

Perspectives of direct searches of neutrino mass
with
very low energy beta spectroscopy of Re-187

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Three ways for measuring $m_{\nu i}$



- β decay: $m_i^2 \neq 0$ can affect spectrum endpoint. Sensitive to the "effective electron neutrino mass":

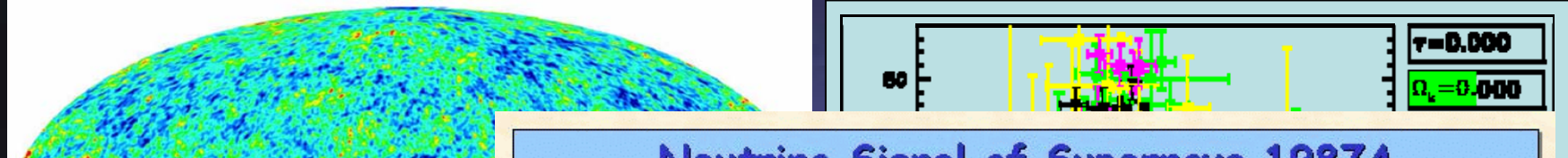
$$m_\beta = [c_{13}^2 c_{12}^2 m_1^2 + c_{13}^2 s_{12}^2 m_2^2 + s_{13}^2 m_3^2]^{\frac{1}{2}}$$

- $0\nu 2\beta$ decay: Can occur if $m_i^2 \neq 0$ and $\nu = \bar{\nu}$. Sensitive to the "effective Majorana mass" (and phases):

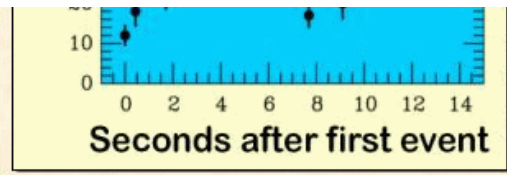
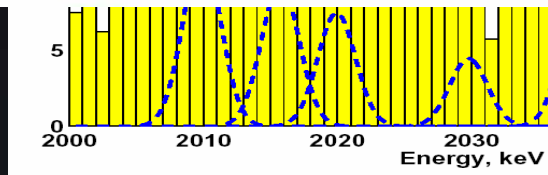
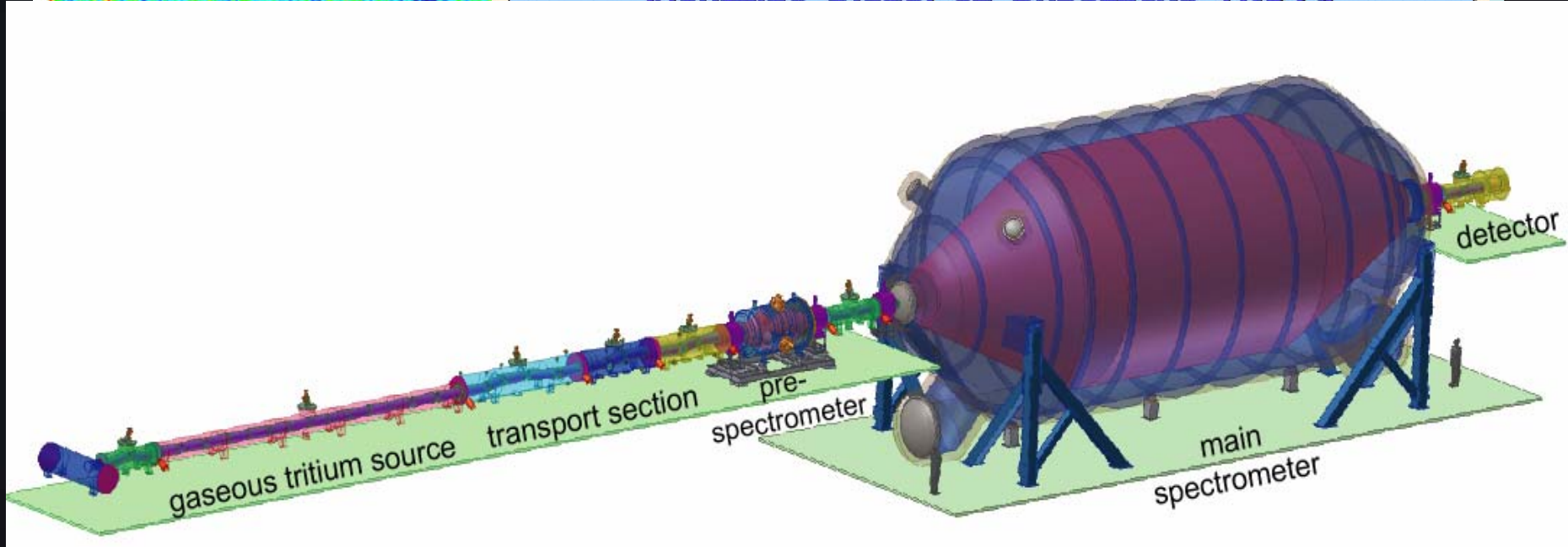
$$m_{\beta\beta} = |c_{13}^2 c_{12}^2 m_1 + c_{13}^2 s_{12}^2 m_2 e^{i\phi_2} + s_{13}^2 m_3 e^{i\phi_3}|$$

- Cosmology: $m_i^2 \neq 0$ can affect large scale structures in (standard) cosmology constrained by CMB+other data. Sensitive to:

$$\Sigma = m_1 + m_2 + m_3$$



Neutrino Signal of Supernova 1987A



Georg Raffelt, Mainz-Planck

The searches via a β decay

◆ $(A,Z) \rightarrow (A,Z+1) + e + \nu \quad (+ Q)$

◆ $dN(\varepsilon) = G^2 F(Z,\varepsilon) |M|^2 S(\varepsilon) p \varepsilon (\varepsilon - \varepsilon_0)^2 d\varepsilon$

⇓

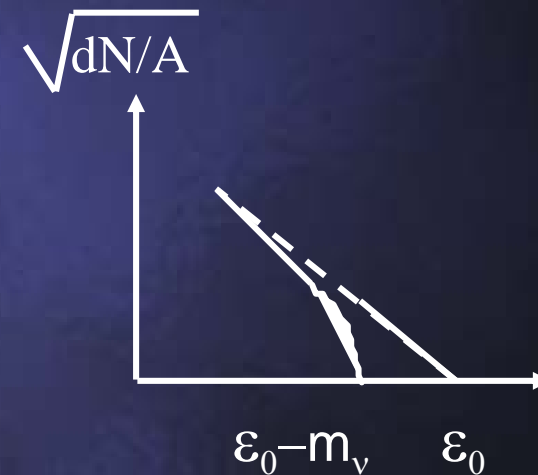
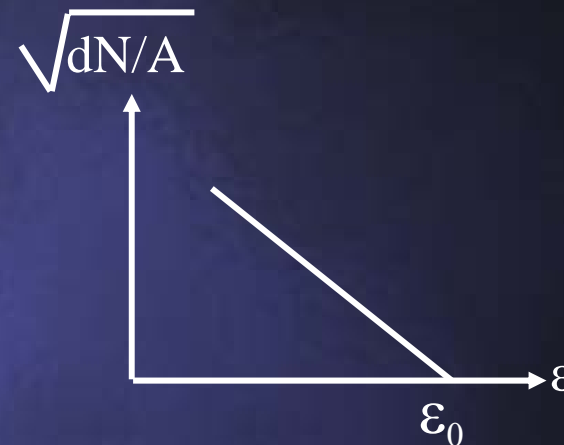
$\varepsilon \rightarrow \varepsilon_0$

