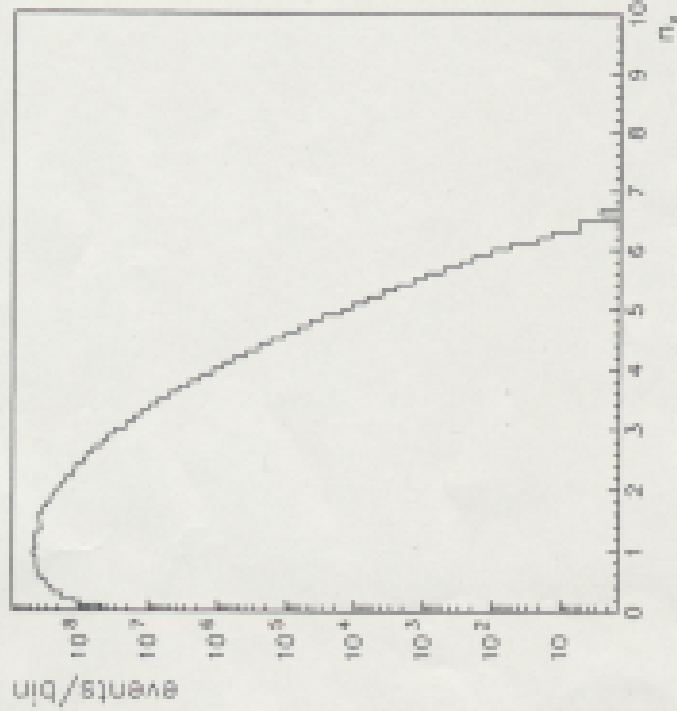


Fig: Phase space distribution at the entrance of the linac
(※ Only electrons within 3σ are plotted here.)

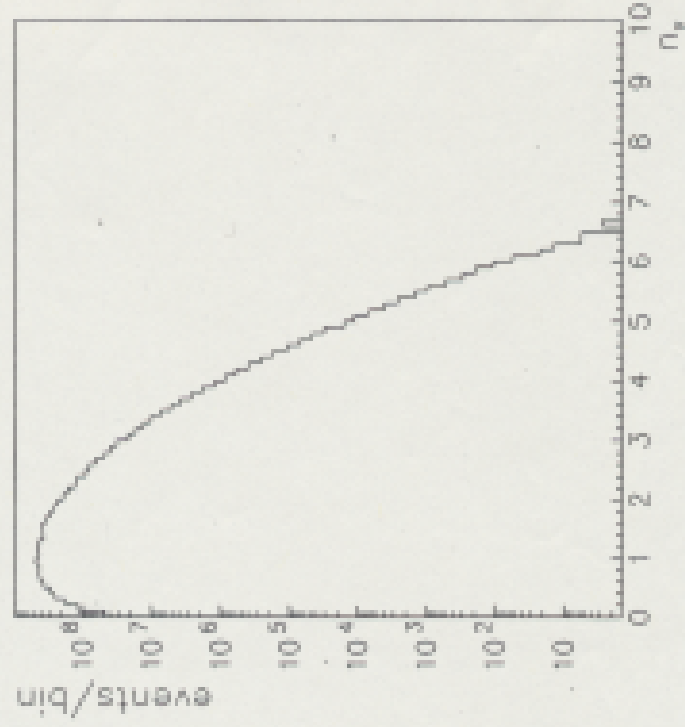
Initial State of the Beam Core (2)

Density distribution in phase space

- We assume **Gaussian distribution**.
- Tail : Up to **6-7 σ** is considered.



Fig(a): n_x distribution



Fig(b): n_y distribution

($n_x[n_y]$: number of sigma in $x-x'[y-y']$ plane)

Emittance $\epsilon_x, \epsilon_y \dots$ parameters which show the spatial spread of the beam

$$\epsilon_x = (\sigma_x^2 \sigma_{x'}^2 - \sigma_{xx'}^2)^{1/2}$$

$$\sigma_x^2 = \langle x^2 \rangle - \langle x \rangle^2$$

$$\sigma_{x'}^2 = \langle x'^2 \rangle - \langle x' \rangle^2$$

$$\sigma_{xx'} = \langle (x - \langle x \rangle) (x' - \langle x' \rangle) \rangle$$

normalized emittance ... expected to be conserved

$$\epsilon_N = \gamma \beta \epsilon \quad (\epsilon \sim 1 / \gamma \beta)$$

Non-linear Magnetic Field

In each quadrupole magnet, **non-linear magnetic field** can be produced by **manufacturing error of the magnet.**



In the simulation, we investigate the effect of sextupole magnet .
(the largest contribution of all non-linear components)

Simulation for Sextupole Magnetic Field

- **Sextupole magnet** is installed (virtually) just before each quadrupole magnet.

- Strength: $\left| \frac{B_3}{B_2} \right| = 10^{-3}$

- Direction of the sextupole: **random**

General form of the magnetic field

$$B_x + i B_y = \sum_{n \geq 2} B_n \cdot \left(\frac{x + i y}{a} \right)^{n-1} \sim \sum B_n \cdot \left(-\frac{r}{a} \right)^{n-1}$$

$$\left[\begin{array}{l} B_n : \text{complex constant} \\ a = 20\text{mm} : \text{aperture} \times 0.5 \end{array} \right]$$

$n=m$ \longleftrightarrow (2m)-pole field

($n=2$: quadrupole $n \geq 3$: non-linear)