$dN(\varepsilon) = A (\varepsilon - \varepsilon_0)^2 d\varepsilon$

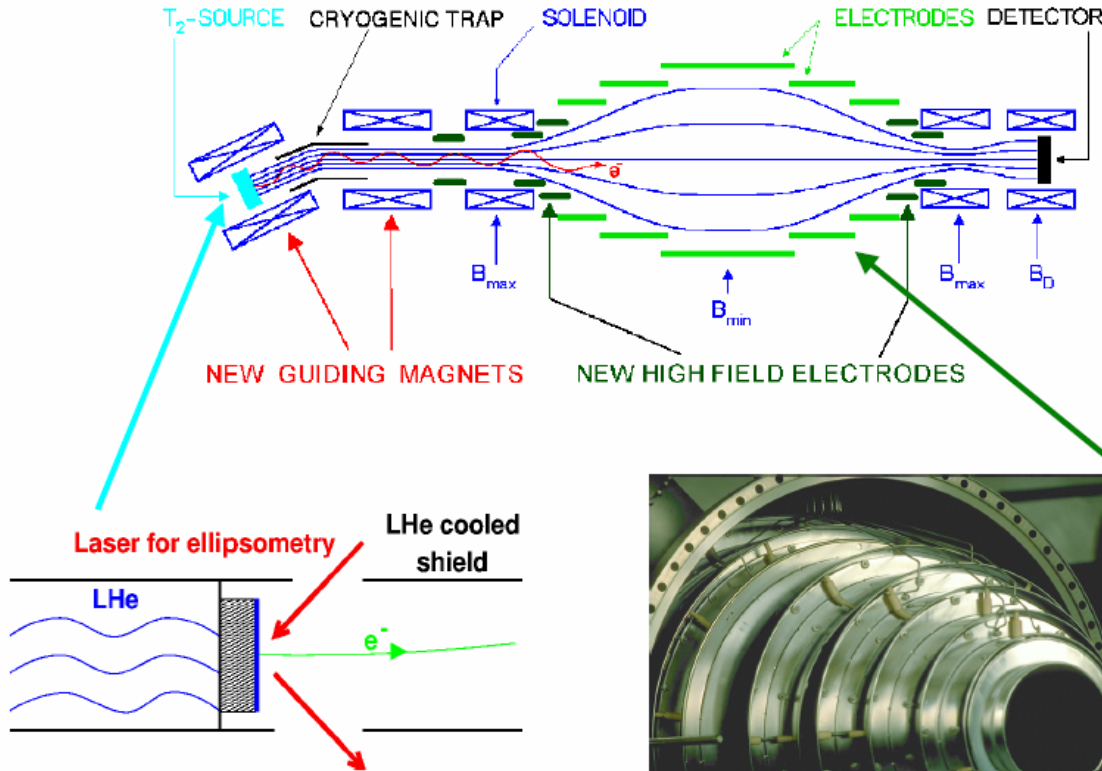
◆ $dN(\varepsilon) = A (\varepsilon - \varepsilon_0)^2 [1 - m_\nu^2 c^4 / (\varepsilon - \varepsilon_0)^2]^{1/2} d\varepsilon$

$\Rightarrow \varepsilon_{\max} = \varepsilon_0 - m_\nu c^2$

$\Rightarrow \text{deficit} \propto [m_\nu c^2 / \varepsilon_0]^3$



1st method: beta impulse spectroscopy



- T₂ Film at 1.86 K
- quench-condensed on graphite (HOPG)
- 45 nm thick (≈130ML), area 2cm²
- Thickness determination by ellipsometry



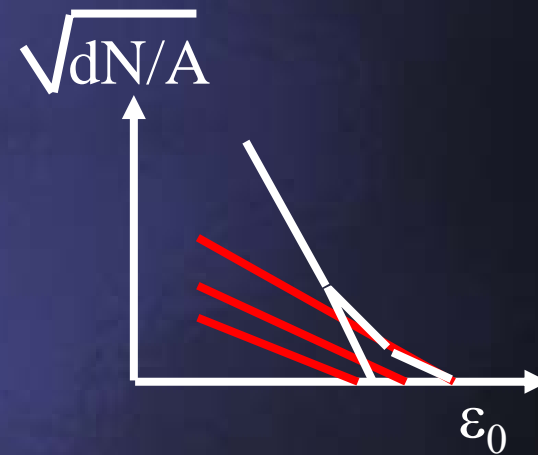
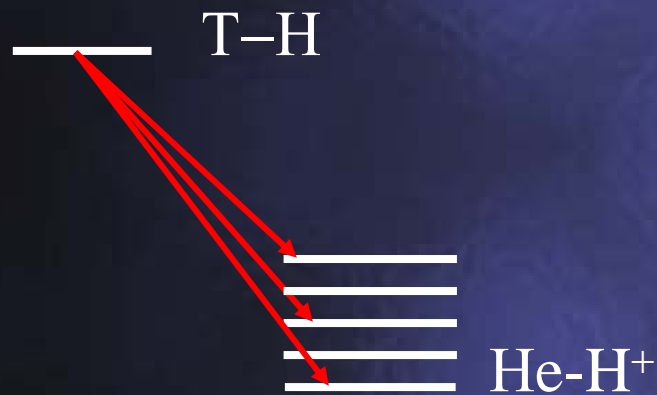
Mainz
v group
2001:

J. Bonn
B. Borschein*
L. Borschein
B. Flatt
Ch. Kraus
B. Müller
E.W. Otten
J.P. Schall
Th. Thümmler**
Ch. Weinheimer**

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- Systematic unknown effects affect the shape at the end-point: in particular the so called "final states effect"



2nd method: calorimetric beta spectroscopy



Advantages:

- Measurement of whole energy of the decay

$$E_i = \varepsilon_i + \Delta_i$$

$$\Rightarrow dN(E) = A \sum_i w_i (E_i - E_0)^2 dE$$

\Rightarrow no model dependent corrections for atomic and molecular final states.

\Rightarrow no correction for nuclear recoil energy and for electron energy losses, ...

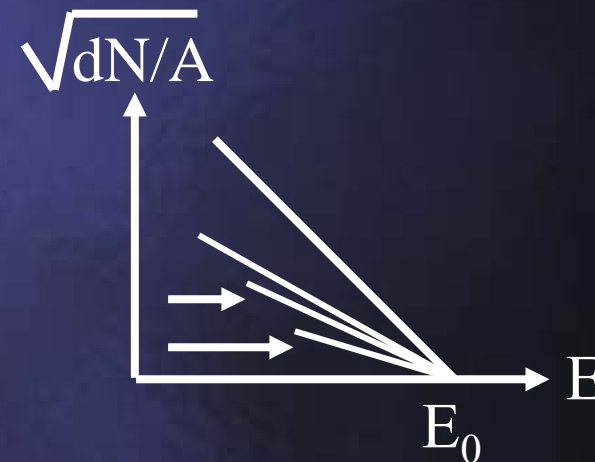
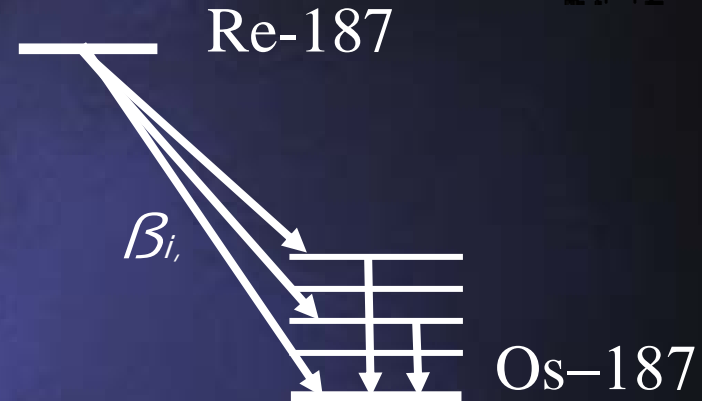
Disadvantages:

- Beta Source inside the detector
 \Rightarrow all spectrum must be acquired: but interesting area proportional to

$$[m_\nu c^2 / E_0]^3$$

\Rightarrow Re¹⁸⁷: lowest Q ~ 2.5 keV.

\Rightarrow Re¹⁸⁷: $[m_\nu c^2 / E_0]^3 \sim 1/400$ of H³



Few Historical hints



- 1985 - First concept of determining neutrino mass (S. Vitalone)
- Re-187 properties re-determined (S. Vitalone)
- Only two measurements of β spectrum in 1967

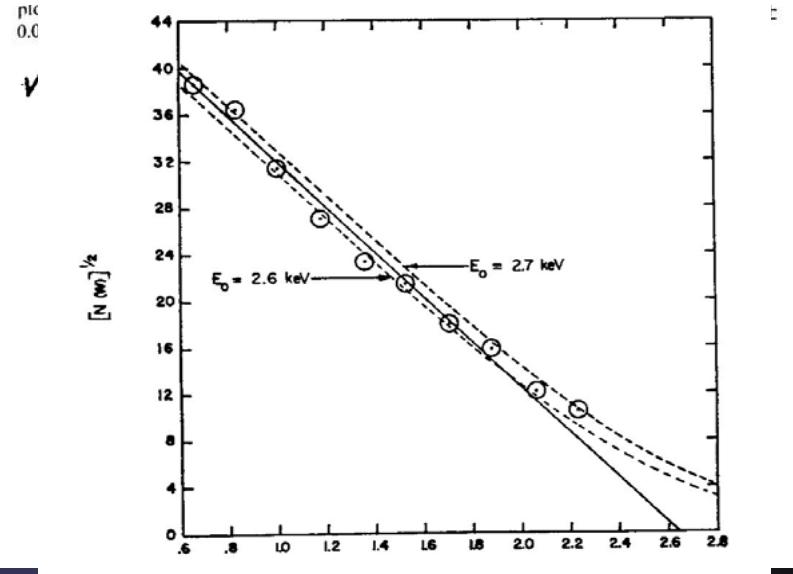


PHYSICAL REVIEW VOLUME 138, NUMBER 6B 21 JUNE 1968

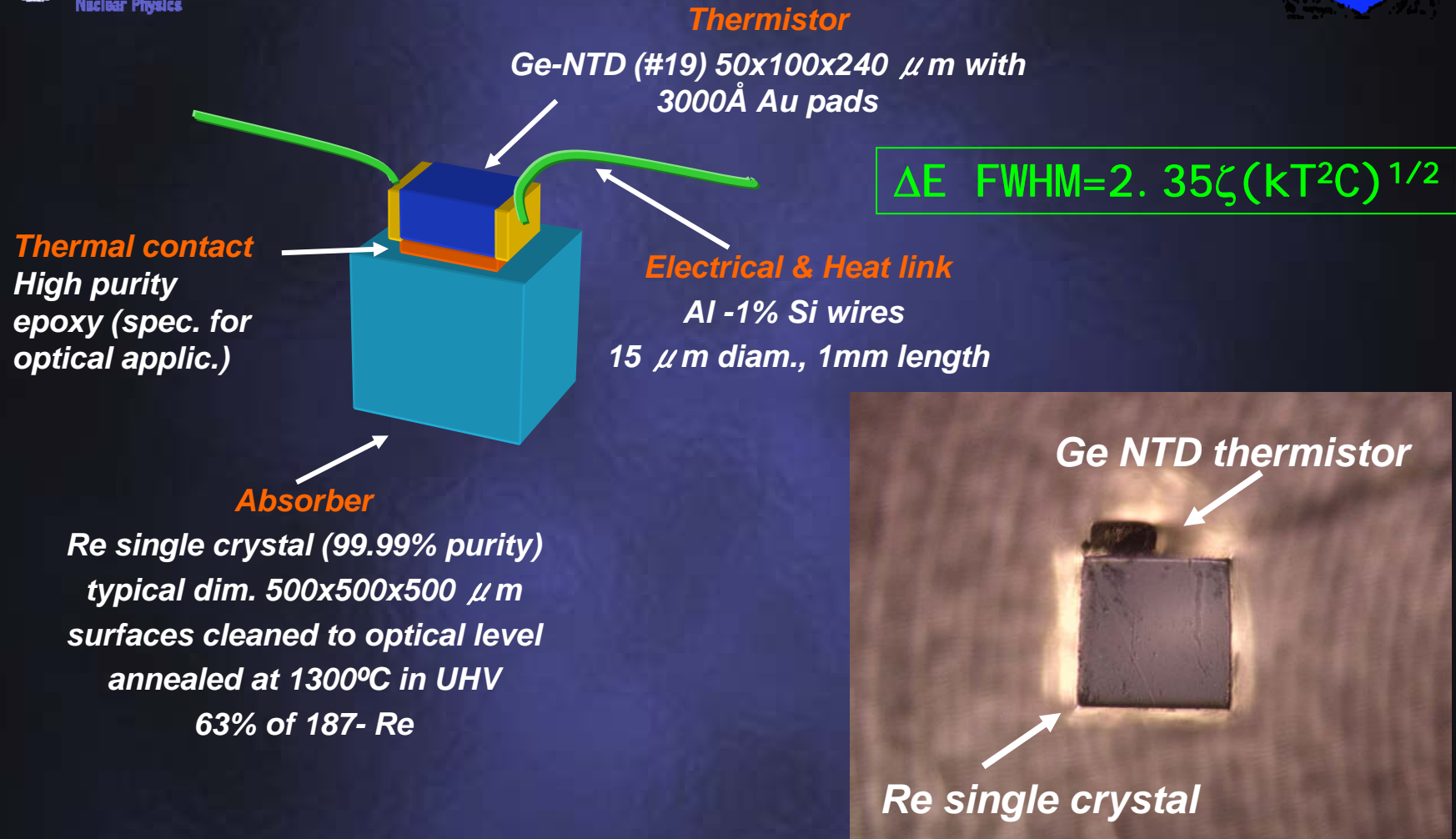
Decay of Rhenium-187

R. L. BRODZINSKI AND D. C. CONWAY*
 Department of Chemistry, Purdue University, Lafayette, Indiana
 (Received 21 September 1964; revised manuscript received 23 February 1965)

The end-point energy for ^{187}Re has been determined to be 2.62 ± 0.09 keV by use of $(\text{C}_2\text{H}_5)_2\text{ReH}$ vapor in a proportional-counter spectrometer. The ^{187}Re half-life determined from the proportional-counting experiments is $(6.6 \pm 1.3) \times 10^{10}$ year. This is identified as the β^- half-life for decay to continuum states only. When this is combined with the total half-life of $(4.3 \pm 0.5) \times 10^{10}$ year obtained from $^{187}\text{Re}/^{187}\text{Os}$ ratios in geologically dated minerals, the ratio of bound-state to continuum-state decays is found to be 0.5 ± 0.3 .