In KEKB and Tristan,

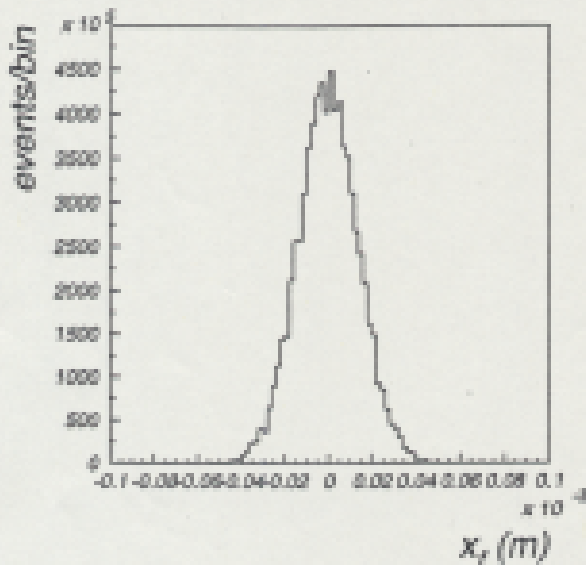
$$\frac{\sum_{n \geq 3} |B_n|}{|B_2|} \lesssim 10^{-3}$$

was achieved. (K.Egawa)

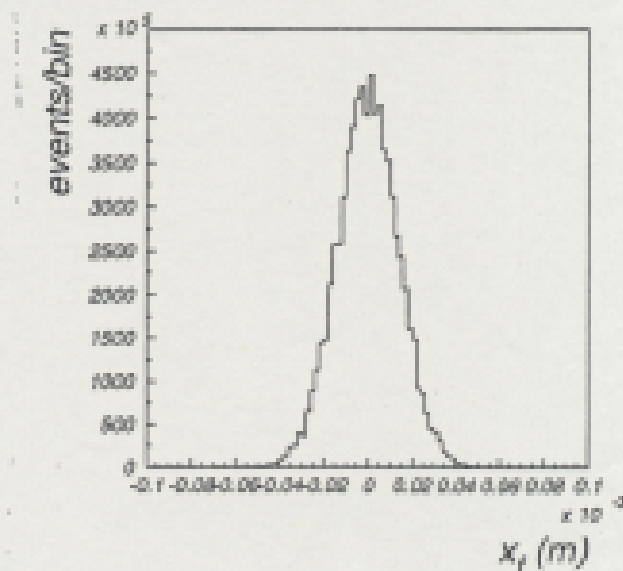
effect on x_f distribution

(x_f ; position in x at the exit)