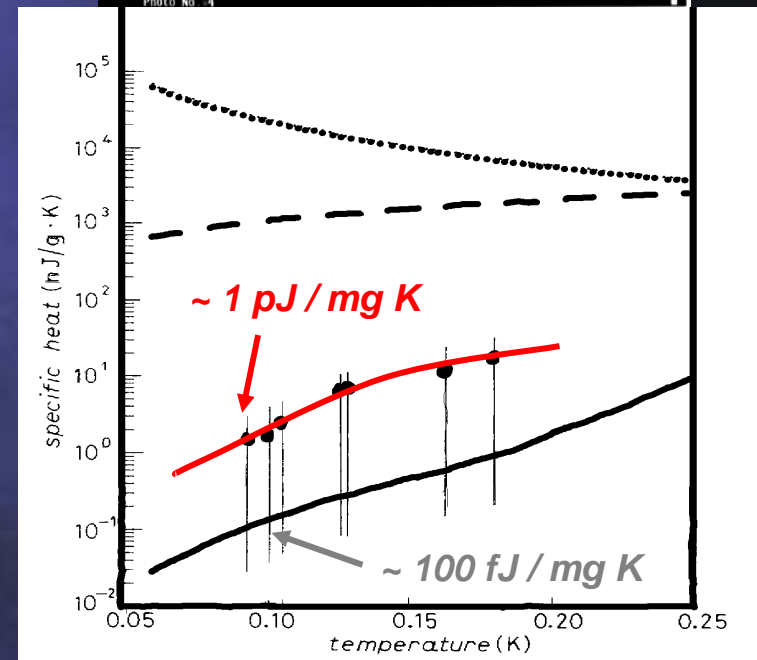
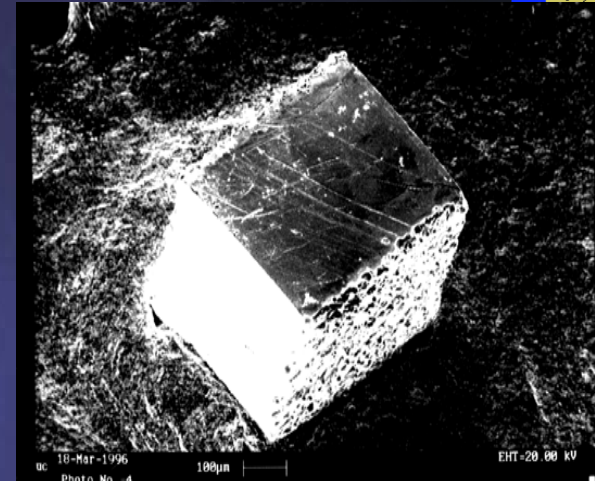


Cryogenic microcalorimeter for β decay



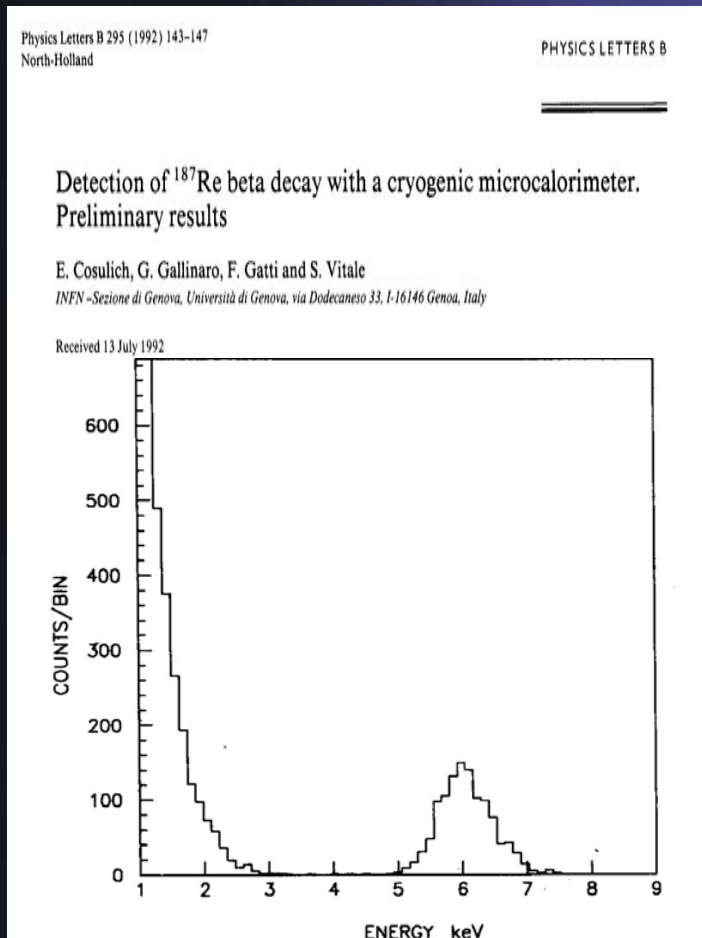
Re properties

- Re metallic superconductor
 - ◆ HCP lattice
 - ◆ $T_c = 1.69$ K
 - ◆ $\rho = 21$ g/cm³
 - ◆ $T(\text{Debye}) = 460$ K
 - ◆ M.P. = 3000 K
 - ◆ $Z = 75$
 - ◆ $A = 185(37\%), 187(63\%)$
 - ◆ $\tau(1/2) \text{ Re-187} = 4 \times 10^{10}$ y
- Specific heat has a Shottky term $1/T^2$, not measured for short thermal pulses in superconducting state
- Specific heat seems contains a small contribution of normal electrons (about 1/1000) measured exciting with ionizing particles.
- Very small electron escape depth L (Ex: $L < 100 \text{ \AA}$ @ 1.6 keV) \Rightarrow negligible corrections for energy escape

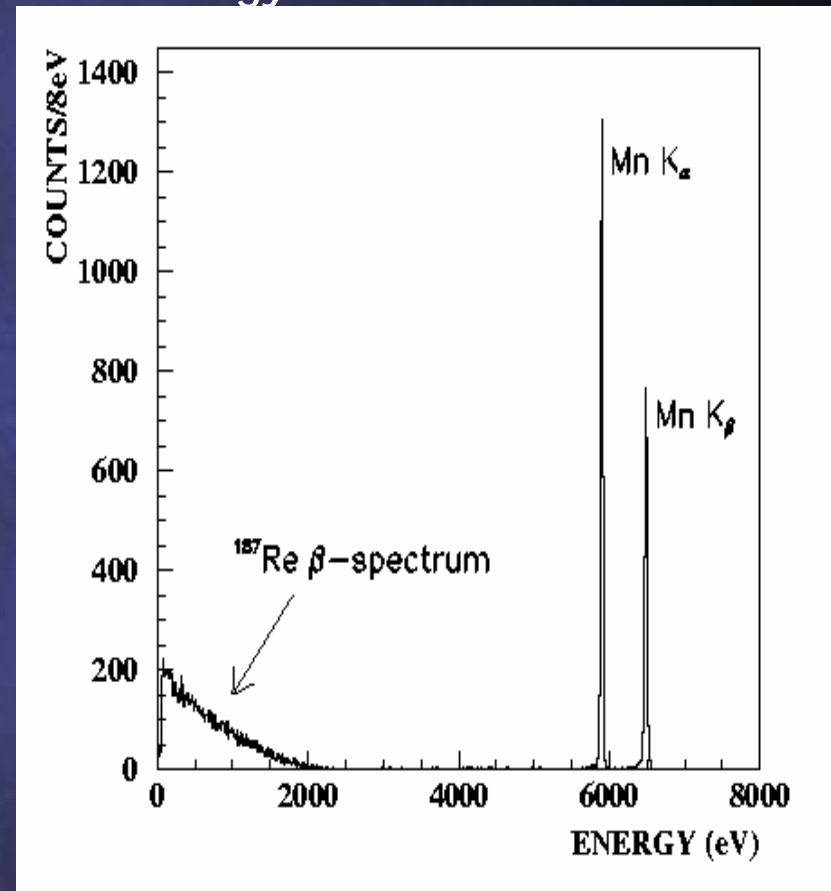


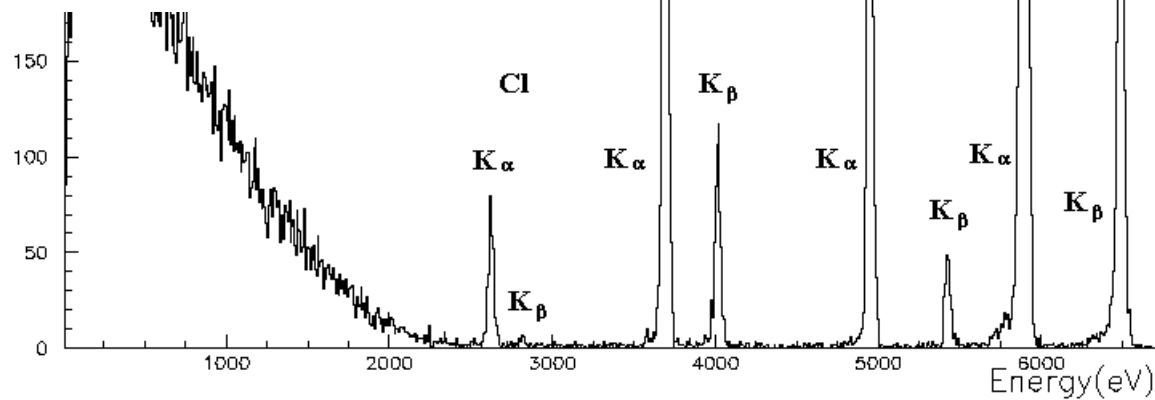
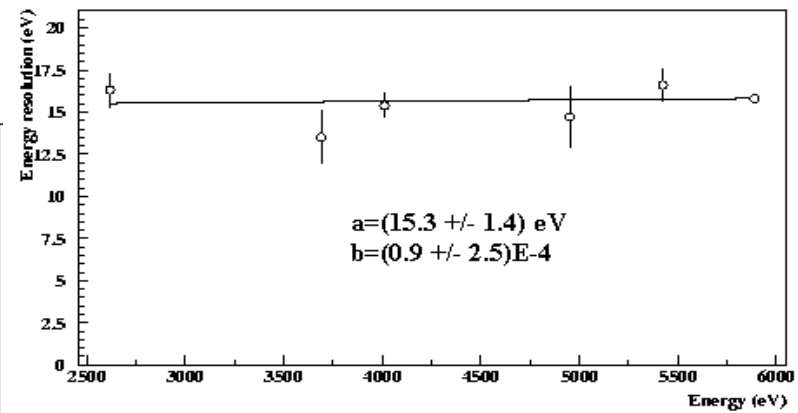
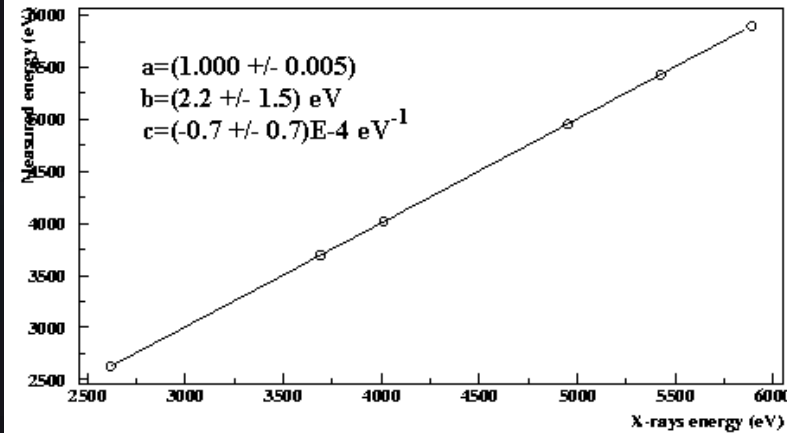
First steps

- 1992 (660 eV FWHM)

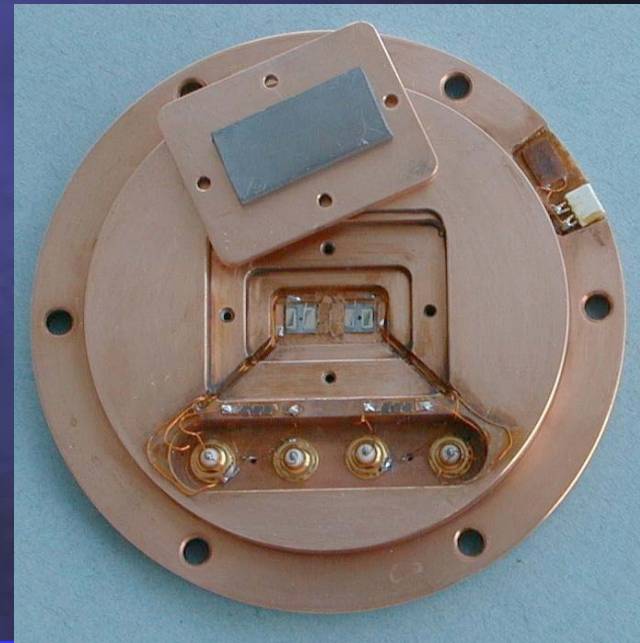
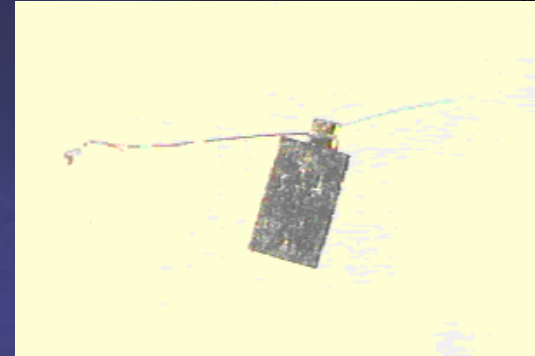
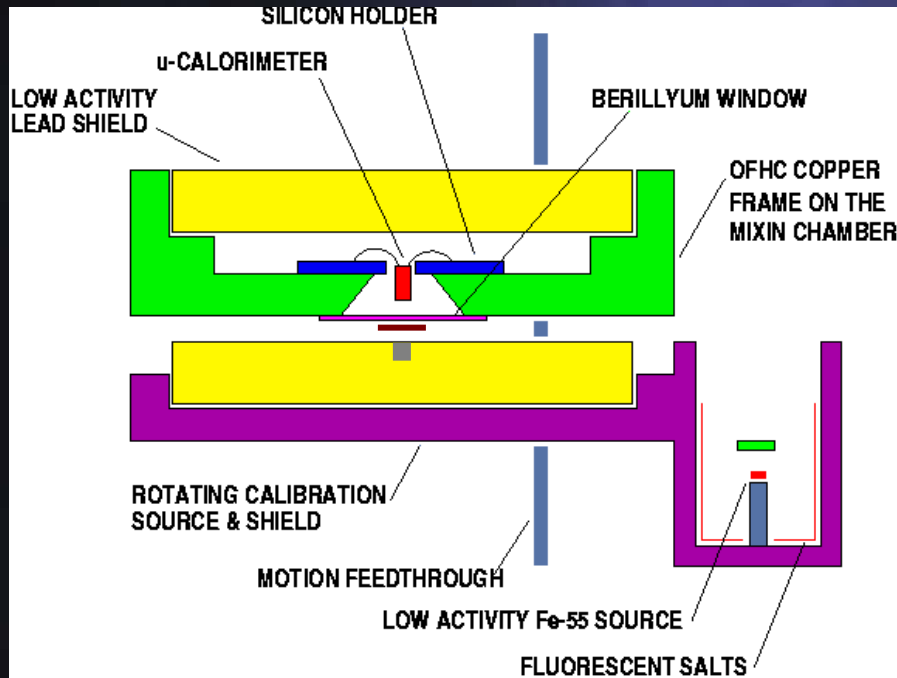