- linear scale for vertical axis

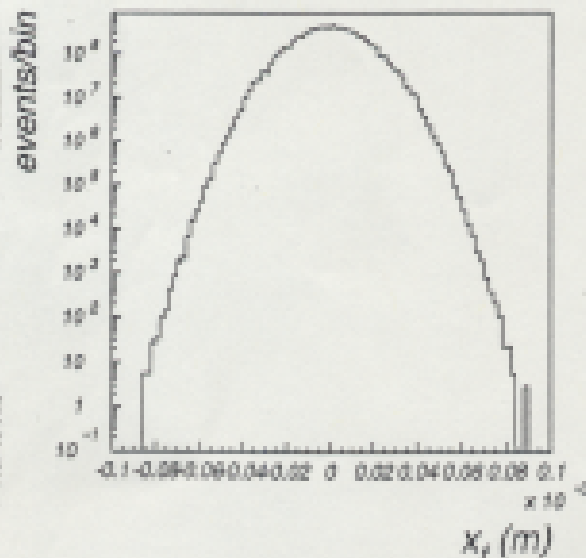


$$|B_3/B_2| = 0$$

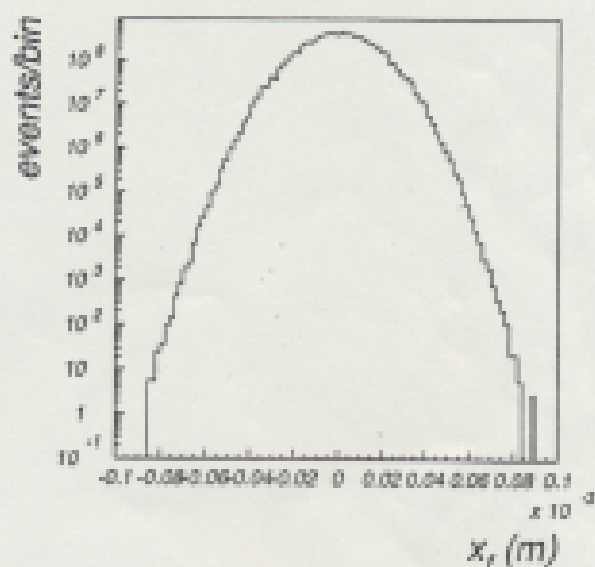


$$|B_3/B_2| = 10^{-3}$$

- log scale for vertical axis



$$|B_3/B_2| = 0$$

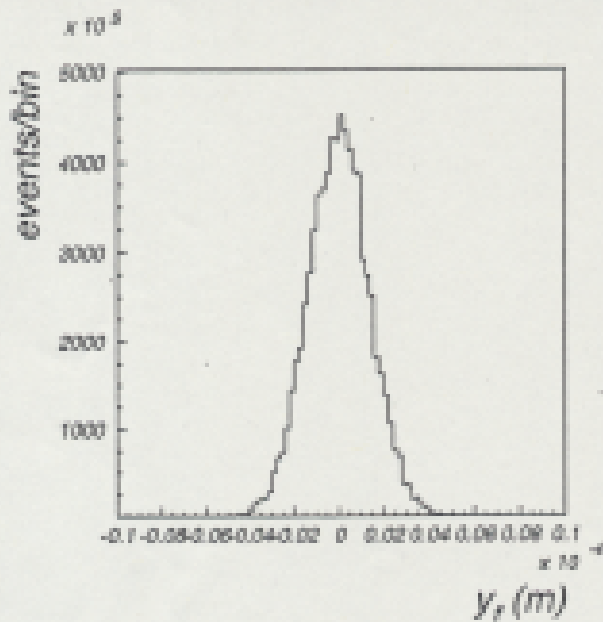


$$|B_3/B_2| = 10^{-3}$$

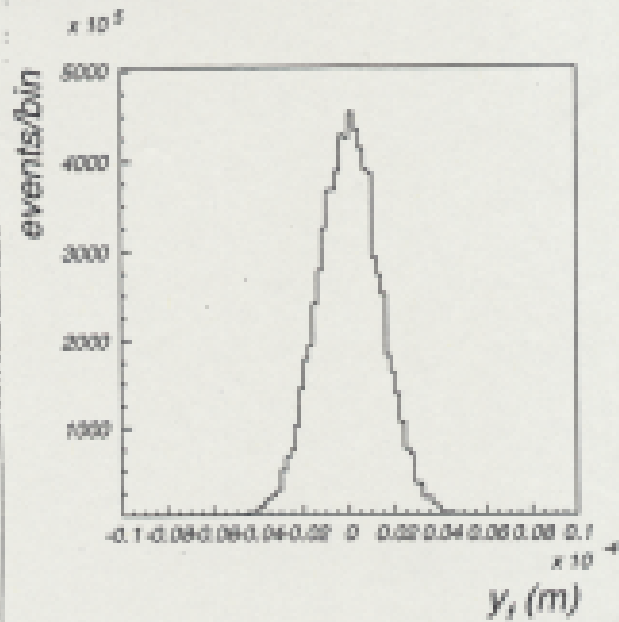
effect on y_f distribution

(y_f ; position in y at the exit)

- linear scale for vertical axis

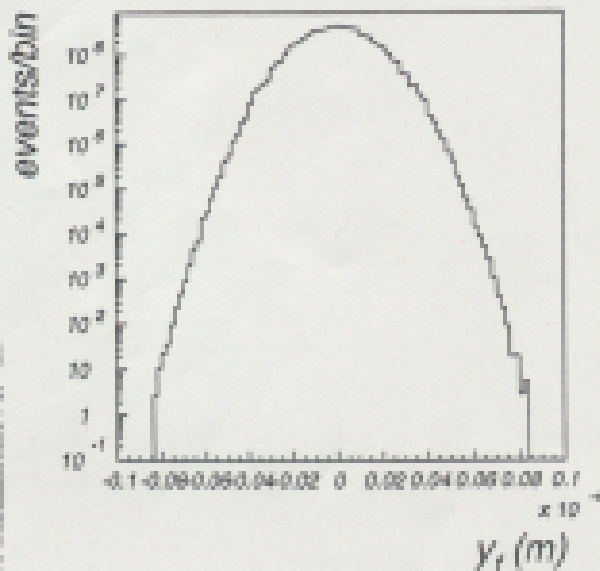


$$|B_3/B_2| = 0$$

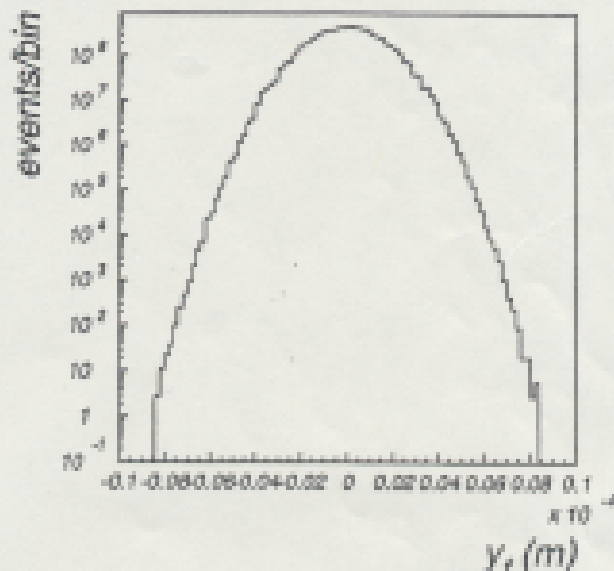


$$|B_3/B_2| = 10^{-3}$$

- log scale for vertical axis



$$|B_3/B_2| = 0$$

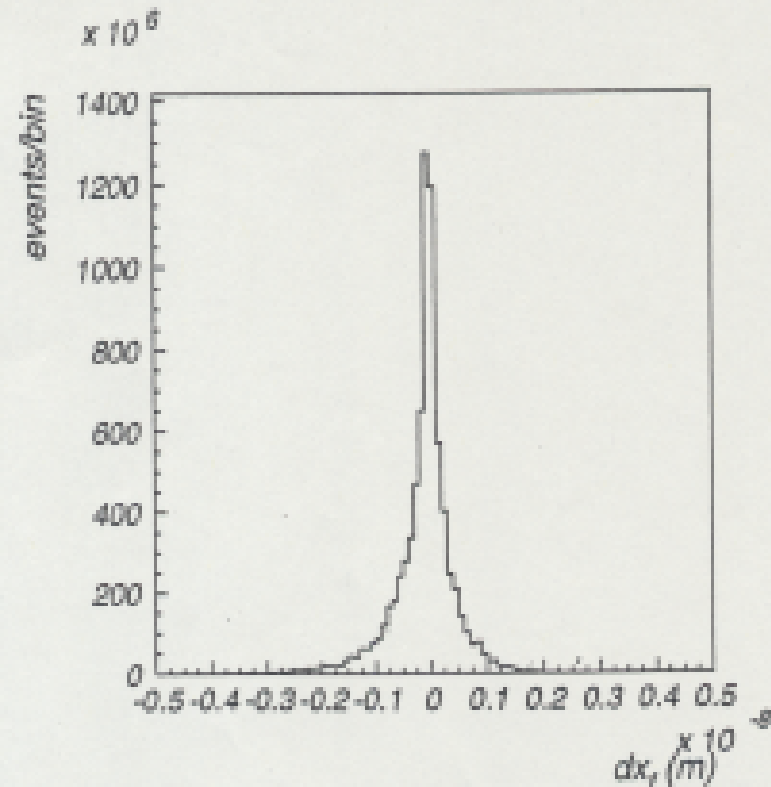


$$|B_3/B_2| = 10^{-3}$$

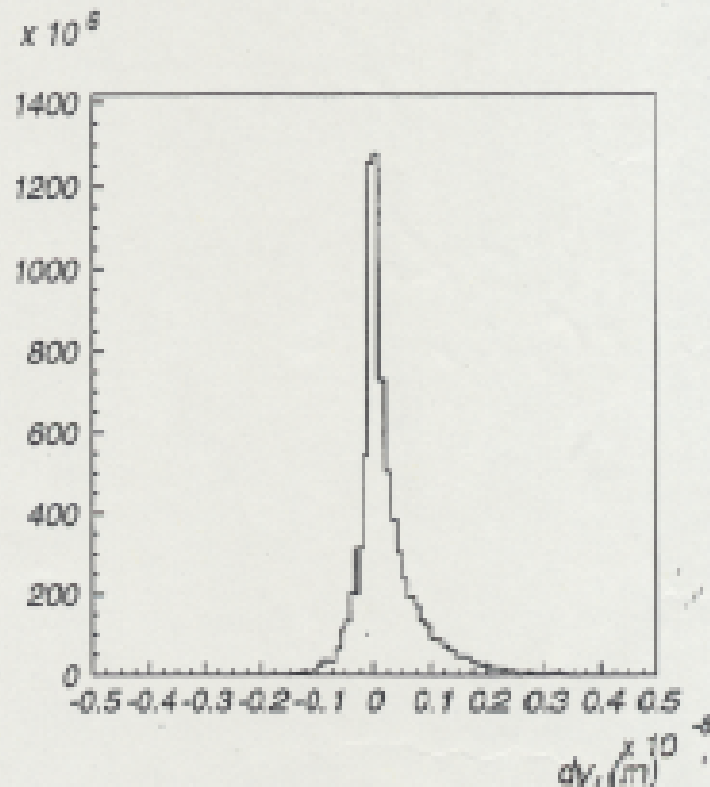
dx_f, dy_f distribution

dx_f, dy_f ; change of x_f, y_f due to the effect of the sextupole magnet

$$\begin{cases} dx_f = X_f(B_3 \neq 0) - X_f(B_3 = 0) \\ dy_f = Y_f(B_3 \neq 0) - Y_f(B_3 = 0) \end{cases}$$

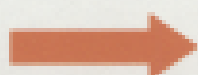


dx_f distribution



dy_f distribution

Distribution spread: $\sim O(1)$ nm



Influence on the beam core with its size of $\sim \mu m$ is almost negligible.

Emittance Blowup by the Sextupole Magnetic Field

The effect is negligible as long as $|B_3/B_2| < 10^{-1}$ - 1.