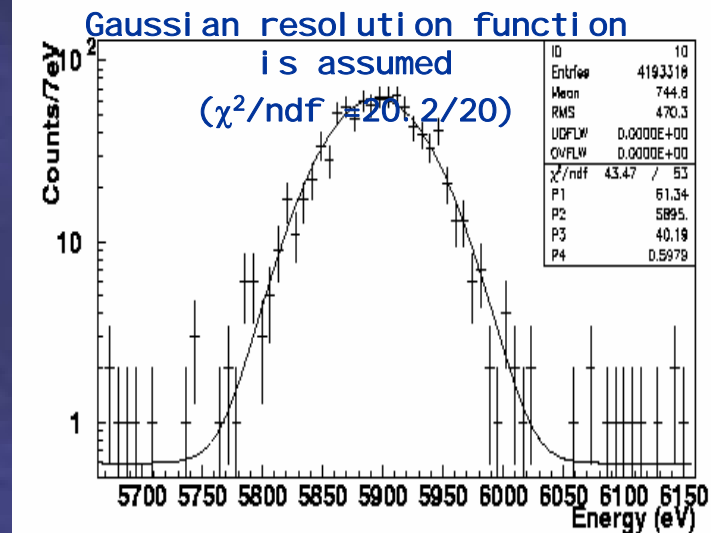
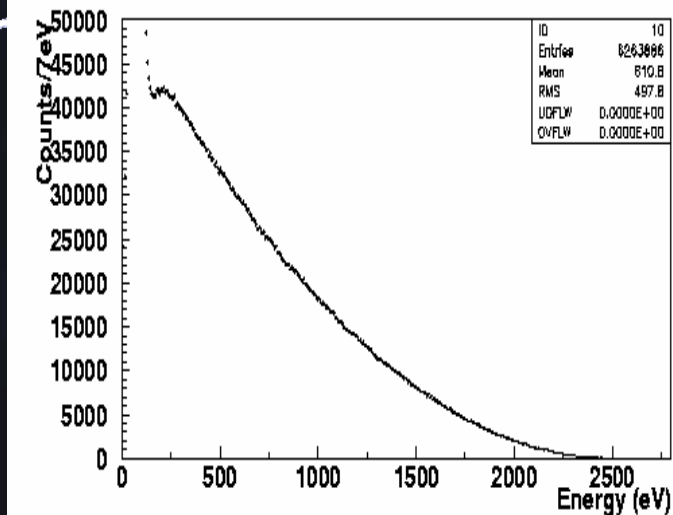
- 1996 (30 eV FWHM)
Energy Threshold = 50 eV





First pilot experiment ('99)

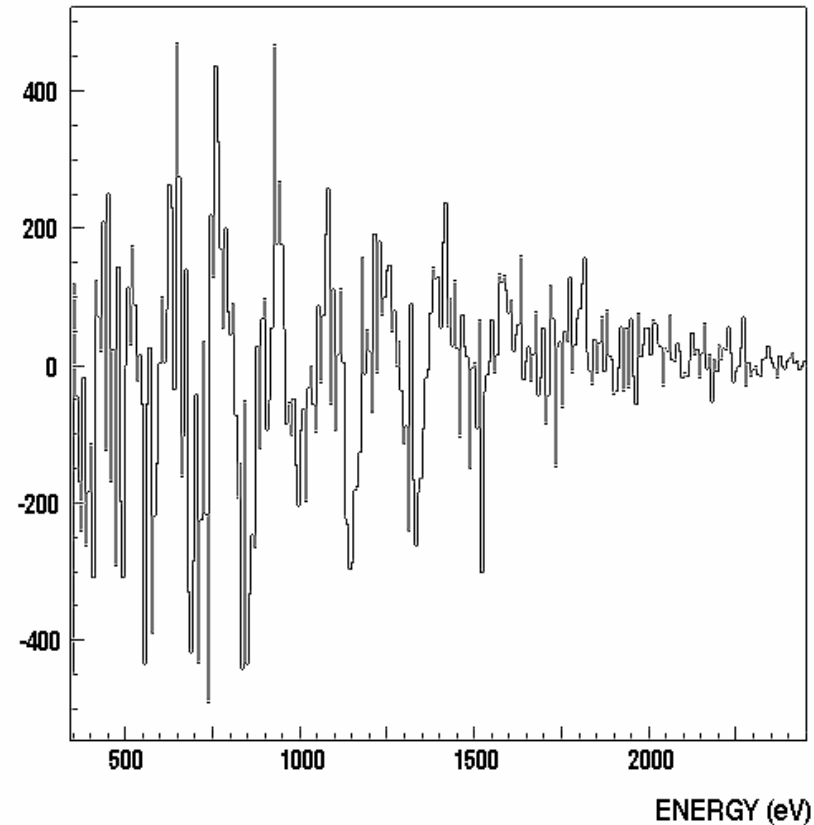




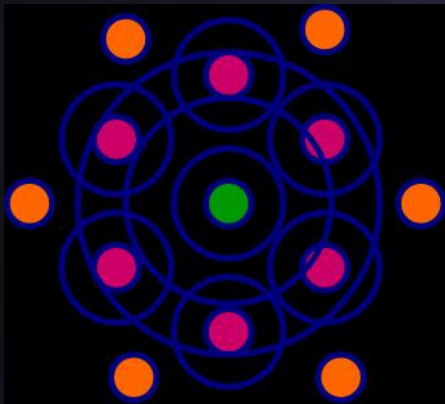
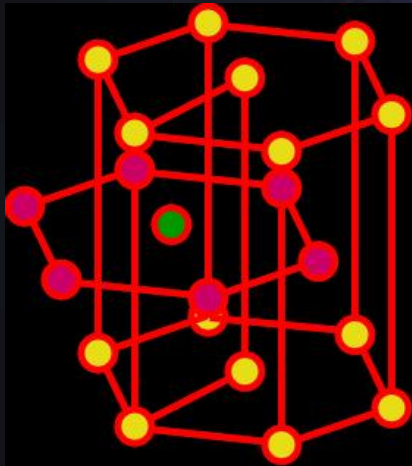
- ✓ End point = $(2470 \pm 1 \text{ stat} \pm 4 \text{ sys}) \text{ eV}$
- ✓ Half life = $(4.12 \pm 0.02 \text{ stat} \pm 0.11 \text{ sys}) 10^{10} \text{ y}$
M. Galeazzi, et al. Phys Rev C, 63 (2000)014302.
- ✓ $m_\nu^2 = -462^{+579}_{-679} \text{ eV}^2/c^4$
- ✓ $m_\nu < 26 \text{ eV}/c^2$ 95% CL, $19 \text{ eV}/c^2$ 95% CL
F. Gatti, Nucl Phys B,
- ✓ Improvement of the present limits on massive neutrino admixture
M. Galeazzi, et al. Phys Rev Lett 86 (2001) 1978.