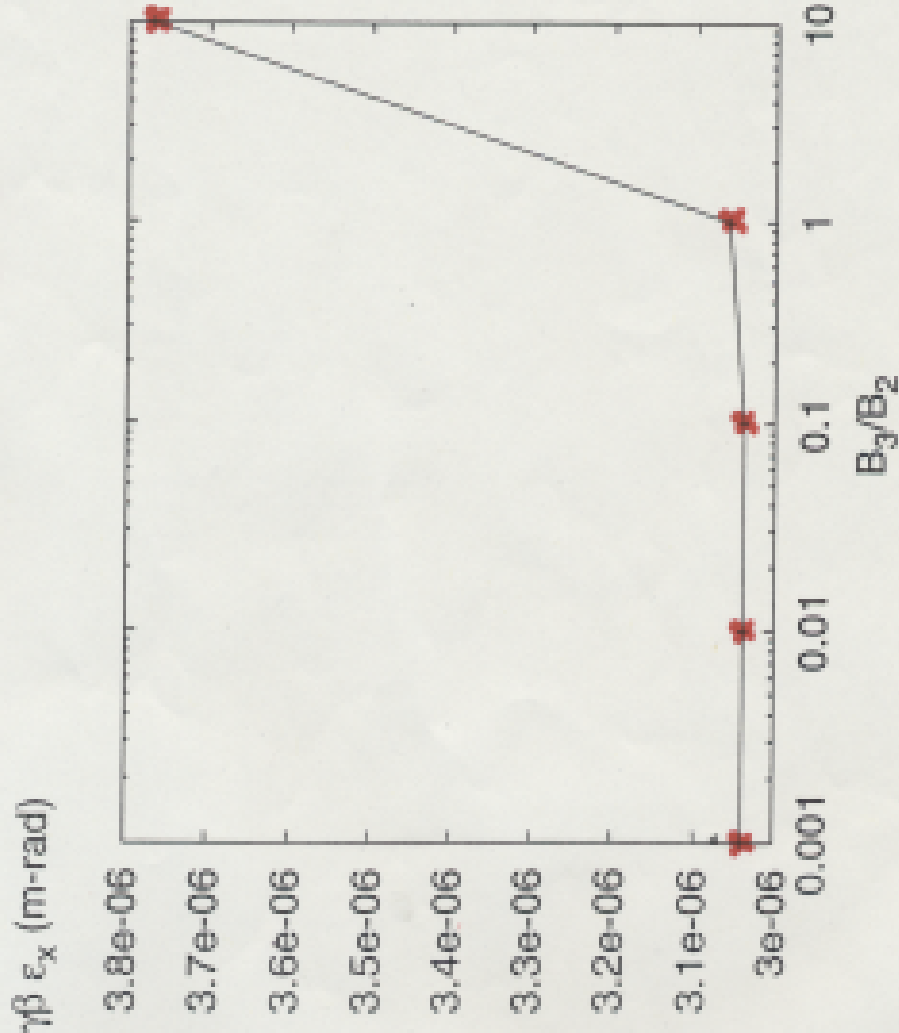


Fig: $\gamma \beta \epsilon_x$ -blowup

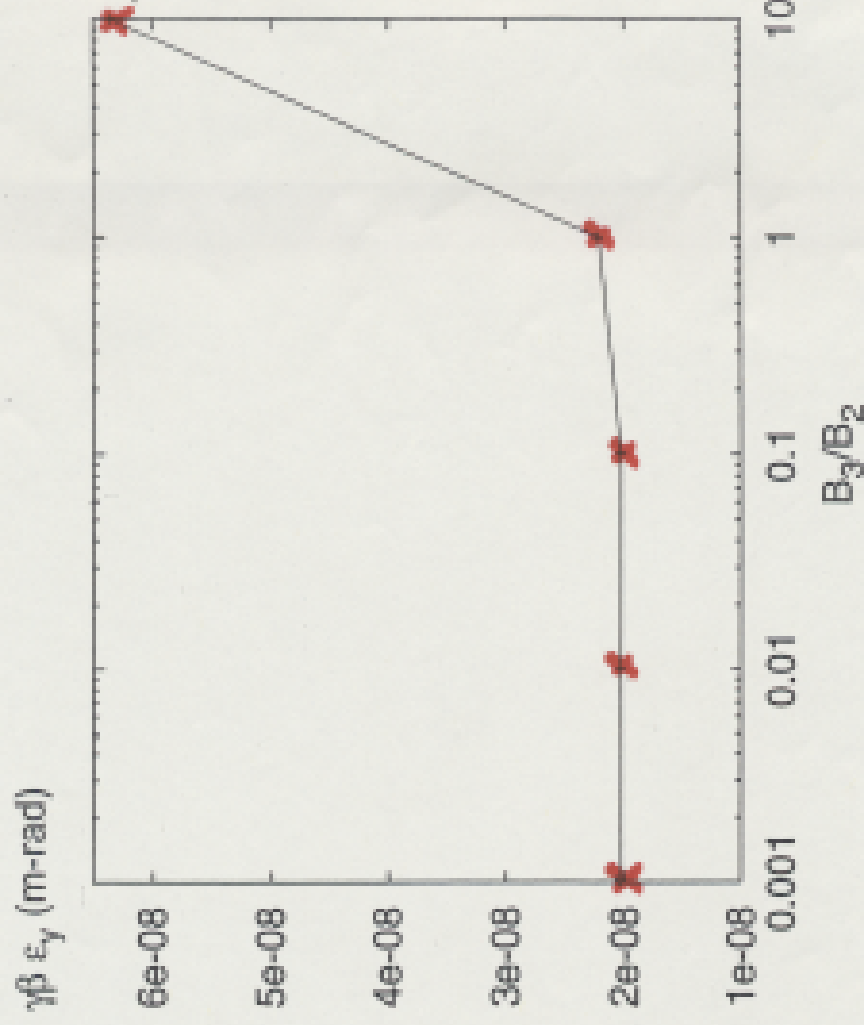


Fig: $\gamma \beta \epsilon_y$ -blowup

Beam Scattering

© Scattering process with large contribution in the linac

○ Interaction with residual gas

- elastic scattering
- bremsstrahlung

○ Compton scattering

on thermal photon

© Assumptions

- Pressure . . . 10nTorr
- Temperature . . . 300K
- Gas component . . . N_2 only

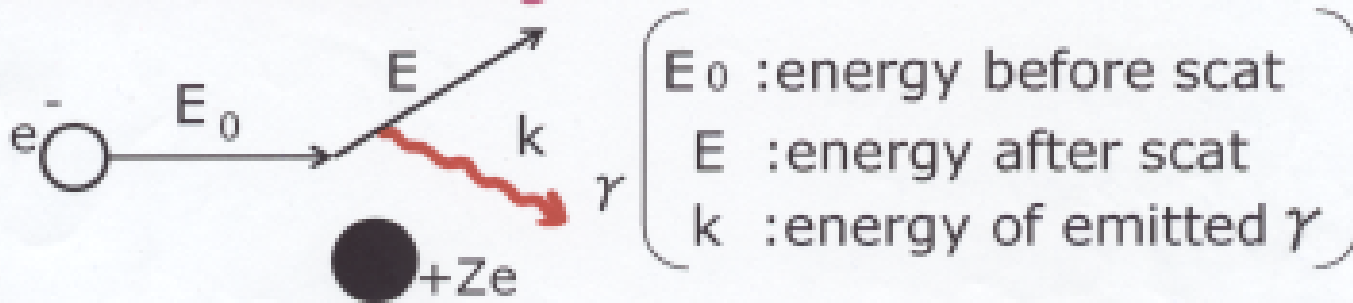
Number of Scattered Particles per Bunch in the Linac

©Elastic scattering

$$\sigma_{el} \sim 0(10^5) \text{ barn}$$

$$N_{el} = 6.04 \times 10^5$$

©Bremsstrahlung



Infrared divergence ;

$$\sigma_{brems} \sim \ln(x_{min})$$

$$\text{At } x_{min} = 10^{-2}, \quad \left(x = \frac{k}{E_0} : \text{energy loss rate} \right)$$

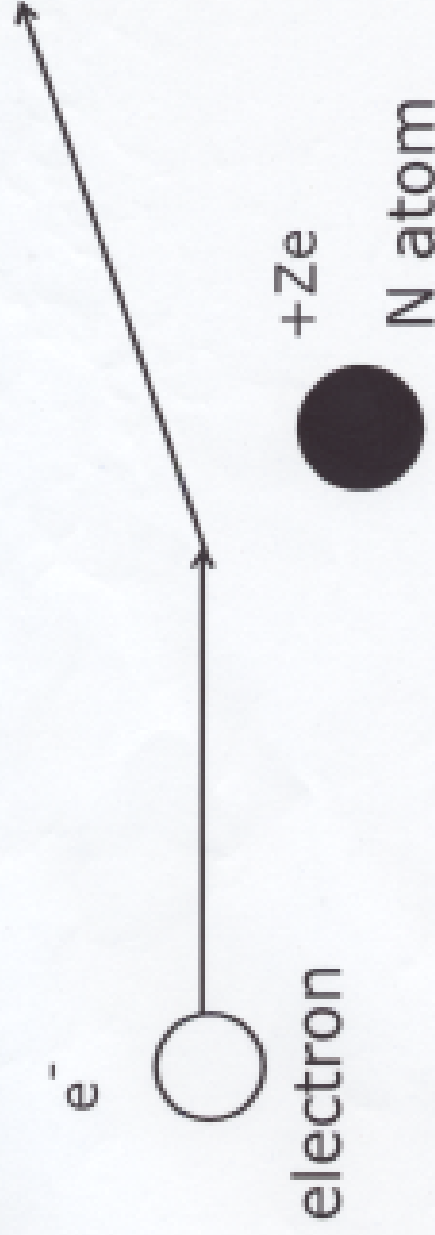
$$\sigma_{brems} \sim 0(1) \text{ barn}$$

$$N_{brems} \sim 0(10)$$

©Compton Scattering

$$N_{comp} = 3.15$$

Simulation for Elastic Scattering



• Elastic scattering occurs uniformly in the linac.

$$\frac{d\sigma}{d\theta} \sim \frac{1}{\theta^3}, \quad \theta_{\min} = \frac{h}{pa} \quad \begin{array}{l} \text{due to} \\ \text{screening effect} \end{array}$$

$$(\leftrightarrow \theta \sim h/pb)$$

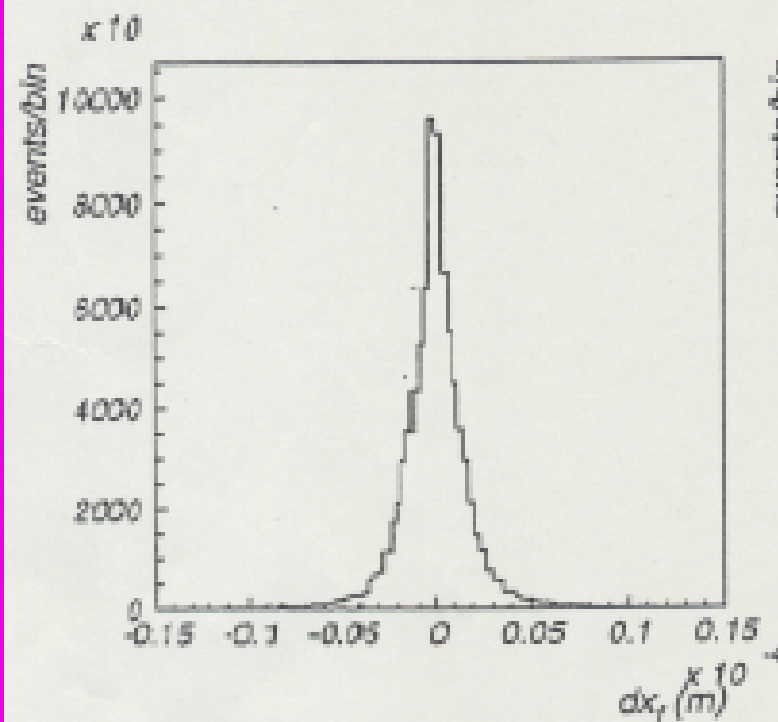
b: impact parameter)

(a: atomic radius of nitrogen
p: incident momentum of e^-)

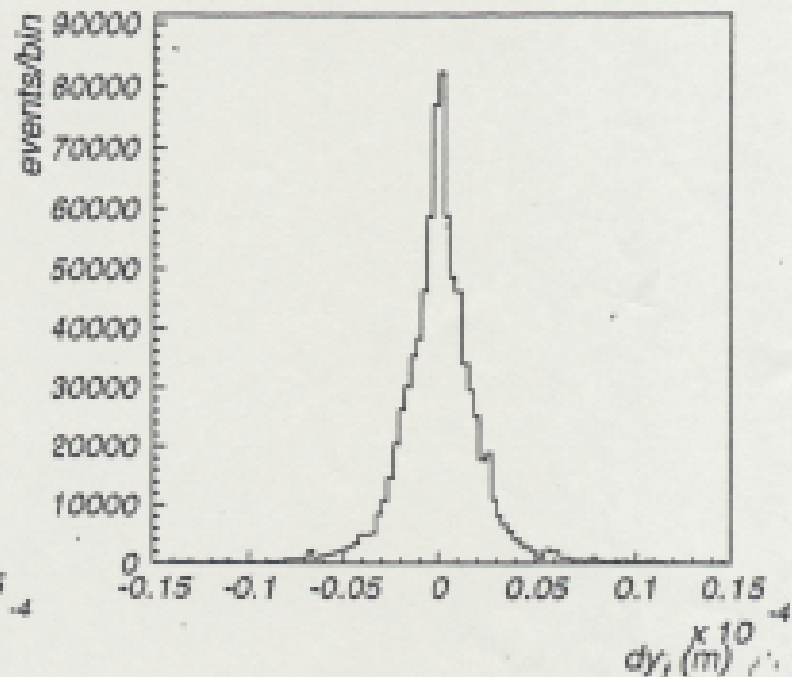
dx_f, dy_f distribution of scattered electrons

dx_f, dy_f : change of x_f, y_f due to the scattering

$$\begin{cases} dx_f = X_{f, \text{scat-occurred}} - X_{f, \text{scat-not-occurred}} \\ dy_f = Y_{f, \text{scat-occurred}} - Y_{f, \text{scat-not-occurred}} \end{cases}$$



dx_f distribution



dy_f distribution

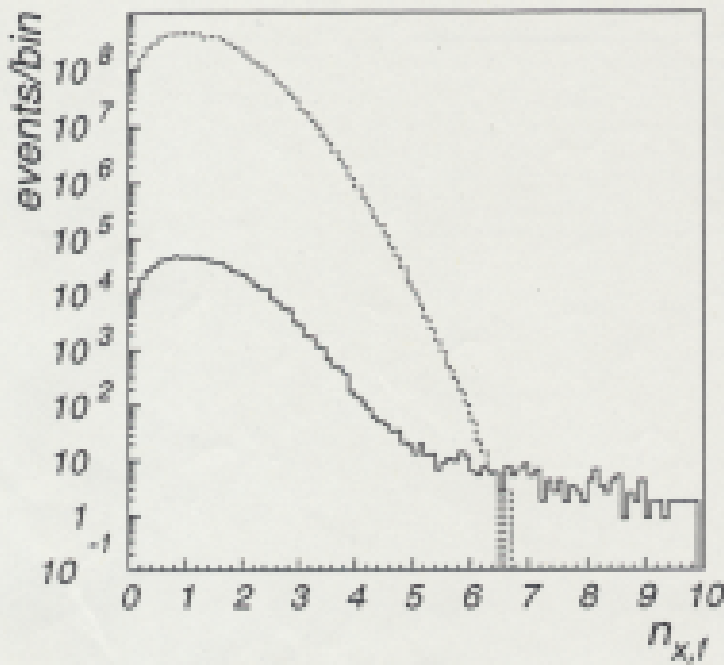
The distribution spread reaches as much as $\mu\text{m scale}$.