But something unexplained

- The residuals of the beta spectra fit with the theoretical function show correlation well beyond the admitted statistic fluctuation.
- This has been found to be caused by a physical effect



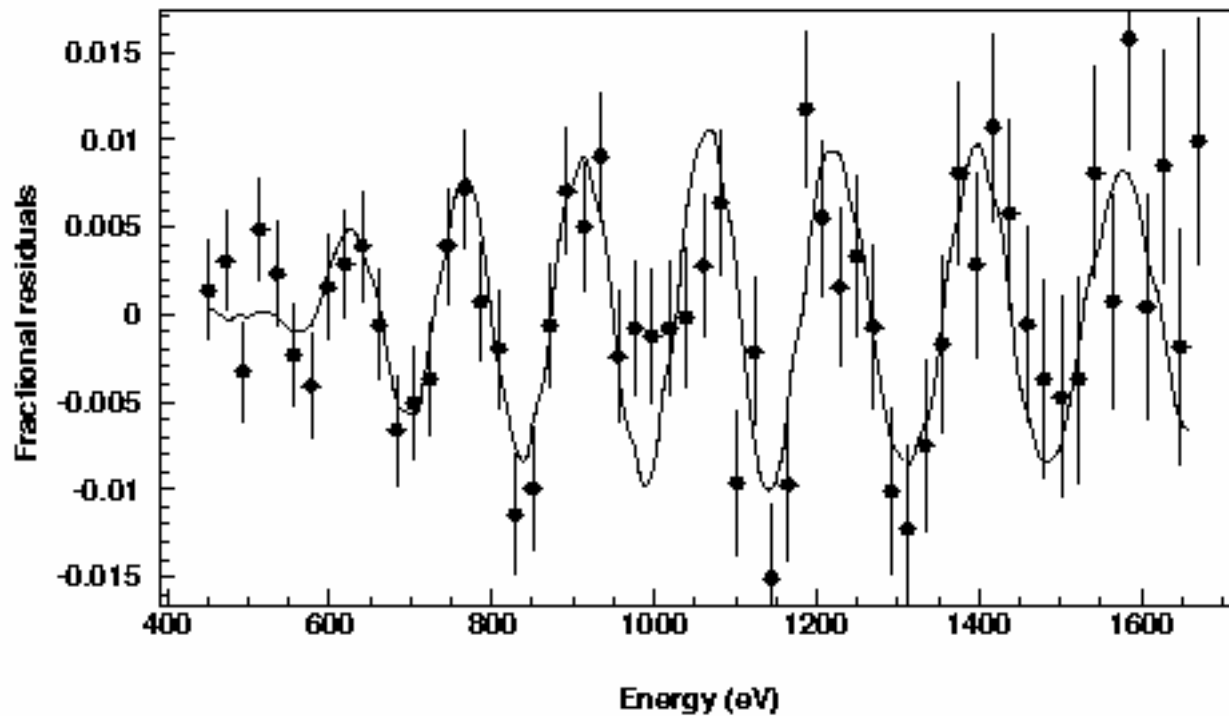
Why the Beta Spectrum Oscillates



- $E(\beta) \gg E(\text{Fermi}) \Rightarrow$ beta electrons interacts with atomic cores.
- $k^2 = 2m(E(\beta) - V)/\hbar^2$, $V \sim -15 \text{ eV}$,
 $\lambda(100) \sim 0.2 \text{ \AA}$, $\lambda(1000) \sim 0.04 \text{ \AA}$
 $a = 2.76 \text{ \AA}$, $c = 4.45 \text{ \AA}$, $c/a = 1.61 \text{ \AA}$ (1.63 \text{ \AA}).
- Self interference of outgoing and reflected waves from atomic shells:
 \Rightarrow (backscattering amplitude) \times (self-interference amplitude on Re nucleus from each atomic shell) \times (number of atoms of shell)
- Thermal motion energy: $T \rightarrow 0 \sim \exp(2k^2 / M\Theta_D)$
- β wave attenuation (“range”): $\exp(-\gamma R)$, $\gamma(\epsilon) \sim 3-20 \text{ \AA}$
- First hypothesis: S.E.Koonin in 91(Nature 354,486), never observed.

$$\chi(T_e) = \sum_i \frac{N_i |f_i(\pi)|}{kR_i^2} \sin(2kR_i + \phi_i + 2\delta_0) e^{-\gamma \cdot R_i} e^{-2\sigma_i^2 k^2}$$

Fine Structure of the Beta Decay



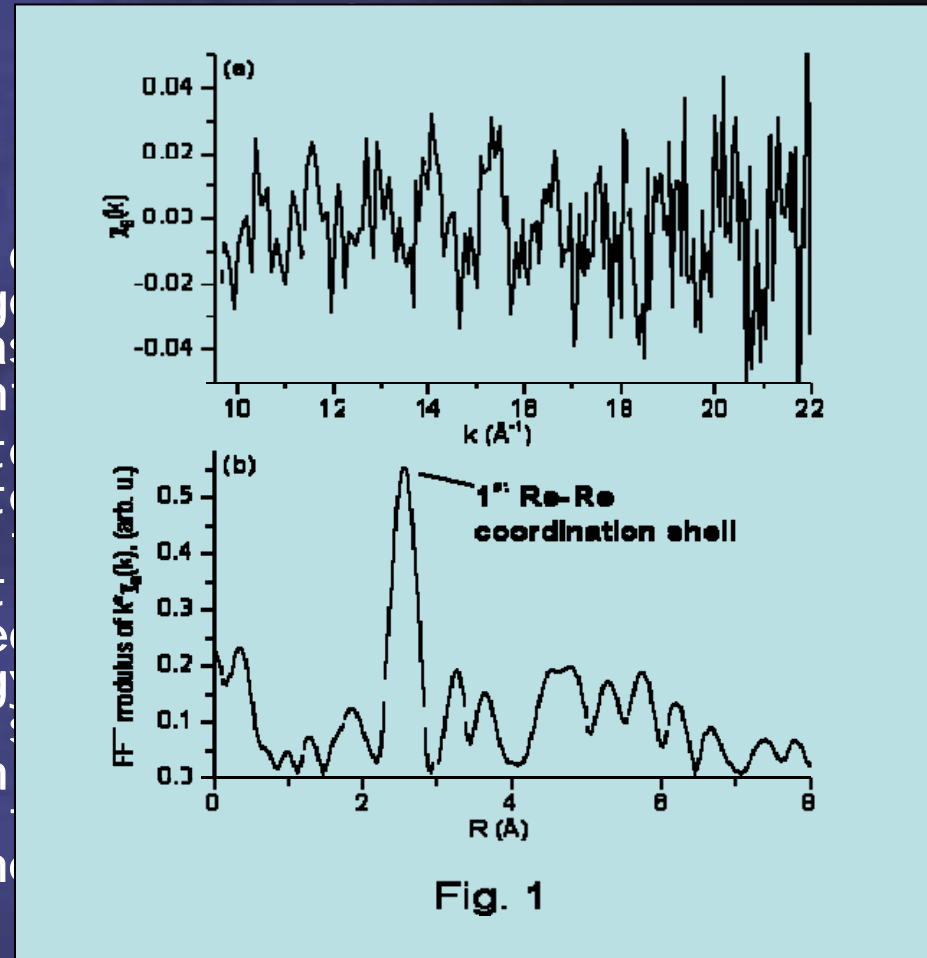
F. Gatti et. Al, NATURE, VOL 397, 14, jan, 1999