cf.) $dx, dy \sim O(1)\text{nm}$

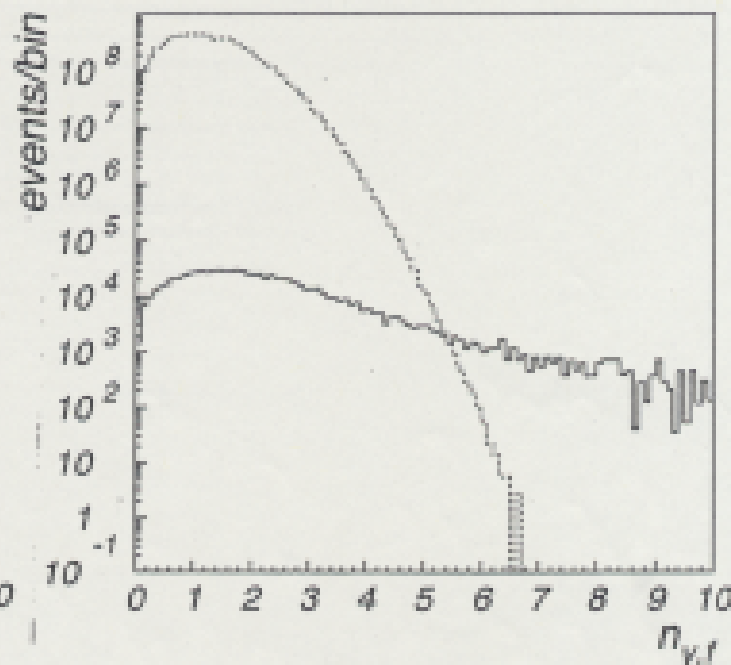
sextupole's effect

$n_{x,f}, n_{y,f}$ distributions

($n_{x,f}, n_{y,f}$: number of sigma at the exit)



$n_{x,f}$ distribution



$n_{y,f}$ distribution

(solid line: scattered electrons)
(dashed line: the others)

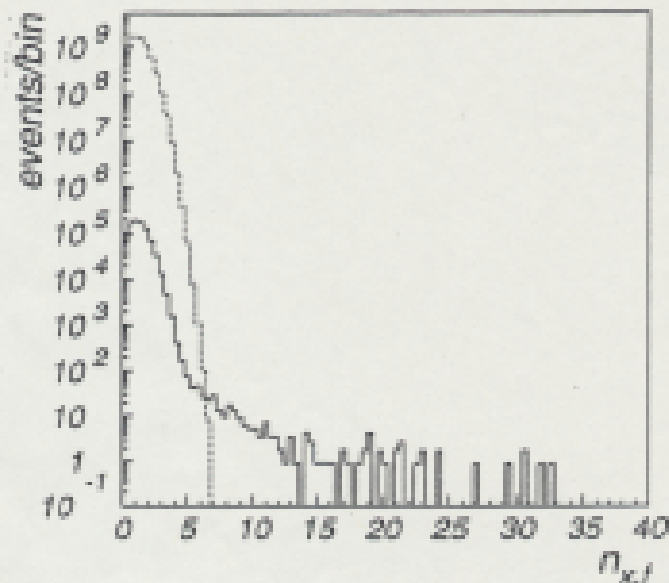
Long tails are produced in the distributions.

$$N(n_{x,f} > 7) = 160$$

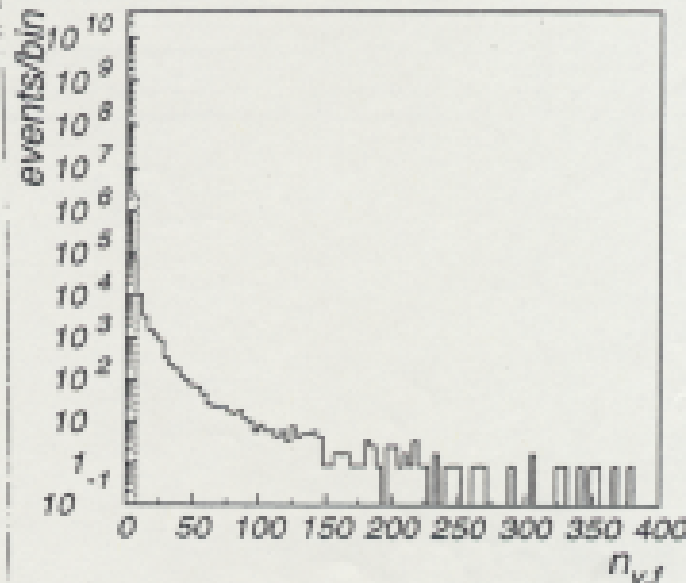
$$N(n_{y,f} > 7) = 24125$$

Tails

(solid line: scattered electrons
dashed line; the others)



$n_{x,f}$ distribution



$n_{y,f}$ distribution

The tails extend to

$$\sim \begin{cases} 30-40 \sigma & \text{(for } n_x) \\ 300-400 \sigma & \text{(for } n_y) \end{cases}$$

Table; $N(n_{x,f} > n_{\min})$ and $N(n_{y,f} > n_{\min})$

n_{\min}	10	50	100
$N(n_{x,f} > n_{\min})$	75	7	2
$N(n_{y,f} > n_{\min})$	24125	442	127

Emittance Blowup due to Elastic Scattering

Since the scattering probability is very small ($\sim 10^{-4}$), the effect on the whole beam does not become large.

	unscattered electrons	scattered electrons	all electrons
Number of electrons	7.50×10^9	6.02×10^5	7.50×10^9
$\gamma \beta \epsilon_x$ (m-rad)	3.06×10^{-6}	3.14×10^{-6}	3.06×10^{-6}
$\gamma \beta \epsilon_y$ (m-rad)	2.03×10^{-8}	6.00×10^{-8}	2.03×10^{-8}

Emittance-blowup vs. P_{linac} (P_{linac} : pressure in the linac)

The effect is negligible for $P_{\text{linac}} \lesssim 10^{-6} - 10^{-5}$ Torr

$\gamma \beta \epsilon_x$
(m-rad)

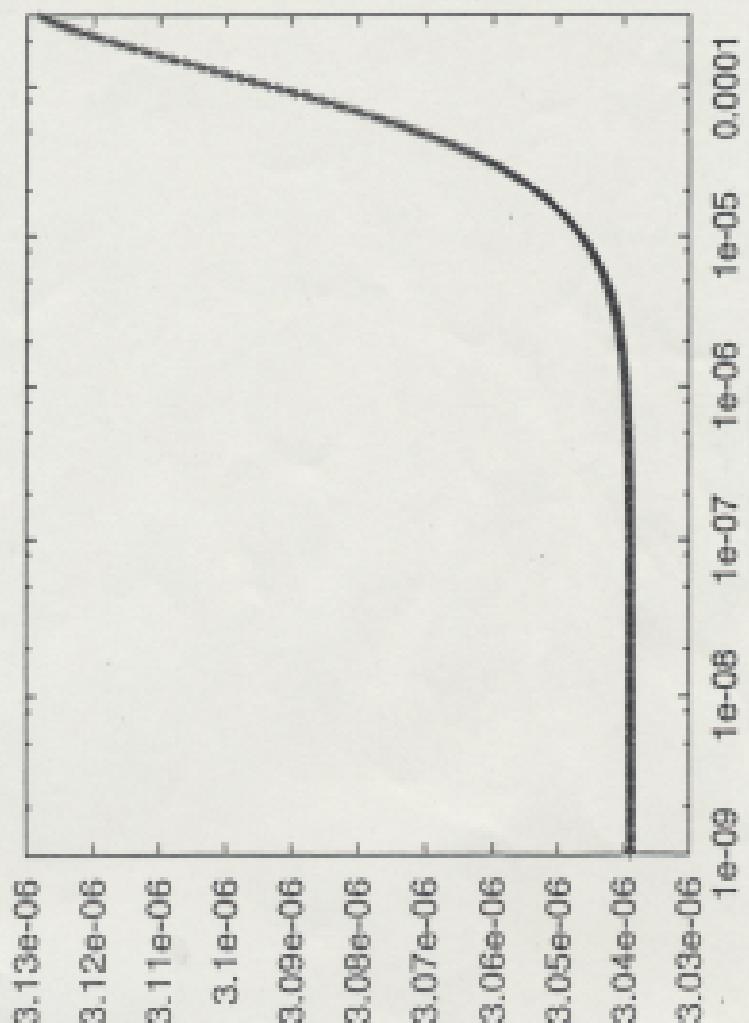
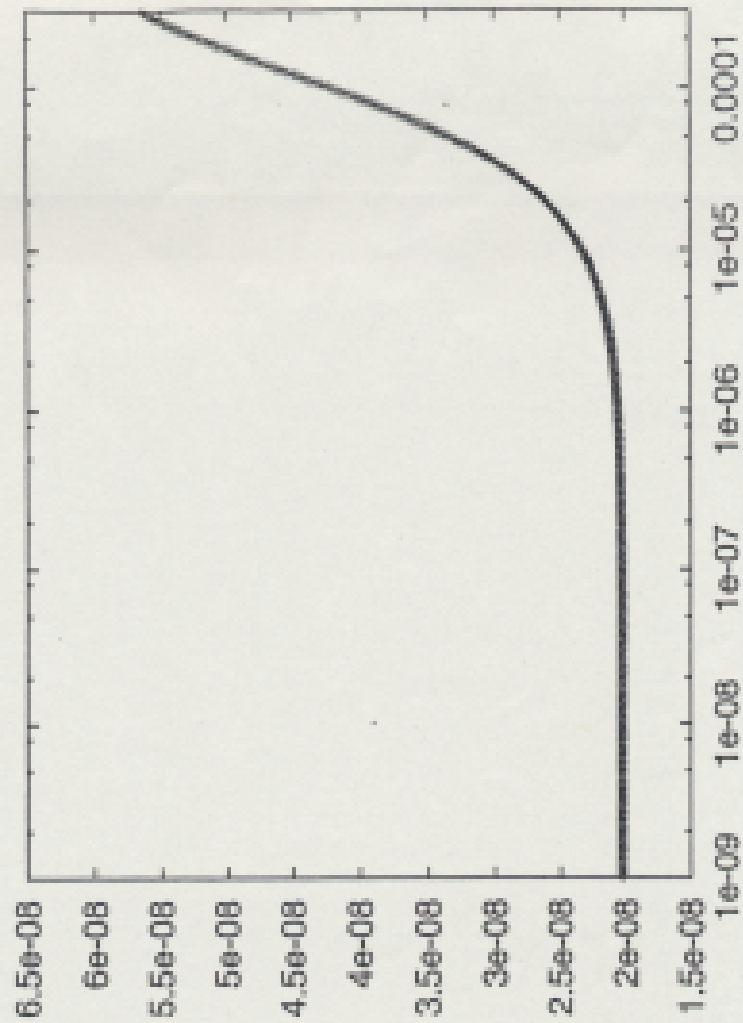


Fig: $\gamma \beta \epsilon_x$ vs P_{linac}

$\gamma \beta \epsilon_y$
(m-rad)



P_{linac} (Torr)

Fig: $\gamma \beta \epsilon_y$ vs P_{linac}

Summary

Simulation for the Electron Beam in the Main Linac

• Effect of sextupole magnet

Scale of the orbit spread : $\sim O(1)$ nm

• Effect of elastic scattering

Although the scale of the orbit spread reaches μm -order, the scattering probability is very small ($\sim 10^{-4}$).



Neither of them cannot have a large contribution

to emittance-blowup.