BEFS as probe of the crystal structure

BEFS

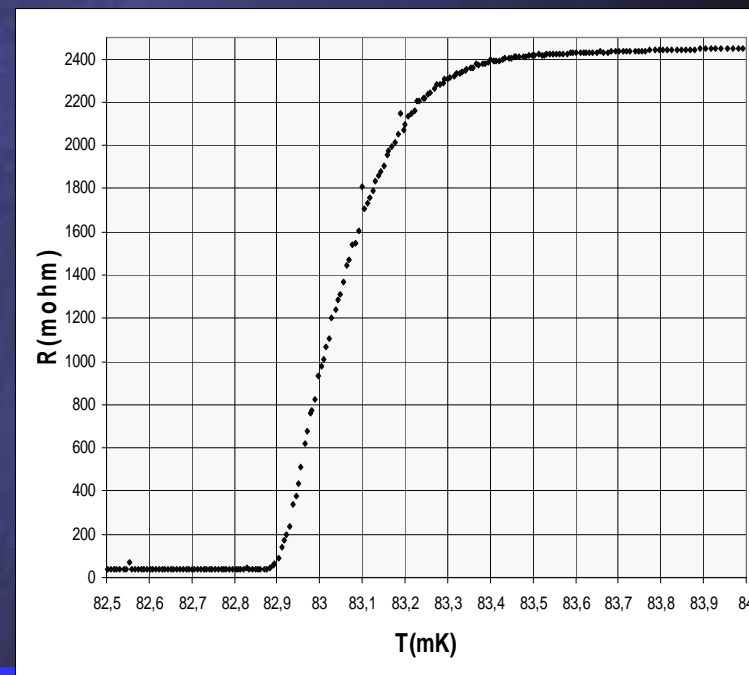
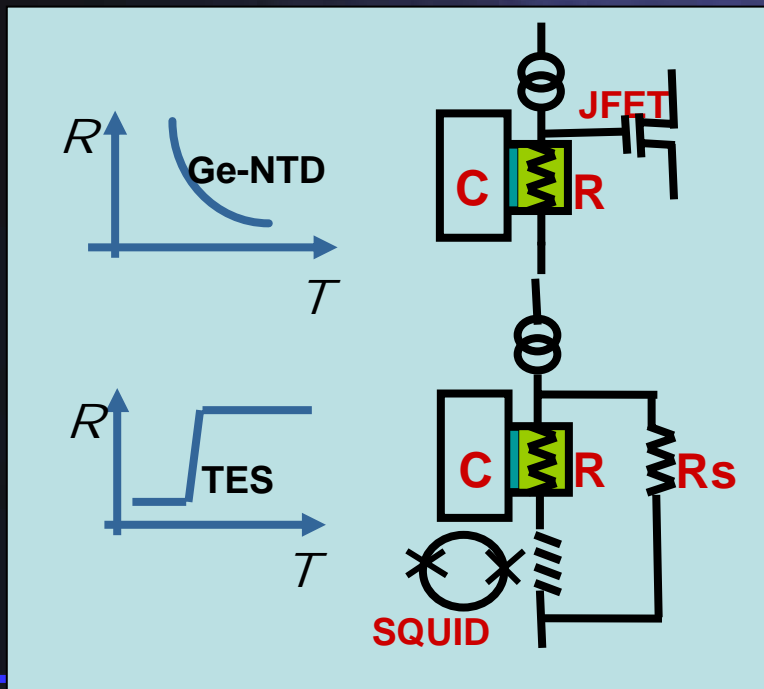
	th. phase ampl.	exp. phase ampl.
R_1 (Å)	2.76 ± 0.02	2.75 ± 0.02
R_2 (Å)	2.78 ± 0.02	

- The BEFS spectroscopy is a powerful tool for studying hydrogen-charged systems, a new interest in the last few years. The researches of efficient materials for hydrogen storage are of great interest.
- Indeed, the atomic site occupancy and the coordination number can be investigated efficiently using the BEFS spectroscopy. Actually this information can be obtained using EXAFS spectroscopy, but the electron binding energy of the core level is recorded of an EXAFS spectrum. The tiny cross section of the core level (hundreds of eV coming from the core level) rise in general to a noise level.

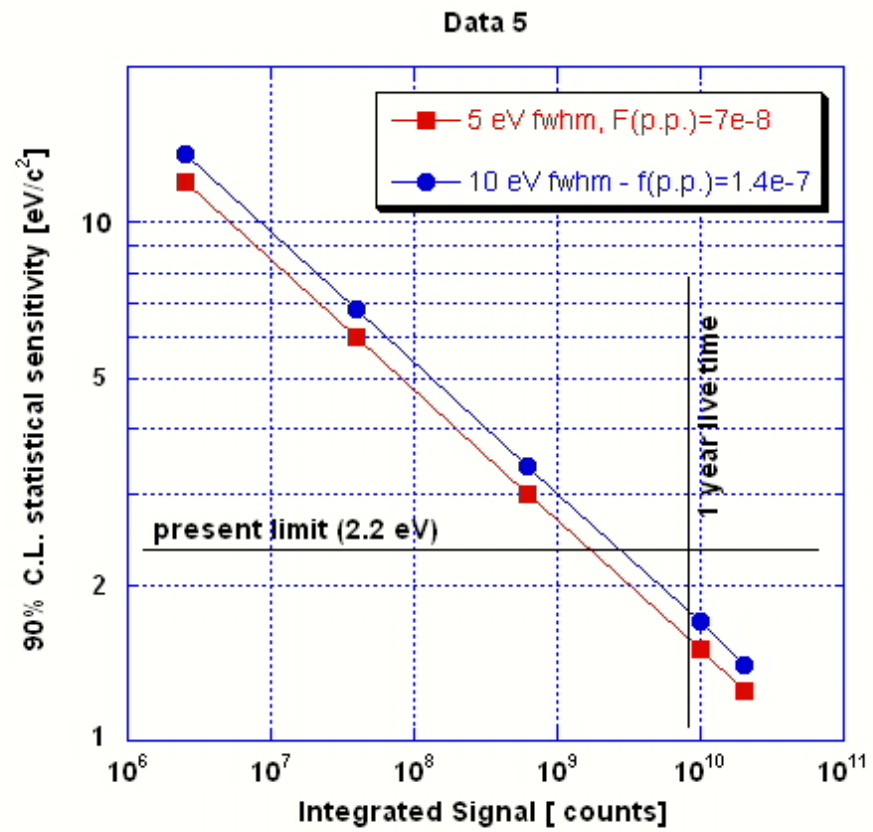
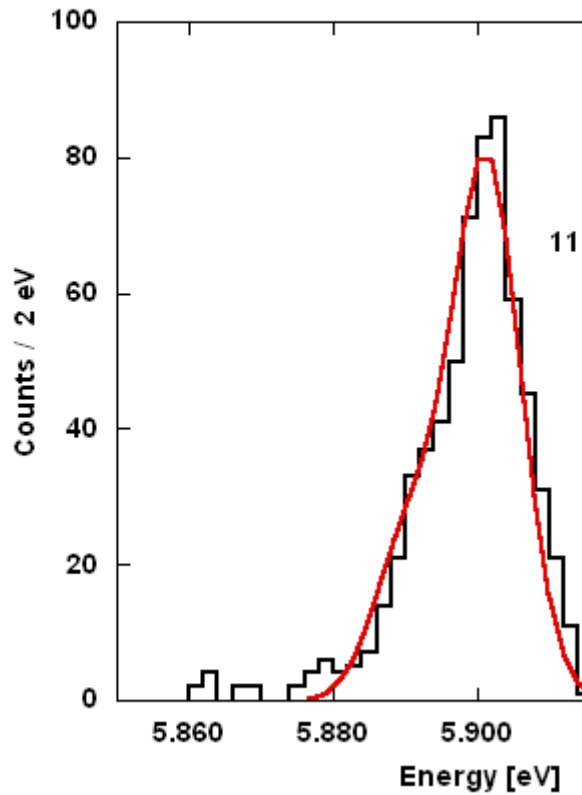


A second generation experiment for neutrino mass search

- New Technology for sensor we have been developed in the last years allow to perform an experiment with about 1 eV/c² sensitivity.
- The result of Mainz & Troitszk can be reproduced with a different methodology and hopefully improved.



A 300 Re detectors experiment



Next Generation Experiment (proposal MARE) down to 0.1 eV 2

